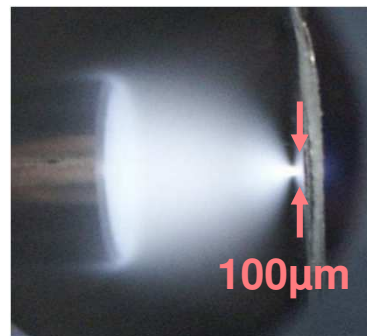


Cours Microplasmas: physique et applications

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Outline of the lecture

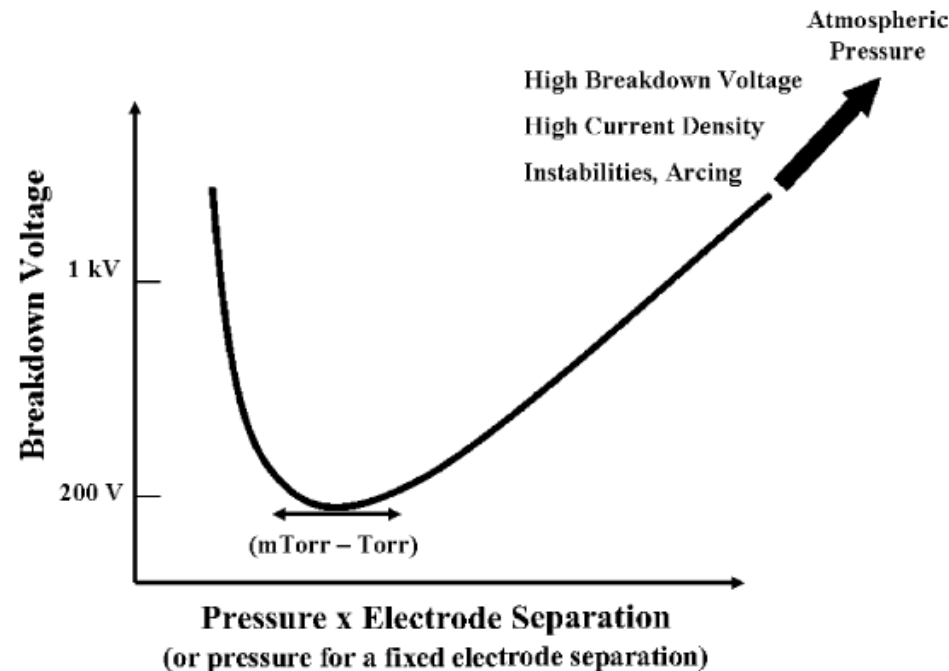


- Microplasmas: definition, interest and drawback
- Microplasma configurations
- Experimental characterization
Parameters of interest: T_g , n_e , EEDF, reactive species density
- Microplasma modeling
PIC simulation, fluid model, global model
- Applications of microplasmas: material synthesis and environment

Microplasmas: definition/applications



- Microplasma = plasma with submillimeter scale
- First works in the mid-90's (Schoenbach *et al.*)
- High pressure stable glow discharge



- Various applications: sterilization, treatment of human skin, light sources, micro-propulsion, synthesis of nanomaterials..

Interest and drawback of microplasmas



- Stable glow discharge at high pressure (intermediate to atmospheric pressure)
- Production of high density of reactive species → interesting for various applications such as lighting, material synthesis or medicine

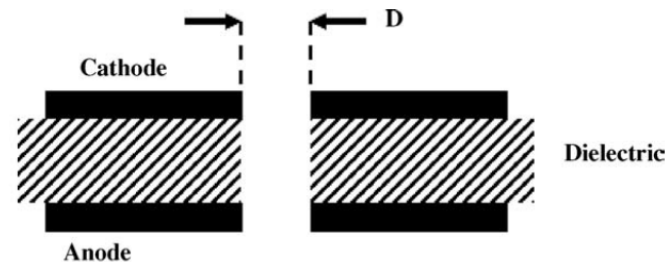
But working surfaces and volumes are weak..

- Solution: microplasmas arranged in arrays or third electrode in the case of MHCD (→ MCSD)

Microplasma configurations (1)

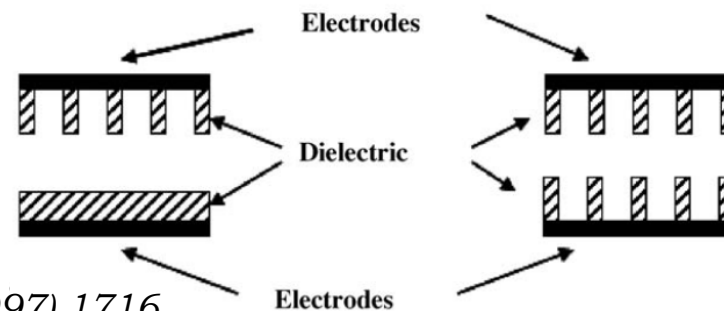


- Micro Hollow Cathode Discharge (MHCD, 1996,USA):



Schoenbach et al. APL **68** (1996) 13
Boeuf et al. APL **86** (2005) 071501
Aubert et al. PSST **16** (2007) 23

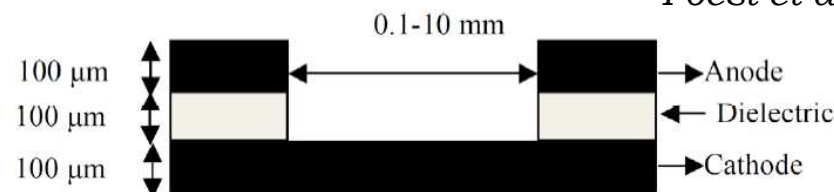
- Capillarity plasma electrode Discharge (CPED, 1997, USA):



- $D = 0,01-1$ mm
- $L/D = 10:1 - 1:1$
- Pulsed DC or AC voltages

Kunhardt et al. BAMS **42** (1997) 1716

- Cathode Boundary Layer (CBL, USA):



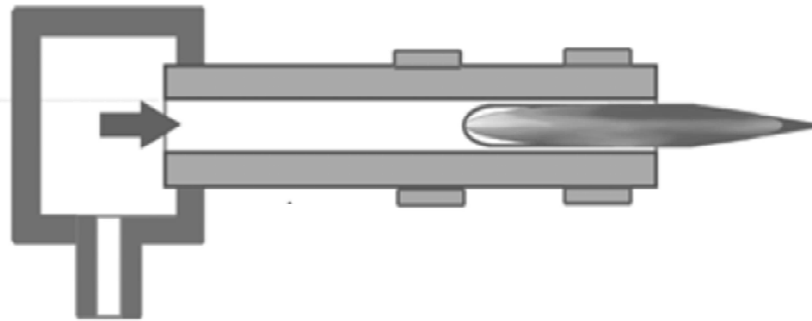
Foest et al. PPCF **47** (2005) B525

- Efficient source of excimer radiation
- Matrix easy to ignite

Microplasma configurations (2)



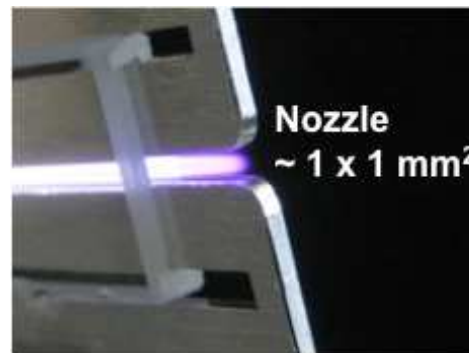
- Dielectric Barrier Discharge Microplasma (DBD, USA)



- RF excitation
- Surface modification

Moselhy et al. JAP 95 (2004) 1642

- Atmospheric pressure plasma jet (APPJ, Bochum/Belfast):



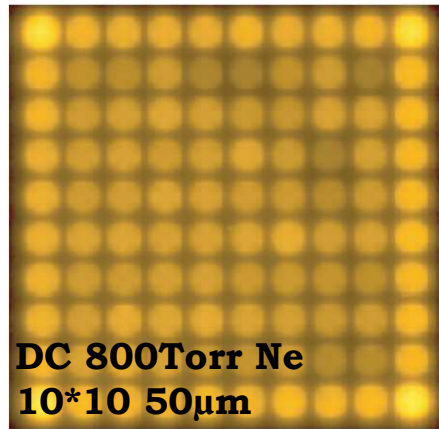
- RF excitation
- Electrode spacing : 1 mm

V. Schulz-von der Gathen et al., J. Phys. D 41 (2008) 194004

Increasing the plasma surface / volume

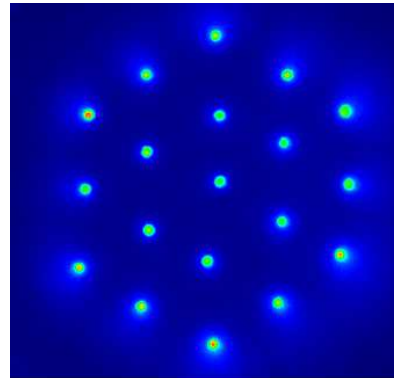


- Array of microplasmas:



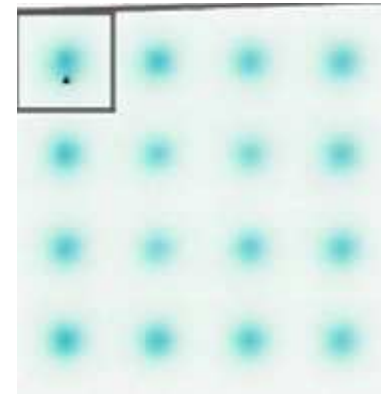
Eden et al. *JPhysD* **38** (2005) 1644

CBL
50mbar Ar
19*800µm



Martin et al. *ICPIG* 2012

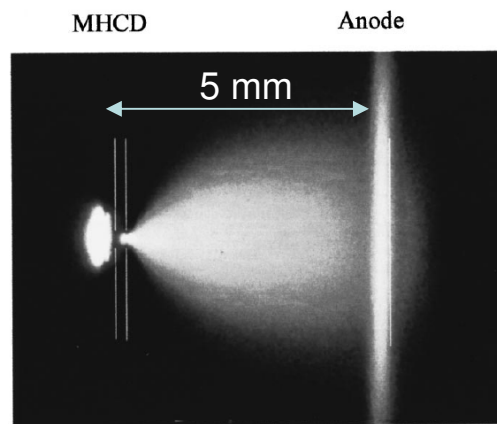
Si/SiO₂/Ni MHCD
350mbar He
16*150µm



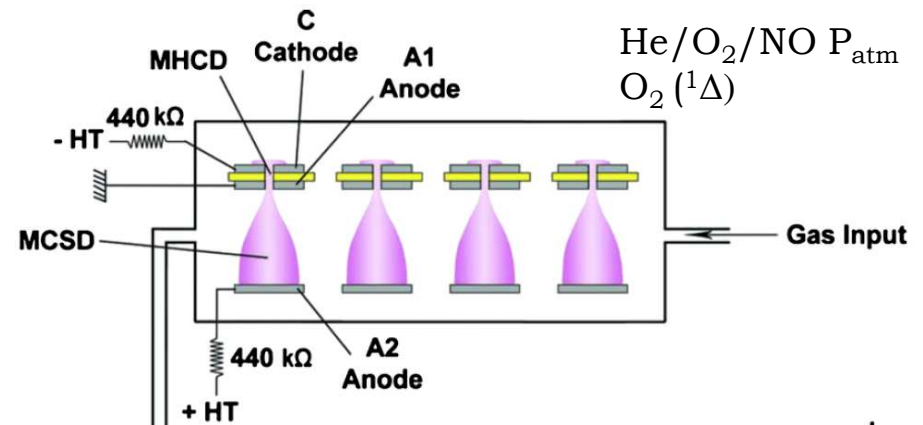
Felix et al. *PSST* **25** (2016) 025021

- MCSD

Ar
160 Torr



Eden et al. *JAP* **85** (1999) 2075
Makasheva et al. *PPCF* **49** (2007) B233



Santos Sousa et al. *APL* **97** (2010) 141502

Experimental characterization of microplasmas

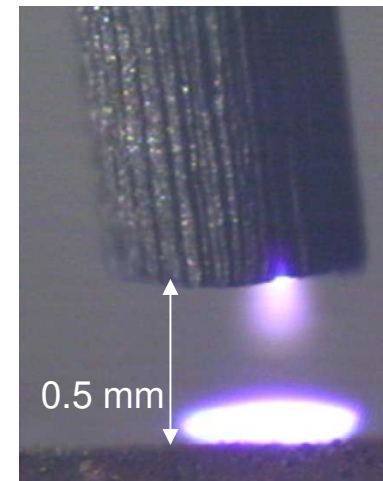
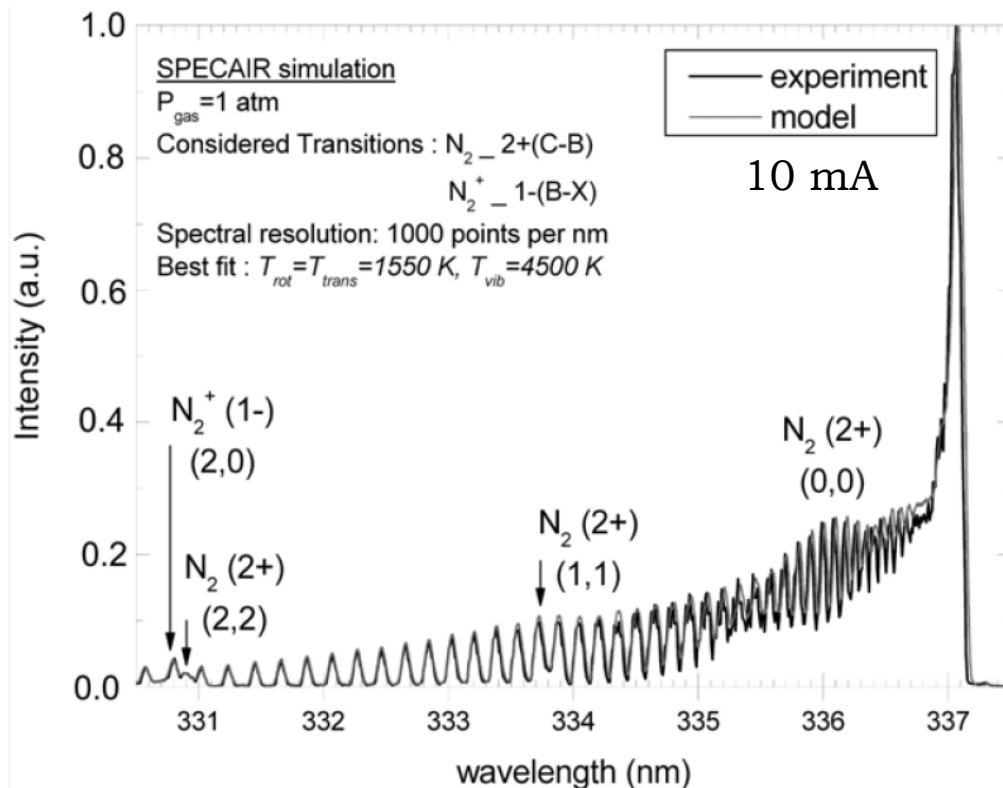


- Diagnostics on microplasmas are challenging because of their small size
- Optical diagnostic methods: emission and absorption spectroscopy, interferometry, Thomson scattering

T_g measurement by OES



- Measurement of the spectrum of the first/second positive band of N_2 and comparison to a simulated spectrum (small admixture of N_2 needed if the gas mixture does not contain N_2)
- Example: Atmospheric-pressure air μ discharge between two electrodes (2nd syst. positive)



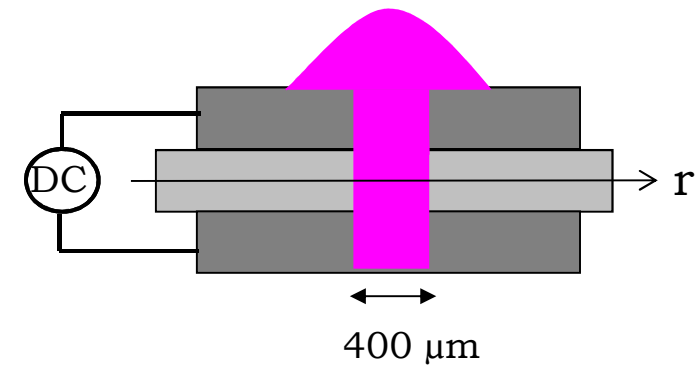
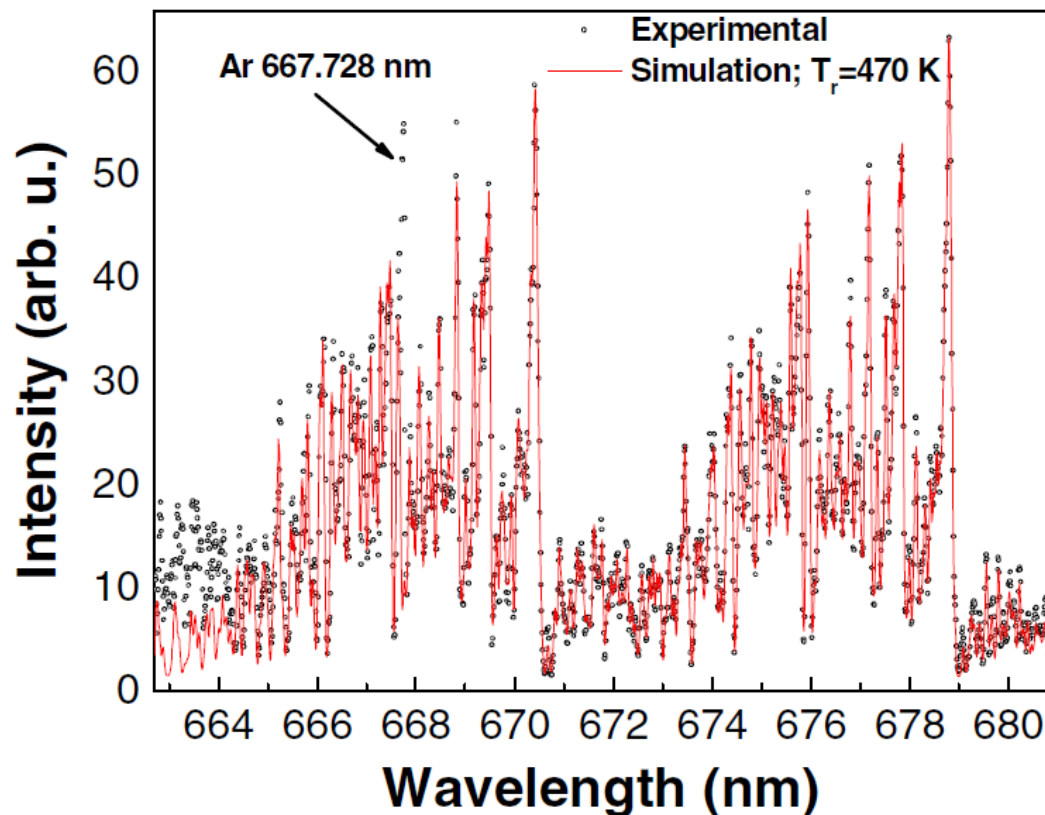
Staak et al. PSST **14** (2005) 700

For a current of 0.4 mA:
 $T_{\text{rot}} = 700 \text{ K}$ and $T_{\text{vib}} = 5000 \text{ K}$

T_g measurement by OES



- Temperature in microdischarges operating in noble gases and/or lower pressure is considerably lower
- Example: MHCD discharge in Ar at 1 mA and 100 Torr (1st system positive $B^3\Pi_g \rightarrow A^3\Sigma_u$)



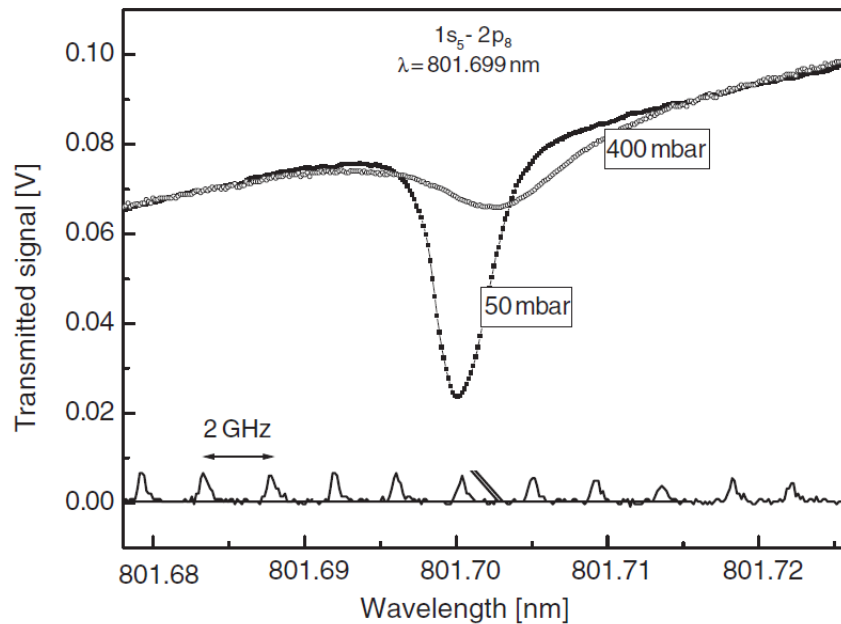
Lazzaroni et al. *JPhysD* **43** (2010) 124008

Radial dependence of T_g
constant
Increase of T_g with the
discharge current

T_g measurement by absorption spectroscopy

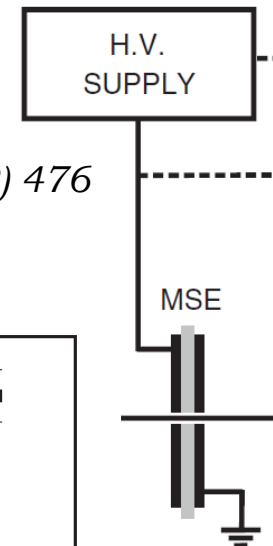
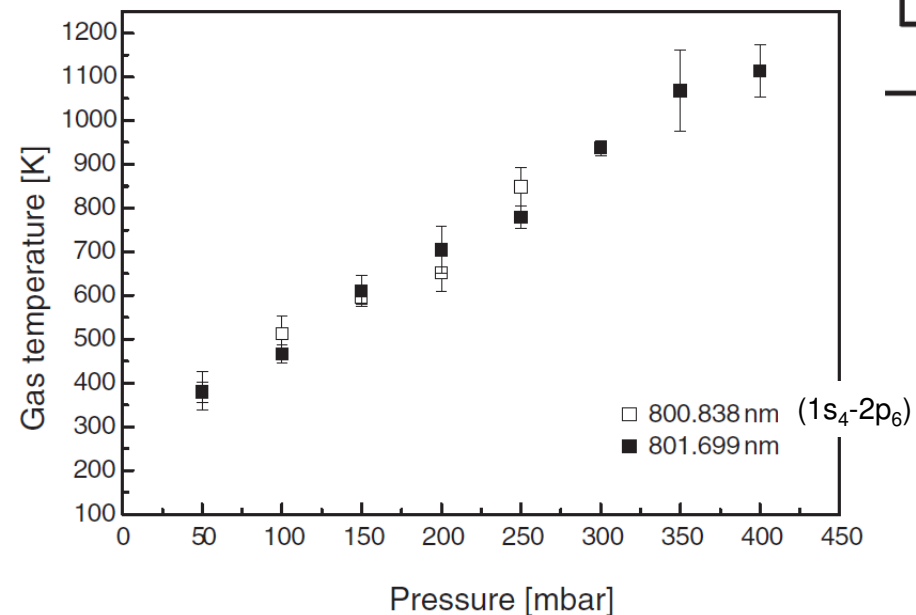


- High-pressure dc glow discharge based on micro-structured-electrode arrays in Ar at 0.5 mA (P = 50 to 400 mbar)
- Diode laser atomic absorption spectroscopy: Doppler broadening



$$\Delta\lambda_G = 7.16 \times 10^{-7} \lambda_0 \sqrt{\frac{T}{M}}$$

Penache et al. PSST 11 (2002) 476

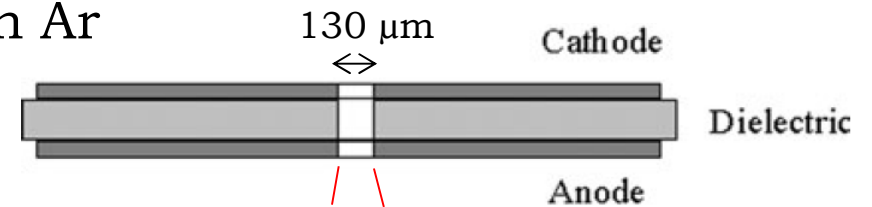


n_e measurement by OES

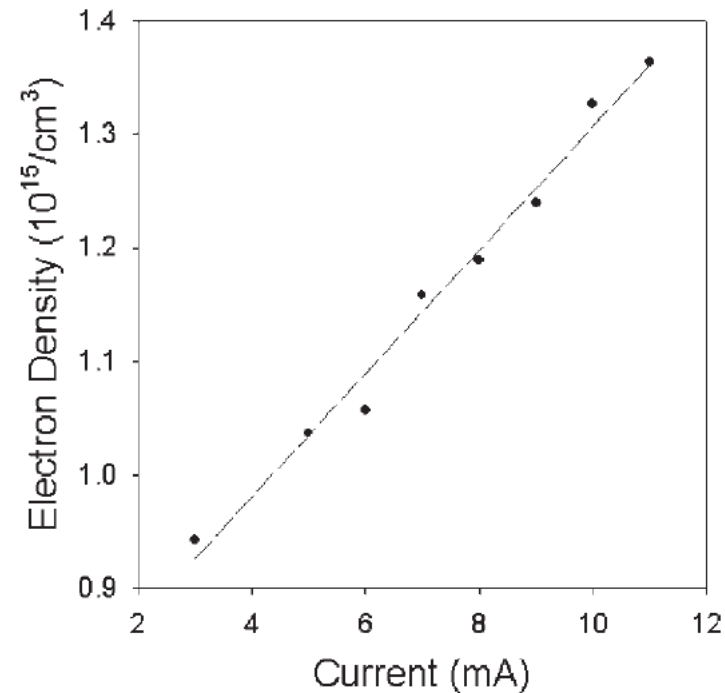
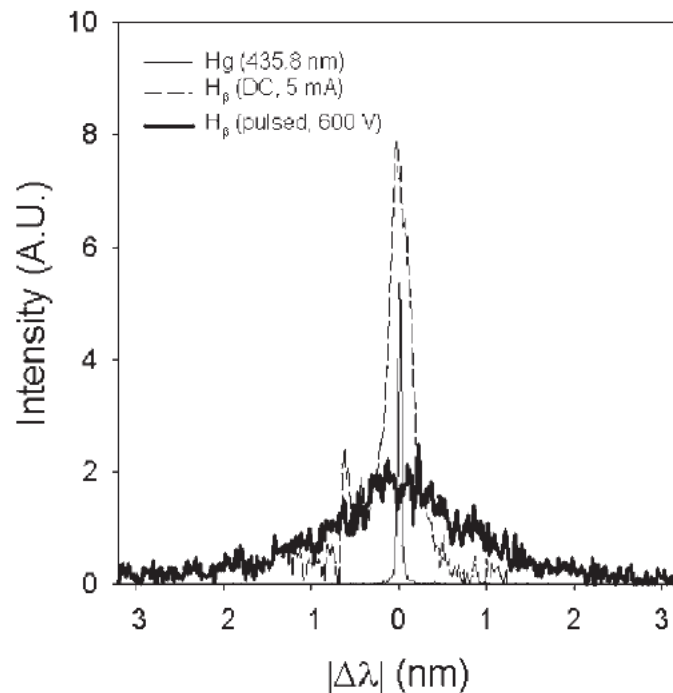


- DC MHCD at atmospheric pressure in Ar

Moselhy et al. JPhysD 36 (2003) 2922



- Small admixture of H₂ to use the Stark broadening of the hydrogen Balmer- β line at 486.1 nm

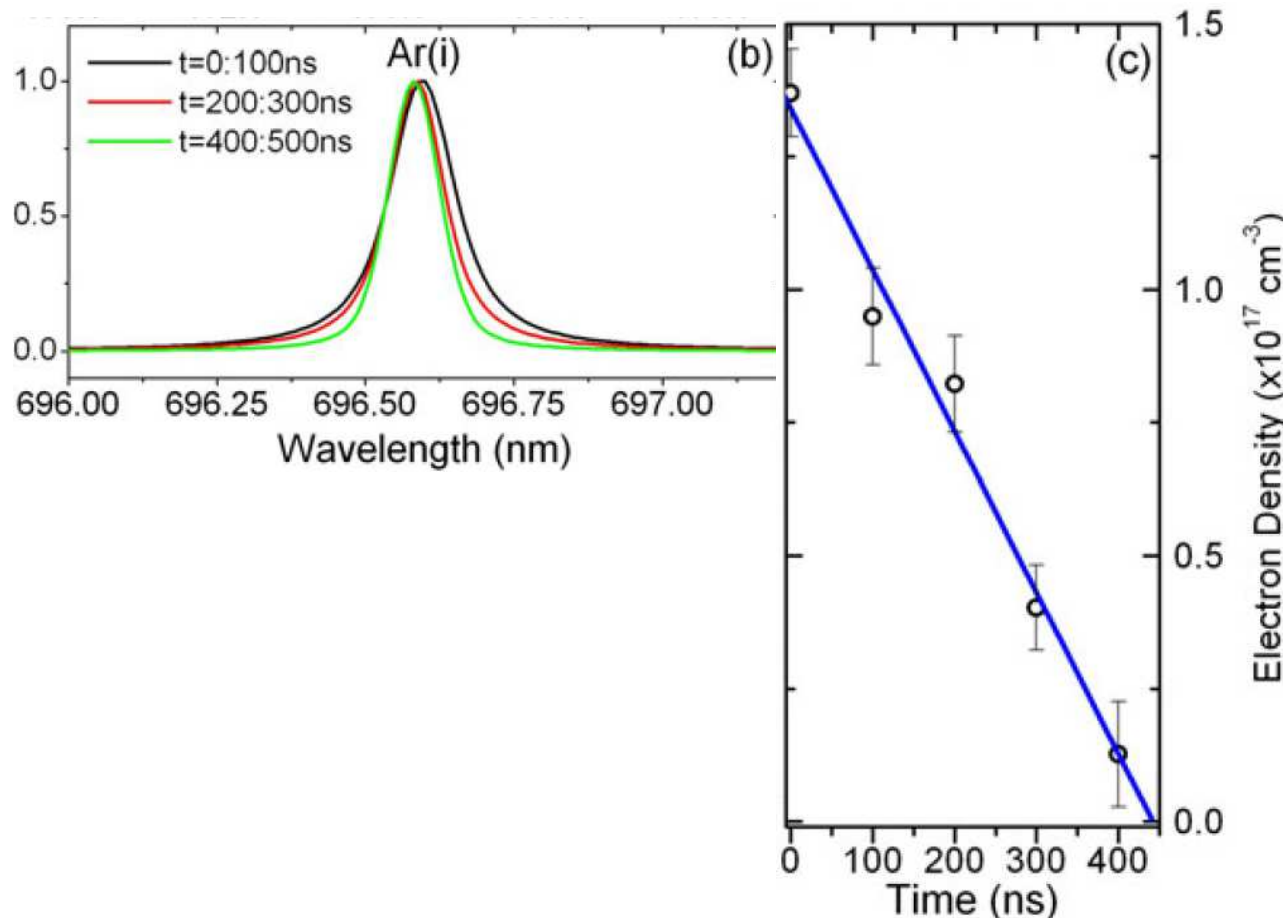


n_e in the pulse mode (600V, 10 ns, 10 mA)
= 5.10^{16} cm⁻³

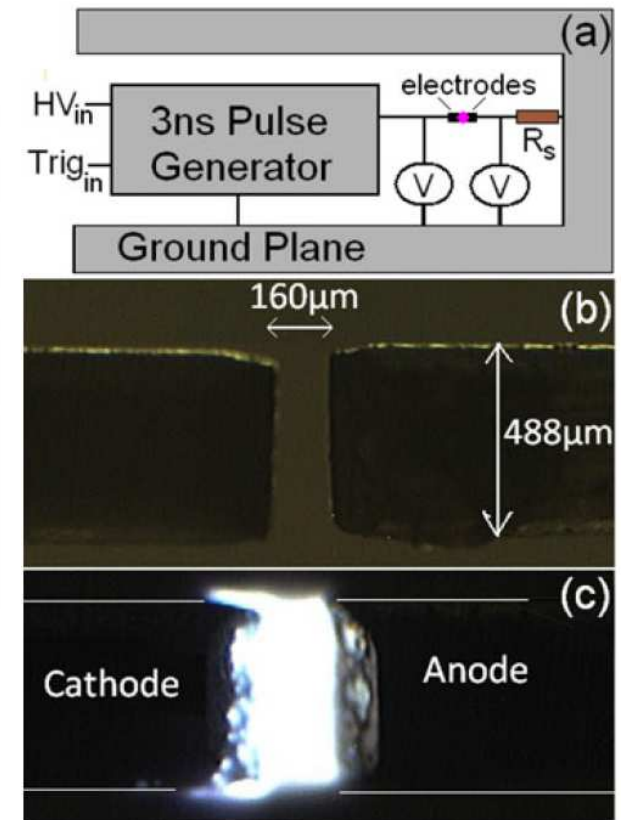
n_e measurement by OES



- 3 nanosecond (2.46 kV, 10 kHz) pulsed atmospheric pressure argon microdischarge. Time-resolved OES
- Stark broadening of the 4p-4s Ar line at 696.54 nm



Walsh et al. EPJD **60** (2010) 523

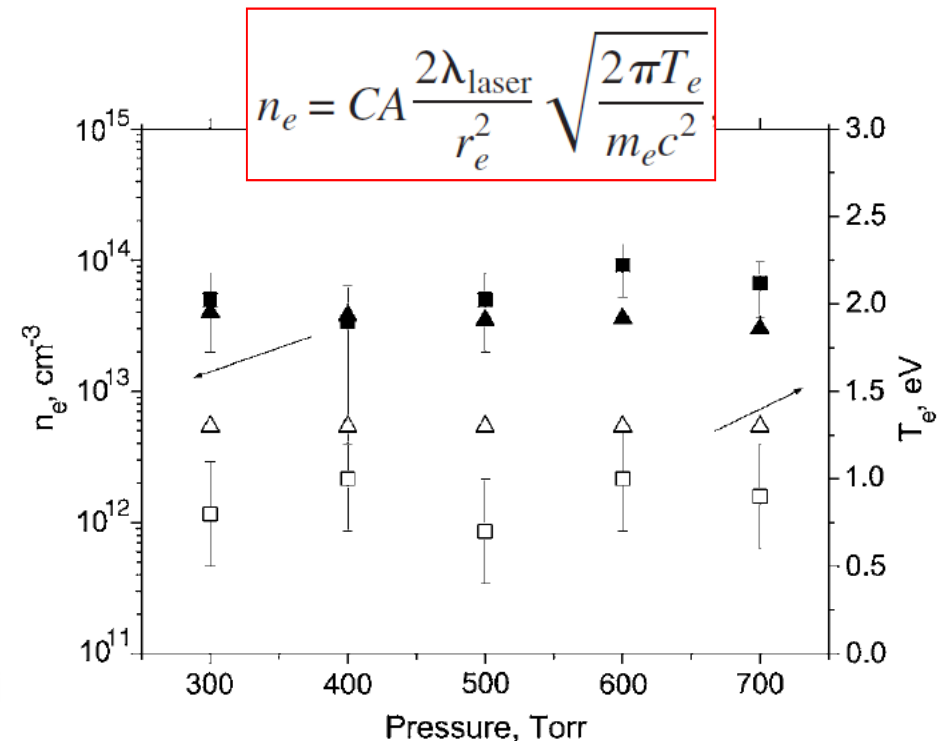
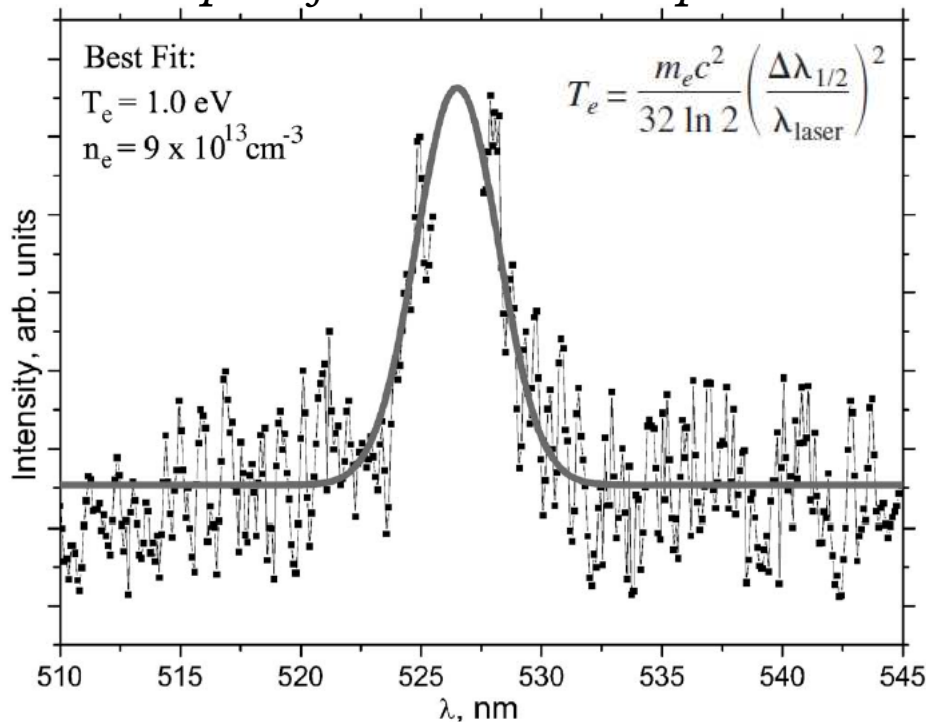


n_e measurement by Laser Thomson scattering

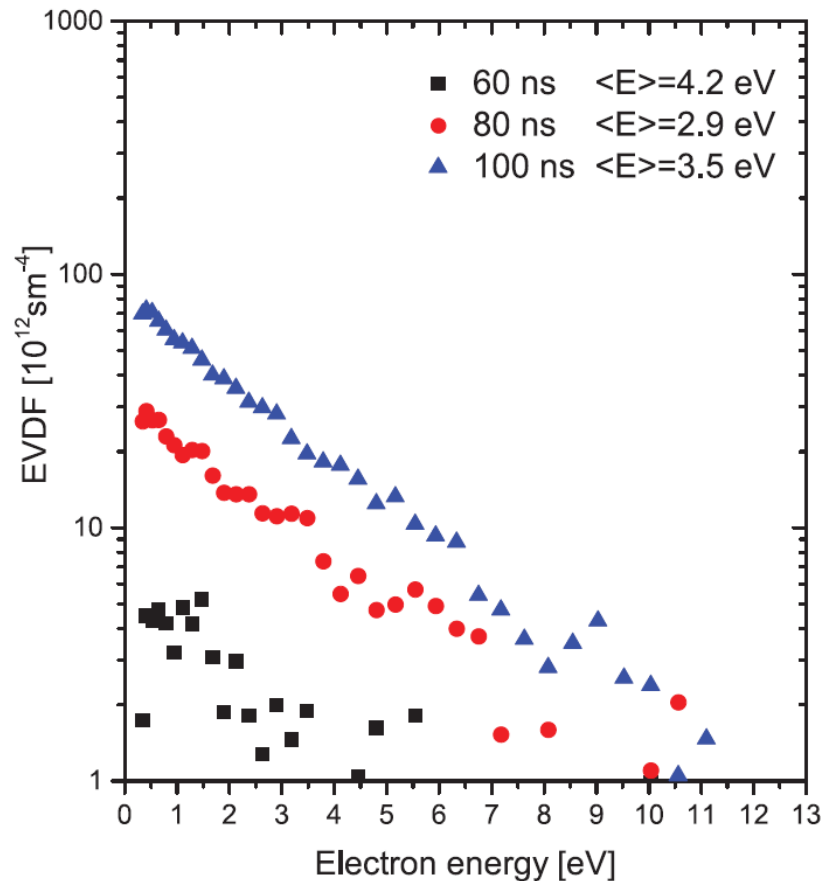


- Microdischarge between plane parallel electrodes (600 μm gap) in argon at intermediate pressure (300-700 Torr) for a current of 50 mA
- Use of the beam of a pulsed frequency-doubled Nd:YLF laser ($\lambda = 526.5 \text{ nm}$, $f = 3 \text{ kHz}$, $\tau = 100 \text{ ns}$, $P = 6 \text{ W}$)

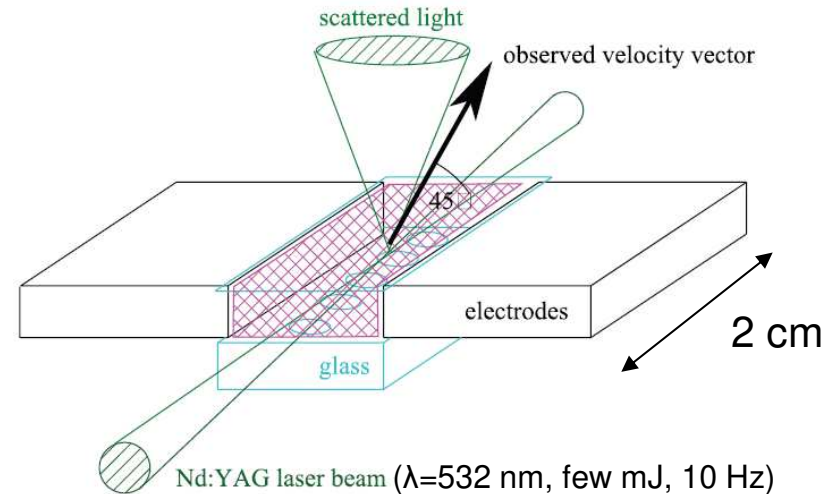
Example of net Thomson spectrum



Electron velocity distribution function



Schregel et al. *PSST* **25** (2016) 054003

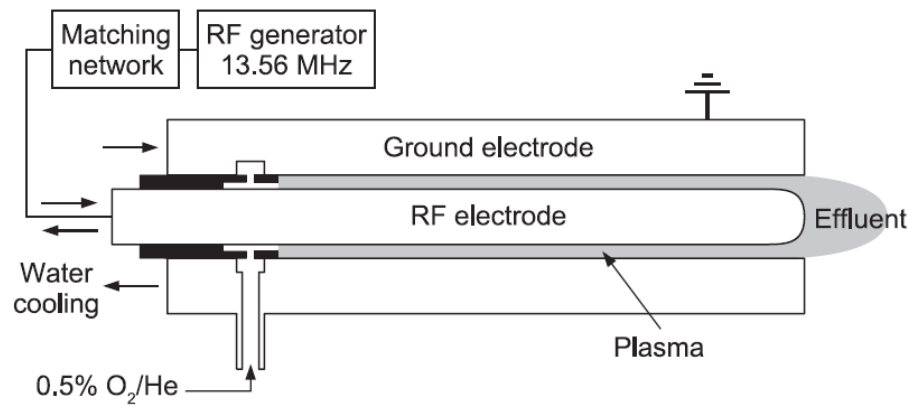


- 0.7 bar He microdischarge between planar electrodes (Mo) separated by a gap of 0.95 mm; 150 ns pulses with amplitude of 1 to 2 kV and $f = 5$ kHz
- Thomson scattering (elastic scattering on free electrons) \rightarrow eedf in a range of energies up to 12 eV
- Maxwellian eedf $\leftrightarrow T_e = 3.5$ eV at 100ns

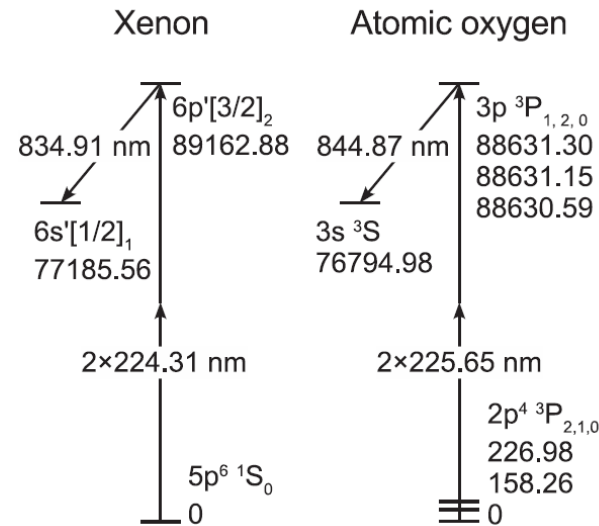
Reactive species density: n_O



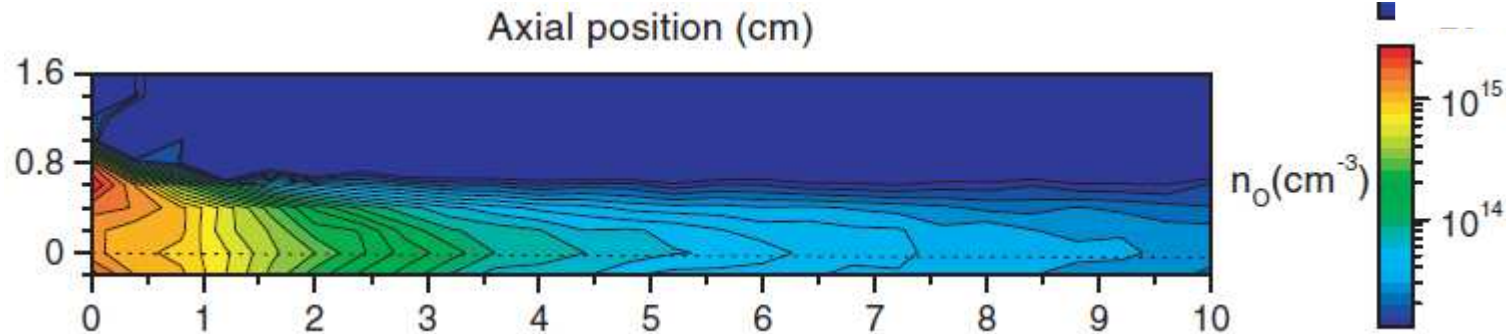
- Measurement of absolute atomic oxygen density by TALIF in a RF-APPJ in He/O₂



Niemi et al. PSST **14** (2005) 375



Two-photon
excitation
schemes

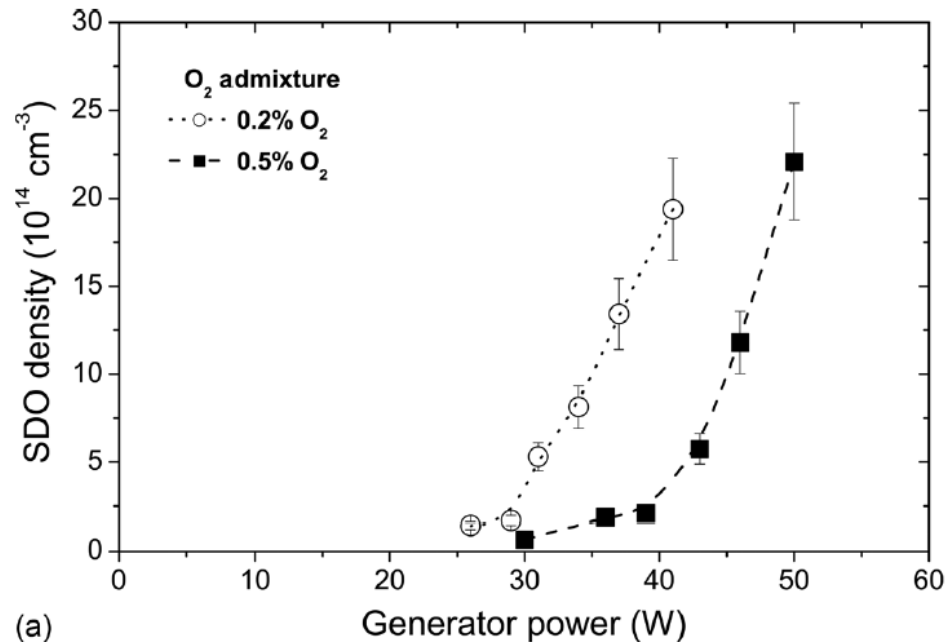
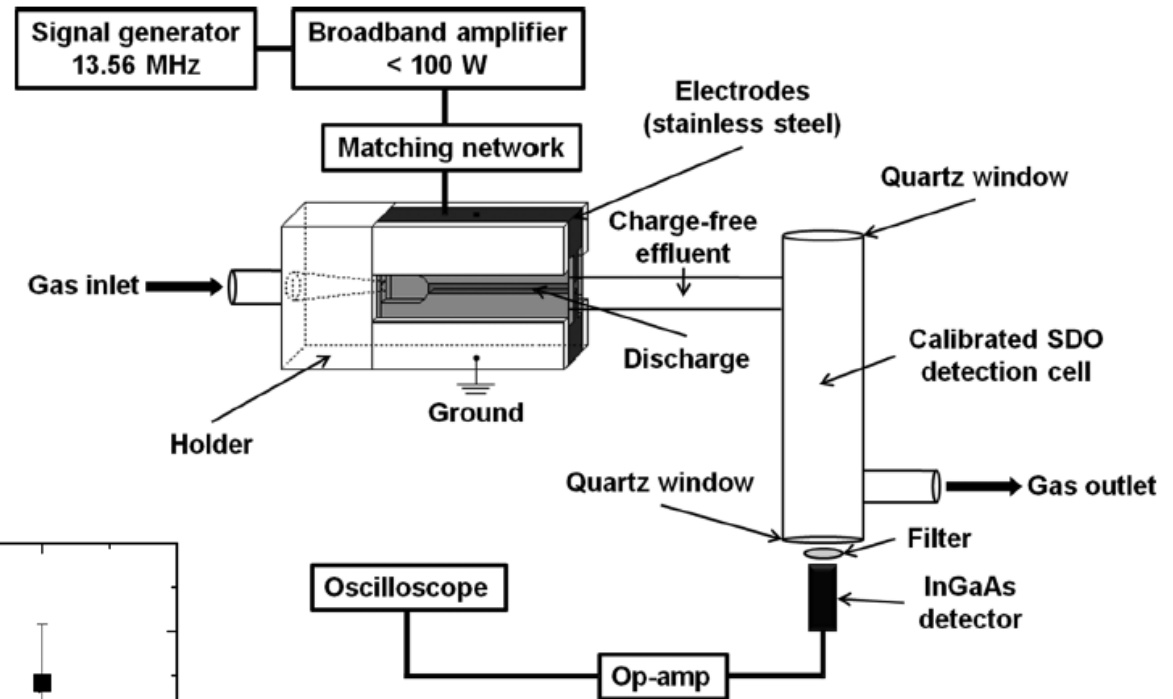


Reactive species density: $n_{O_2^*}$



Measurement of absolute O_2^* density by infrared optical emission spectroscopy in a RF-APPJ in He/ O_2 : emission at $1.27 \mu\text{m}$

Sousa et al. JAP 109 (2011) 123302

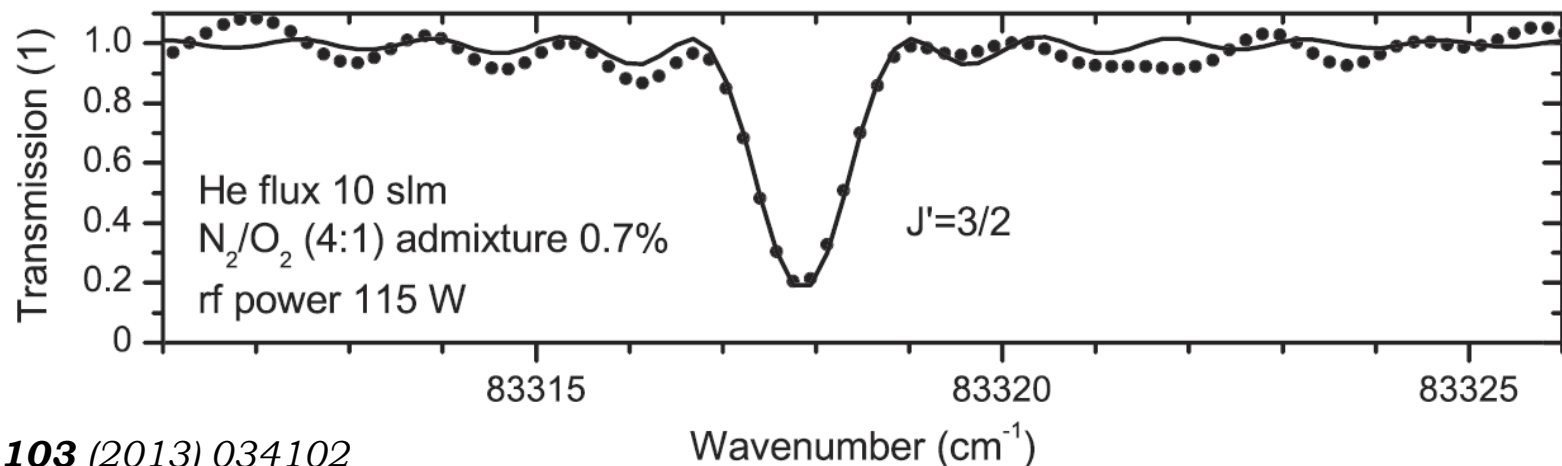
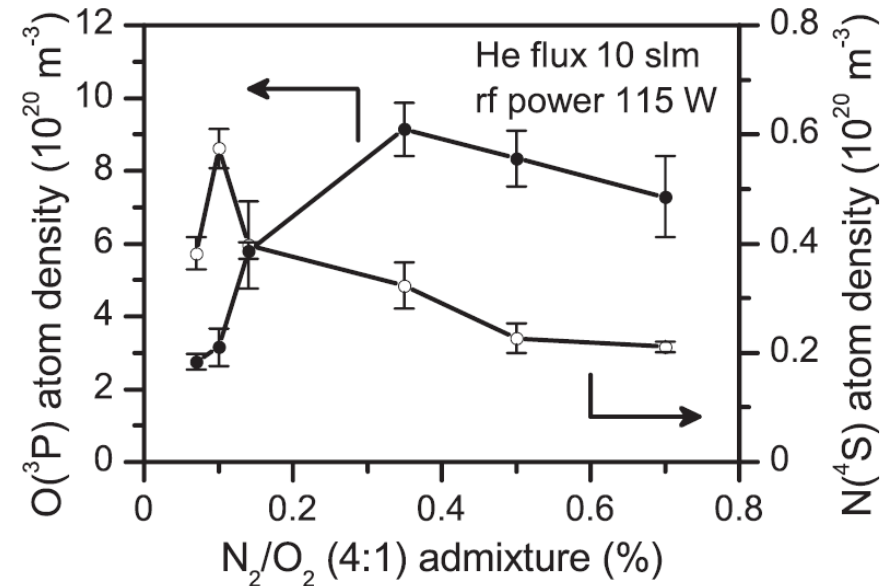


At larger O_2 admixture, shift towards higher power \rightarrow energy consumption through molecular vibrational and rotational excitations, dissociations

Reactive species density: n_N



- Measurement of absolute atomic nitrogen density by VUV FTS
- Atmospheric pressure RF microplasmas in He/N₂/O₂ (d=1mm)
- Transitions from the ground state to the quadruplet state located around 120 nm.



Modeling of microplasmas

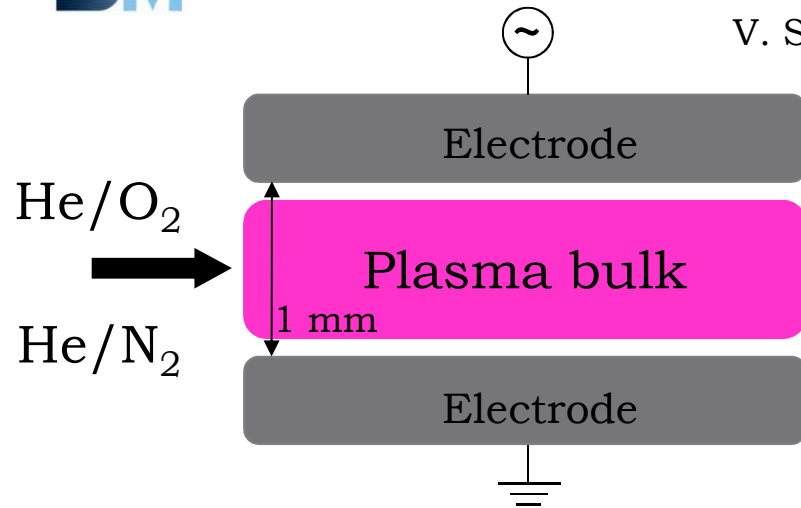


- Difficulty of diagnostics → considerable emphasis on numerical modeling
- Different kind of model: PIC, fluid and global

RF capacitive discharges at atmospheric pressure (APPJ)



V. Schulz-von der Gathen *et al.*, J. Phys. D **41** (2008) 194004

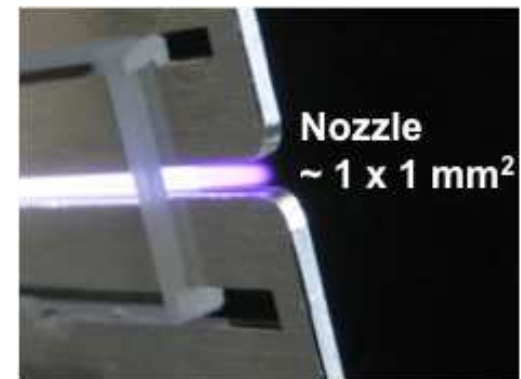


- RF frequency : 13.56-27.12 MHz
- electrode gap : 1 mm
- He/O₂ or He/N₂
- Gas flow to produce the plasma jet

- 1D PIC simulation in He/N₂
(Kawamura *et al.* PSST **23** (2014) 035014)

- 1D Fluid model in He/O₂
(Niemi *et al.* PSST **20** (2011) 055005)

- Global model in He/N₂ and He/O₂
(Lazzaroni *et al.* PSST **21** (2012) 035013)



Plasma kinetic description



- Particle-in Cell (PIC) simulations: Define macro-particle and solve the motion of each of these self-consistently with the fields. Method based on Newton's laws. Calculation of the electric field at every time-step from positions of charged particles (Poisson equation)
- Collisions are treated with Monte Carlo (MC) approach
- Advantage: accurate and self-consistent approach; fields, particle densities and fluxes obtained without making any assumptions about the particle temperatures or velocity distribution
- Disadvantage: long calculation time, no or simple chemistry

PIC-MCC simulation of APPJ



- RF-APPJ in He/0.1%N₂ (electropositive plasma); 1D simulation
- Each computer particle represents a cluster of 10⁷ “real” particles

For each timestep:

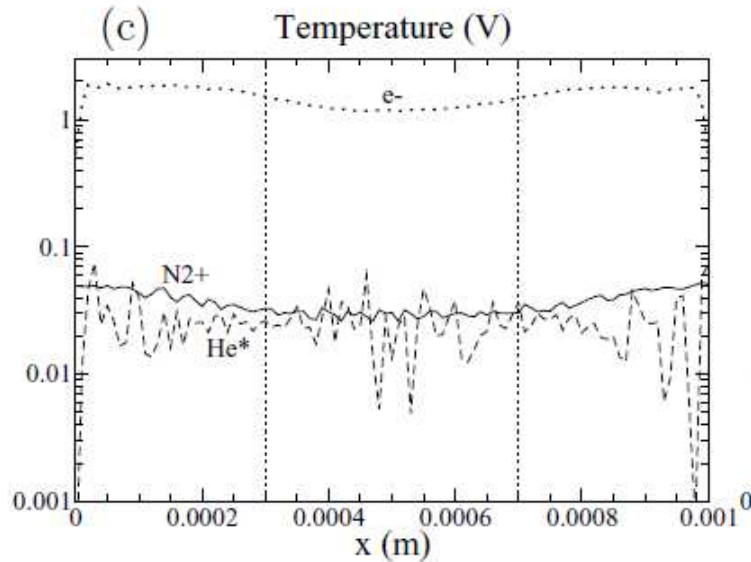
- 1) Particles are linearly weight \rightarrow charge density at the grid points
- 2) Poisson's equation \rightarrow electric field at the grid points
- 3) Determination of the force on each particle
- 4) Motion Newton equation \rightarrow particle new positions and velocities
- 5) Boundaries conditions \rightarrow bound particles are removed and injected particles are introduced (secondary electrons)
- 6) MCC handler \rightarrow collisions \rightarrow particle velocities are adjusted

3 species: N₂⁺, He* and e⁻
8 collisions

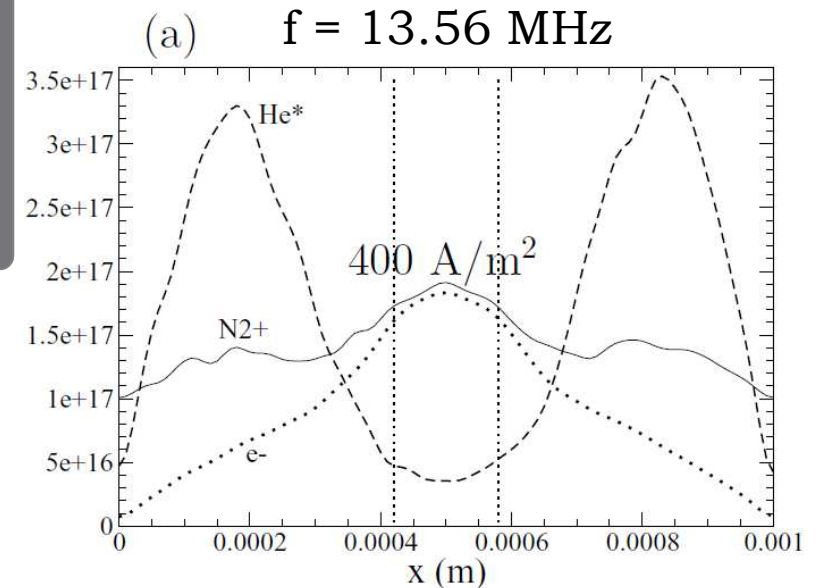
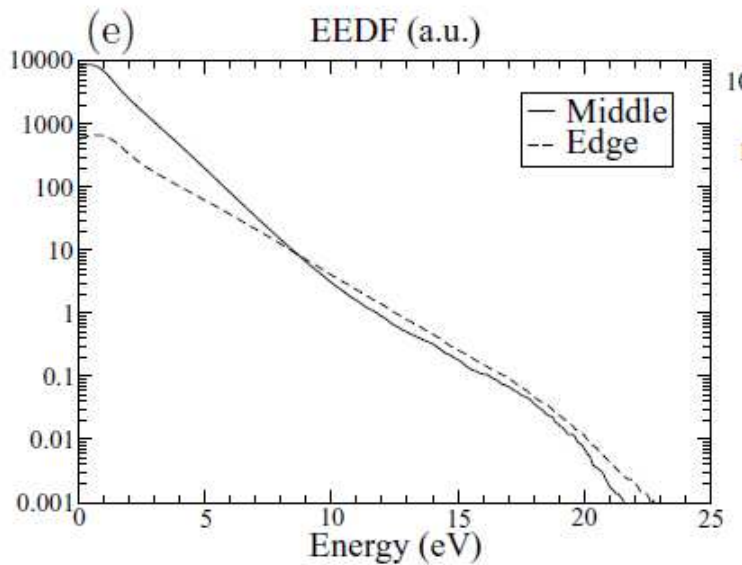
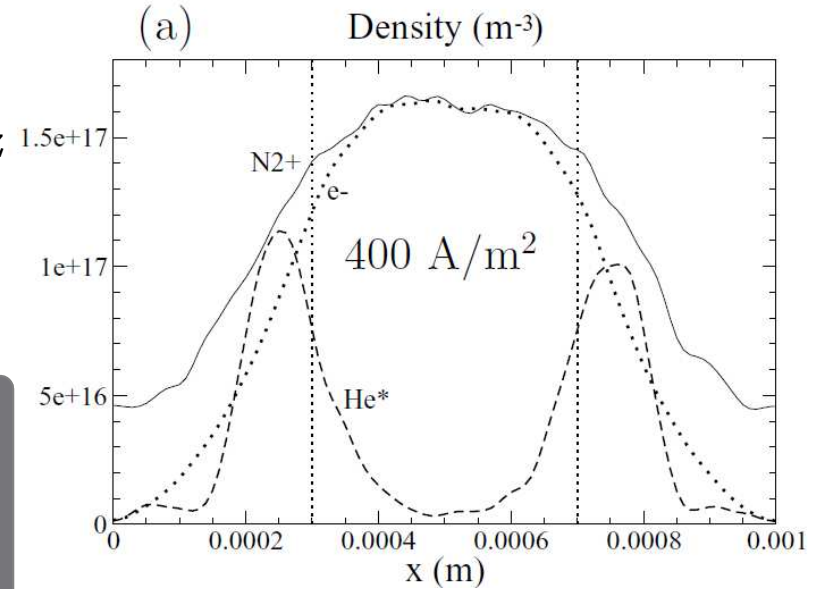
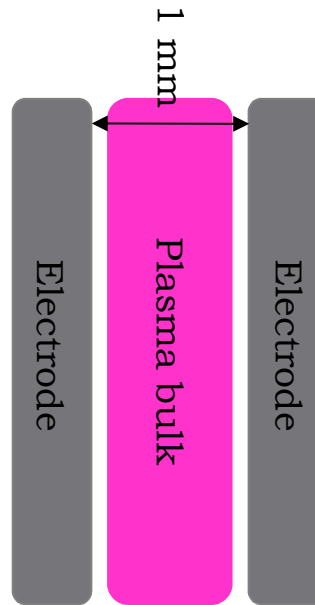
Time calculation ~ day

1. e + He \rightarrow e + He, Elastic Scattering
2. e + N₂ \rightarrow e + e + N₂⁺, Ionization
3. e + N₂⁺ \rightarrow N₂, Recombination
4. N₂⁺ + He \rightarrow N₂⁺ + He, Ion Elastic Scattering
5. e + He \rightarrow e + He*, Metastable Excitation
6. He* + 2He \rightarrow He₂* + He, Loss of He* (He₂* is not followed)
7. He* + N₂ \rightarrow e + N₂⁺ + He, Penning Ionization by He*
8. He* + He \rightarrow He* + He, He* Elastic Scattering

PIC simulation of APPJ



$f = 27.12$ MHz



Fluid description



- Particle conservation equation or continuity equation (obtained by integrating the Boltzmann equation over velocity space):

$$\frac{\partial n}{\partial t} + \nabla \cdot (n\mathbf{u}) = S - L$$

← volume source term
← volume loss term

- Momentum conservation equation (obtained by integrating the Boltzmann equation over velocity space after multiplication by $m\mathbf{v}$), with $\mathbf{B}=0$:

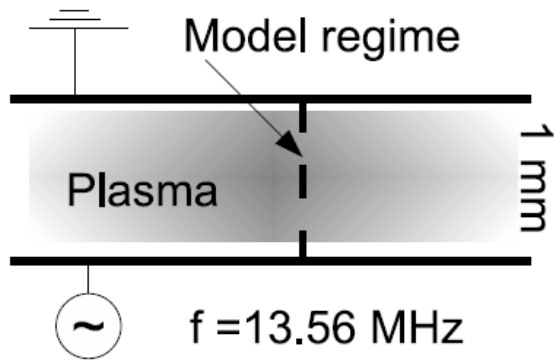
$$nm \left[\frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right] = nq\mathbf{E} - \nabla p - m\mathbf{u} [n\nu_m + S - L]$$

- Energy conservation (obtained by integrating the Boltzmann equation over velocity space after multiplication by $0.5m\mathbf{v}^2$)
- Coupled to Poisson equation for self-consistent E field
- Can be coupled to a Boltzmann's solver for accurate rate coefficient

Fluid model of APPJ



He/O₂
(electronegative plasma)

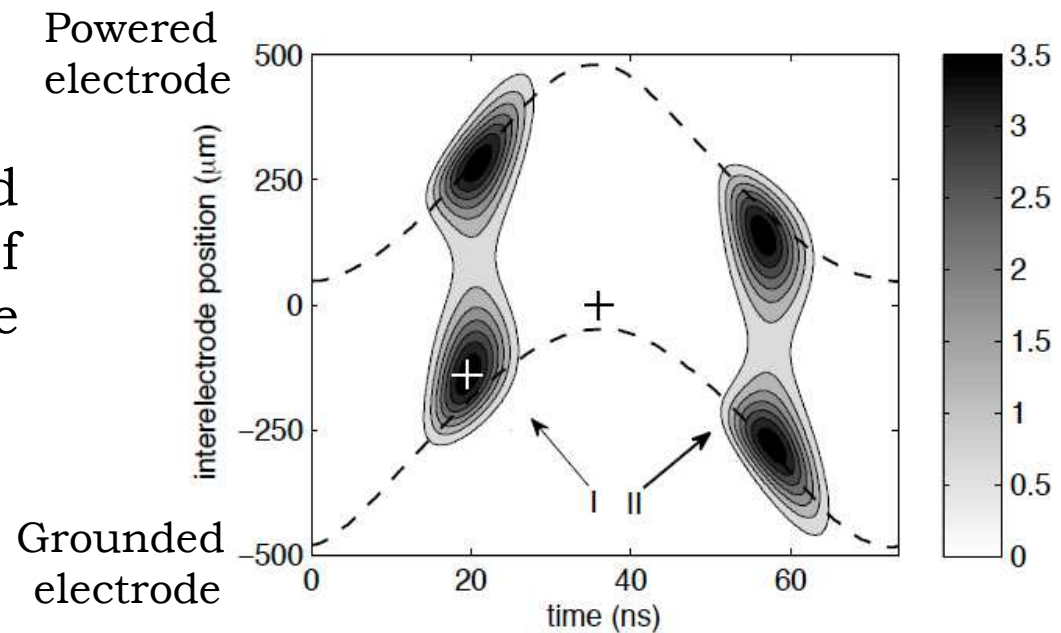


Spatio-temporal simulated electron-impact excitation of helium into the metastable state He

16 species
116 reactions

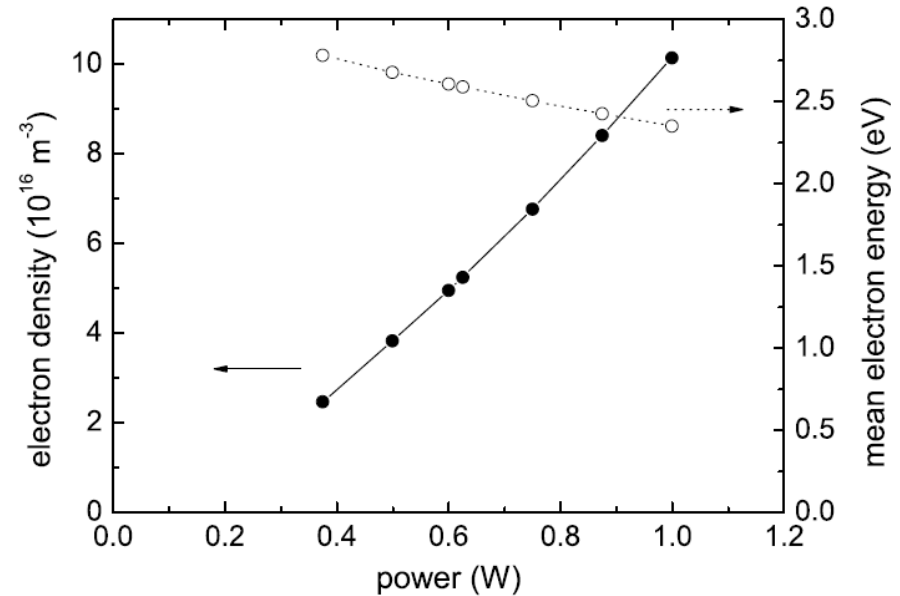
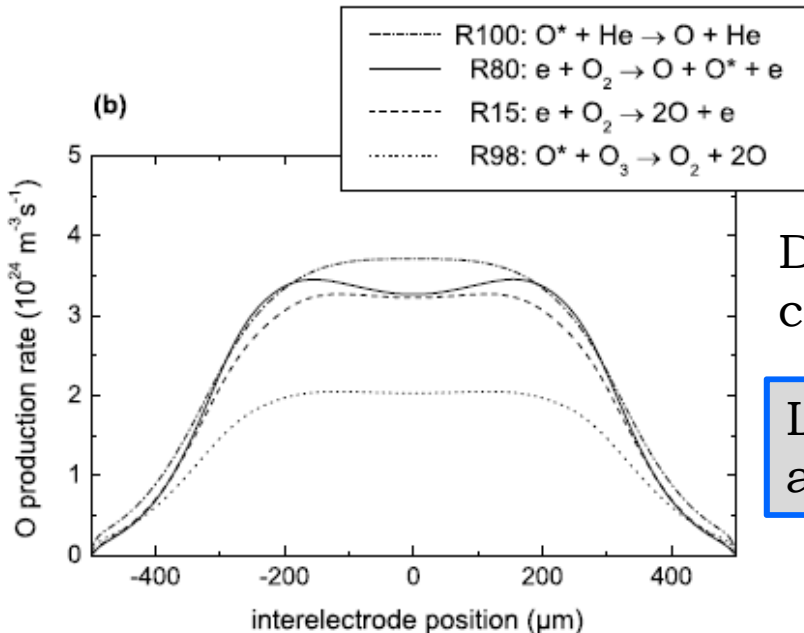
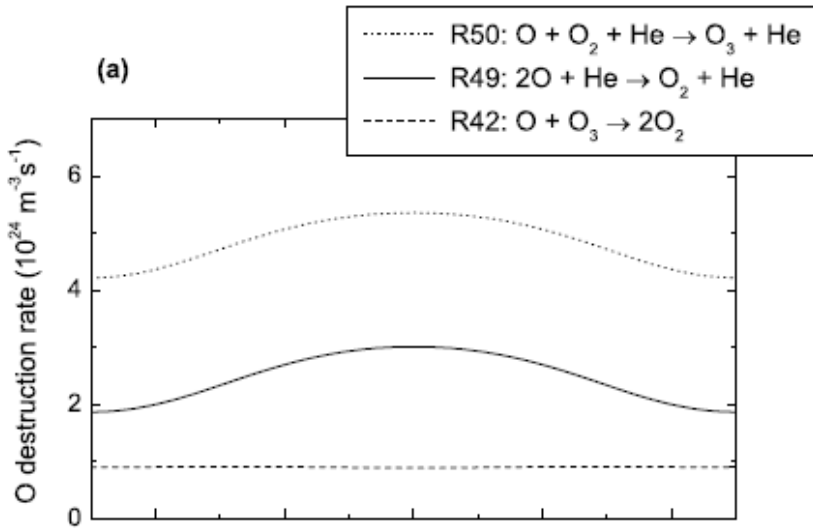
Time calculation ~ few hours

- Importance of electron dynamics
- Electron power strongly non uniform in space and time
- 1D fluid model



low power operation at 0.6 W/cm²

Fluid model of APPJ



Time and space averaged electron density and mean electron energy

Decrease of n_e and increase of T_e with power;
change in the time and space averaged EEDF

Lost of accuracy but faster calculation time
and more complexe chemistry than PIC model

Global description



- Volume-averaged (0D) model: densities and temperatures are uniform in space (obtained by spatial integration of fluid equations)

- Particle balance:
$$\frac{dn_{\alpha}}{dt} = G_{\alpha} - P_{\alpha}$$

- Electron power balance:

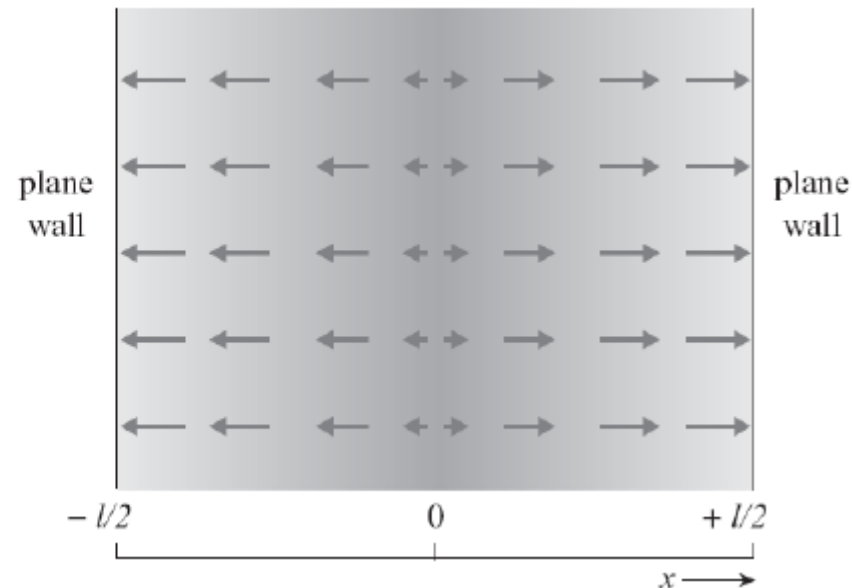
$$\frac{d}{dt} \left(\frac{3}{2} n_e e T_e \right) = P_{\text{abs}} - P_{\text{dis}}$$

Particle balance at low pressure



- Particle losses at the walls. The term P_α contains volume and wall losses. For the electron particle balance at low pressure, wall losses dominate and in one dimension:

$$\frac{d\bar{n}_e}{dt} = \bar{n}_e n_g K_{iz} - \frac{2\Gamma_{\text{wall}}}{l}$$



- A plasma transport theory is required to relate the space-averaged electron density to the flux at the wall

Particle balance: issues in μ plasmas



- In microdischarges, ionization is often non uniform. Classical low-pressure transport theory do not apply
- Fortunately in some instances volume losses dominate (recombination)
- However, evaluation of wall losses is a critical point for high-pressure discharges modeling
- Properly evaluate the reaction rates ($T_e(x,t)$)
- Properly evaluate the electron power absorption

Hybrid analytic-numerical global model of APPJ

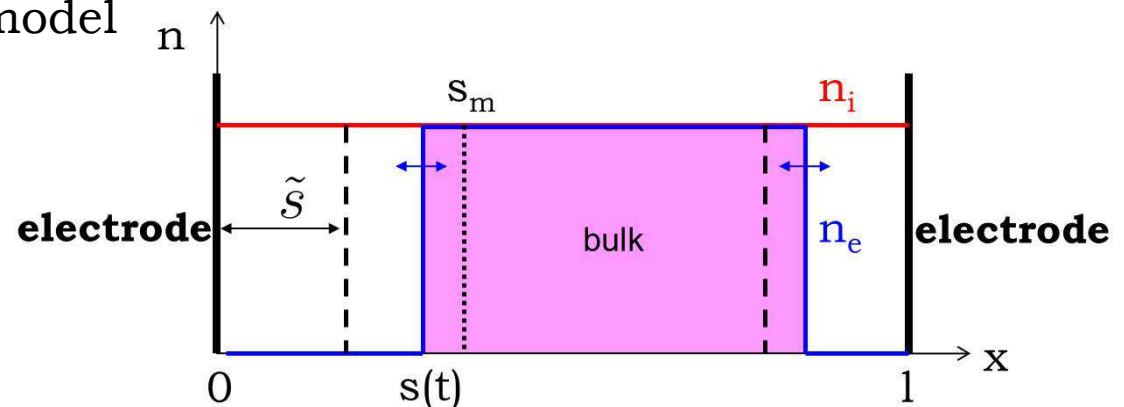


- Fast solution of the discharge equilibrium: exploration of a large parameter space
- Variations of discharge parameters with discharge composition and RF power

Homogeneous discharge model

Principles of Plasma Discharges and Materials Processing, M.A.Lieberman and A.J.Lichtenberg, 2nd Edition, John Wiley and Sons Inc., New York, 2005

Physics of Radiofrequency Plasmas, P. Chabert and N. Braithwaite, Cambridge University Press, 2011



- Three regions: - a quasi-neutral plasma ($n_i = n_e = n_0 = \text{const}$)
- two sheaths ($n_i = n_0$ and $n_e = 0$)
- Ions do not respond to the RF field
- Analytical expressions of $J_C(t)$, $s(t)$, $E(t)$ and $P_C(t)$

Global model



Particle balance

$$\frac{dn_{\alpha}}{dt} = G_{\alpha} - P_{\alpha}$$

numerically

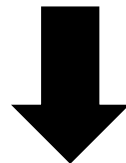
/ Global e⁻ power balance

$$\frac{dT_e}{dt} = \frac{2}{3} \frac{P_c(t)}{en_e} - \frac{2}{3} \nu \frac{3m}{M} T_e - \sum_j \frac{2}{3} \nu_j \mathcal{E}_j$$

analytically

3) Mean rate coefficient

- 1) $P_c(t) \leftarrow$ homogeneous discharge
- 2) $T_e(t) \leftarrow$ analytical integration



Discharge equilibrium

Electron temperature



Electron energy balance:
$$\frac{dT_e}{dt} = \frac{2}{3} \frac{P_c(t)}{en_e} - \frac{2}{3} \nu \frac{3m}{M} T_e$$

$$P_c(t) = \bar{P}_c [1 + \cos(2\omega t - \theta)]$$

Hypothesis: - $T_e(t)$ uniform in the bulk
- n_e uniform and independant of time

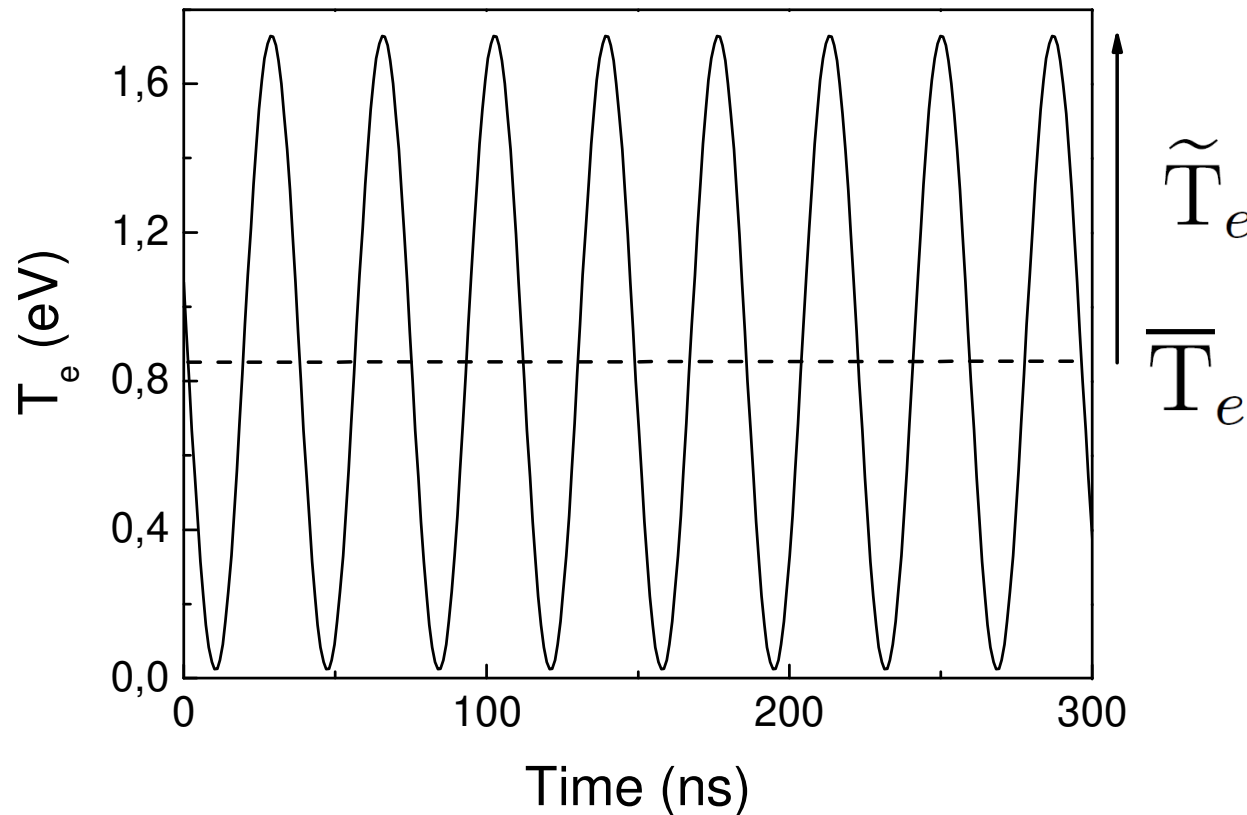
integration

$$T_e(t) = \bar{T}_e + \tilde{T}_e \cos(2\omega t - \phi_0)$$

Electron temperature



$$T_e(t) = \bar{T}_e + \tilde{T}_e \cos(2\omega t - \phi_0)$$



Unlike at low pressure, T_e oscillates during the rf cycle. This point is extremely important for the global modeling.

Effective rate coefficients



- Electron-activated processes strongly affected by $T_e(t)$
- Maxwellian rate coefficients: $K = K_0 \exp(-\mathcal{E}_a/T_e)$
- Averaging over the oscillating T_e
→ enhanced rate coefficient:

$$\bar{K} = K_0(T_{e\max}) \operatorname{erfc} \left(\sqrt{\mathcal{E}_a/2\tilde{T}_e} \right)$$

$$\langle K_0 \exp(-E_a/T_e(t)) \rangle \neq K_0 \exp(-E_a/\langle T_e(t) \rangle)$$

Particle balance: He/N₂ chemistry



Particle balance for each species:

$$\frac{dn_{\alpha}}{dt} = G_{\alpha} - P_{\alpha}$$

Electropositive plasma He-N₂ mixture

- 0.1% of N₂
- 8 species
- 15 reactions in the gas phase
- Surface reactions for all species
- T_g=345 K

The 8 species:

He, He⁺, He₂⁺, He*,
He₂^{*}, N₂, N₂⁺ and e⁻

Time calculation ~ few ms

Particle balance: He/O₂ chemistry



Particle balance for each species:

$$\frac{dn_{\alpha}}{dt} = G_{\alpha} - P_{\alpha}$$

Electronegative plasma He-O₂ mixture

- 0.1-1 % of O₂
- 16 species
- 132 reactions in the gas phase
- Surface reactions for all species
- T_g=345 K

The 16 species:

He, He⁺, He₂⁺, He*,
He₂^{*}, O₂, O, O₃, O⁺,
O₂⁺, O⁻, O₂⁻, O₃⁻, O*,
O₂^{*} and e⁻

Time calculation ~ few seconds

Sheath physics: limitation in %O₂



- Finite sheath widths limit operating regimes in atmospheric pressure discharges.
- Requirement for equilibrium: $\bar{s} < l/2$
- Time average energy balance and magnitude of the oscillating sheath width give ($\nu \gg \omega$):

$$\omega \bar{s} = u_{Bg} (6\bar{\zeta})^{1/2}$$

$u_{Bg} = (e\bar{T}_e/M)^{1/2}$

RF time average energy loss factor

Sheath physics: limitation in %O₂



- Sheath size increases with %O₂
- Limit of %O₂ when sheath size exceeds half-gap

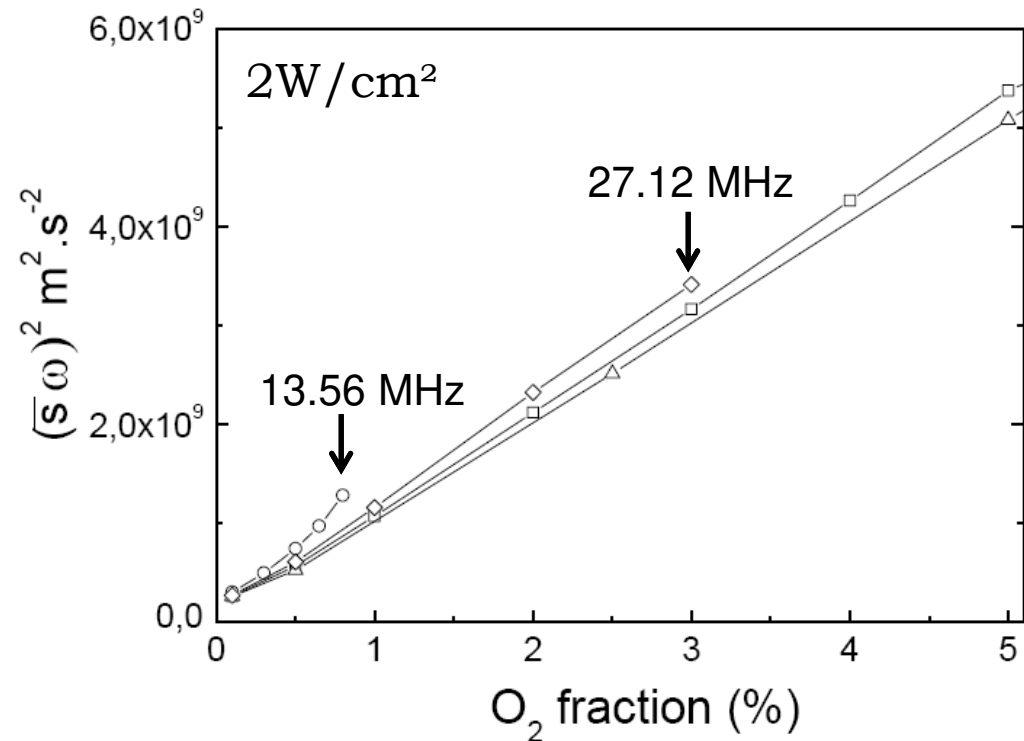
$$\omega \bar{s} = u_{Bg} (6\bar{\zeta})^{1/2}$$

- The RHS increases with the O₂ fraction (inelastic collisions)

- Requirement for equilibrium: $\bar{s} < l/2$

- Therefore there is a limit %O₂ at given frequency

- Higher possible %O₂ at higher frequency

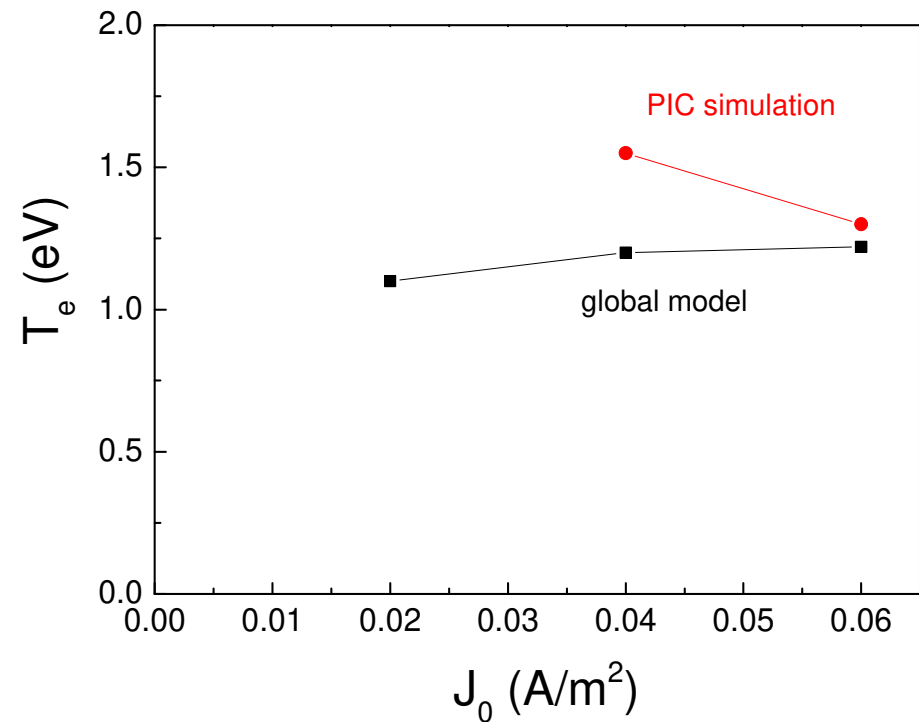
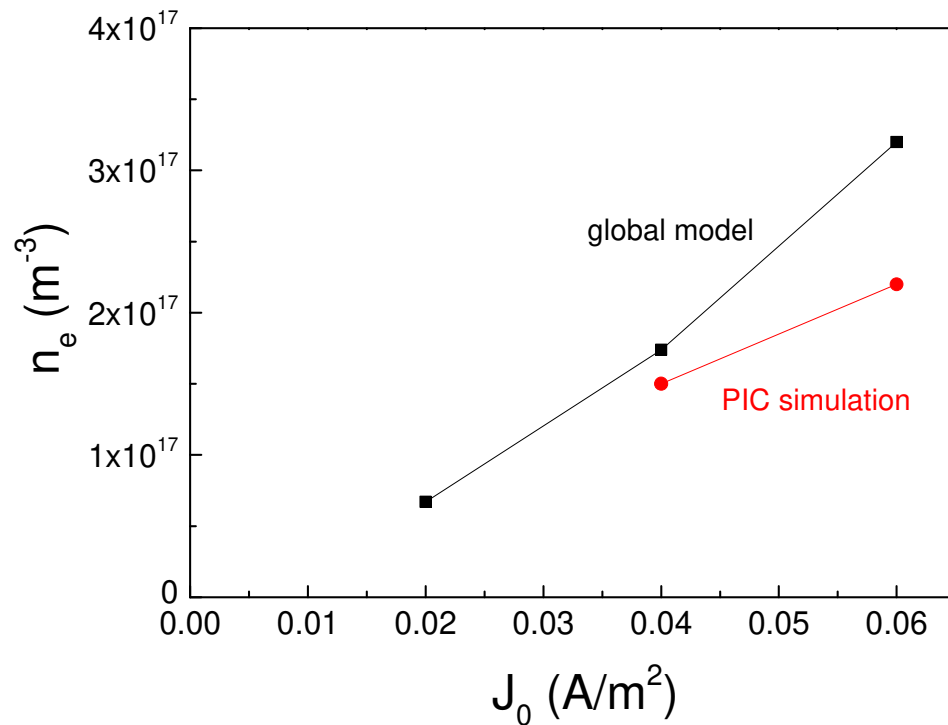


$$O_{2Lim} \% \approx 1.3 \left(\frac{l(\text{mm}) f(\text{MHz})}{13.56} \right)^2$$

Global model vs PIC simulation



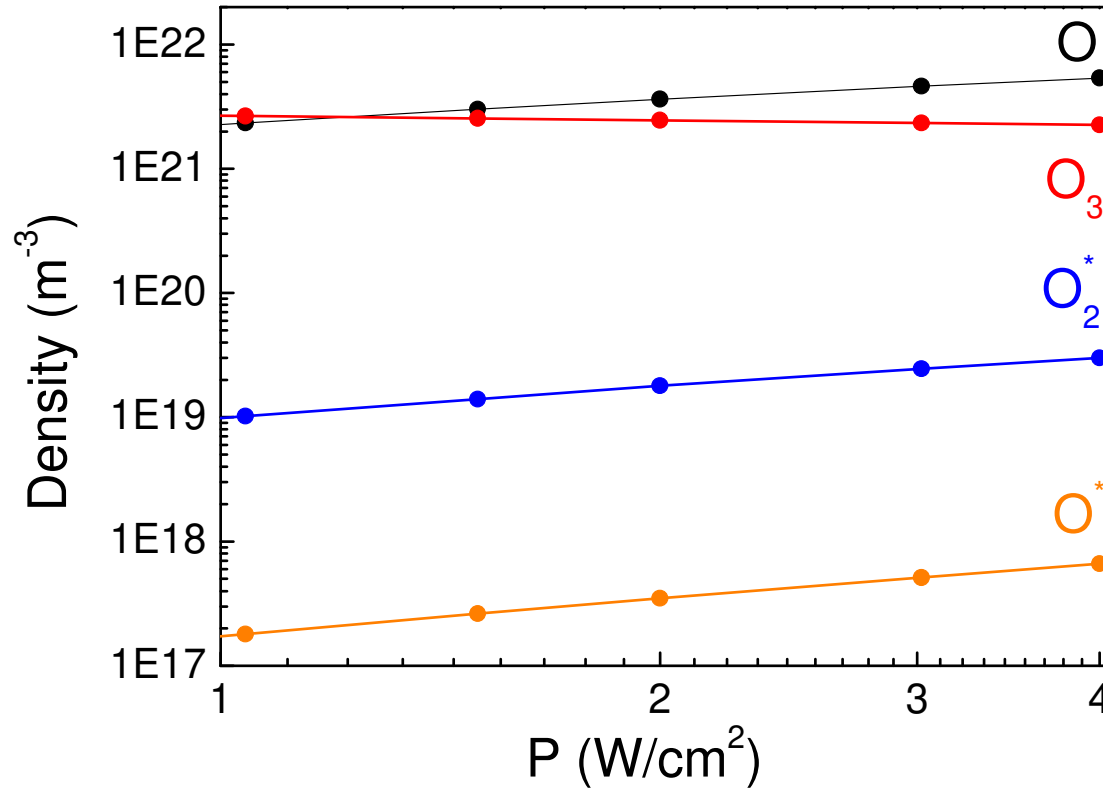
He/N₂



Global model vs fluid model: neutral species



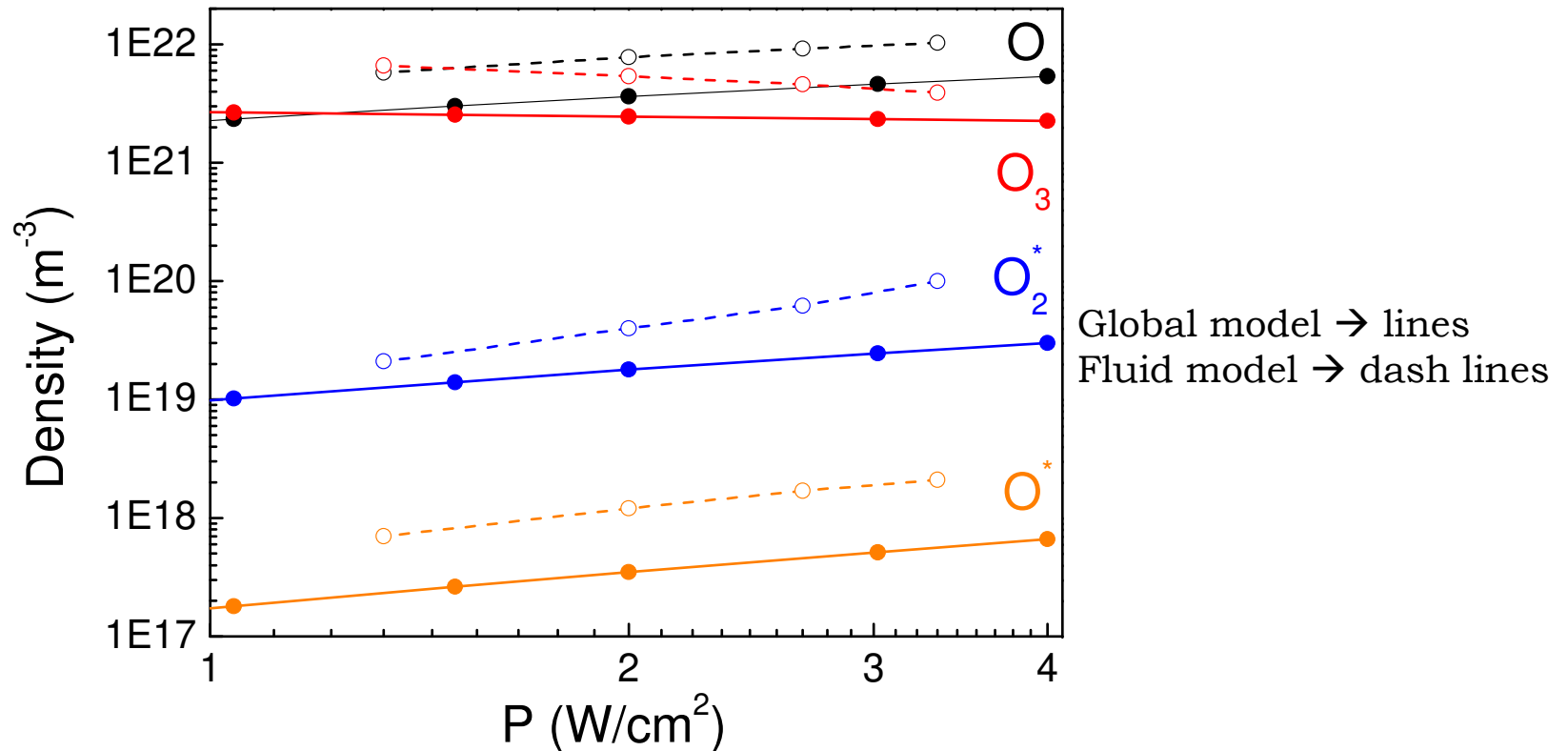
He/O₂



Global model vs fluid model: neutral species



He/O₂

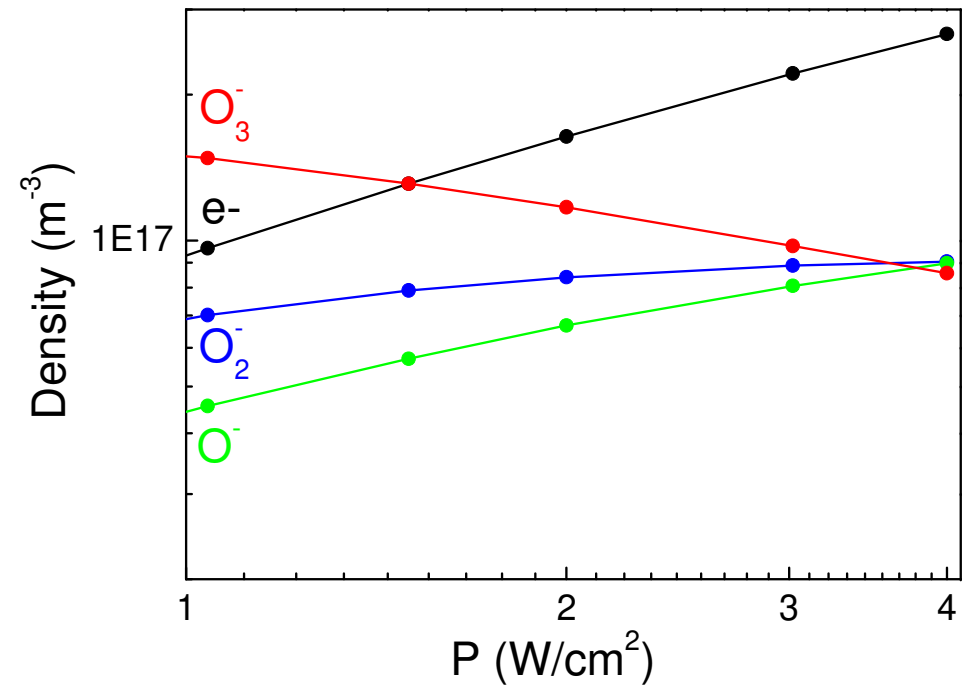
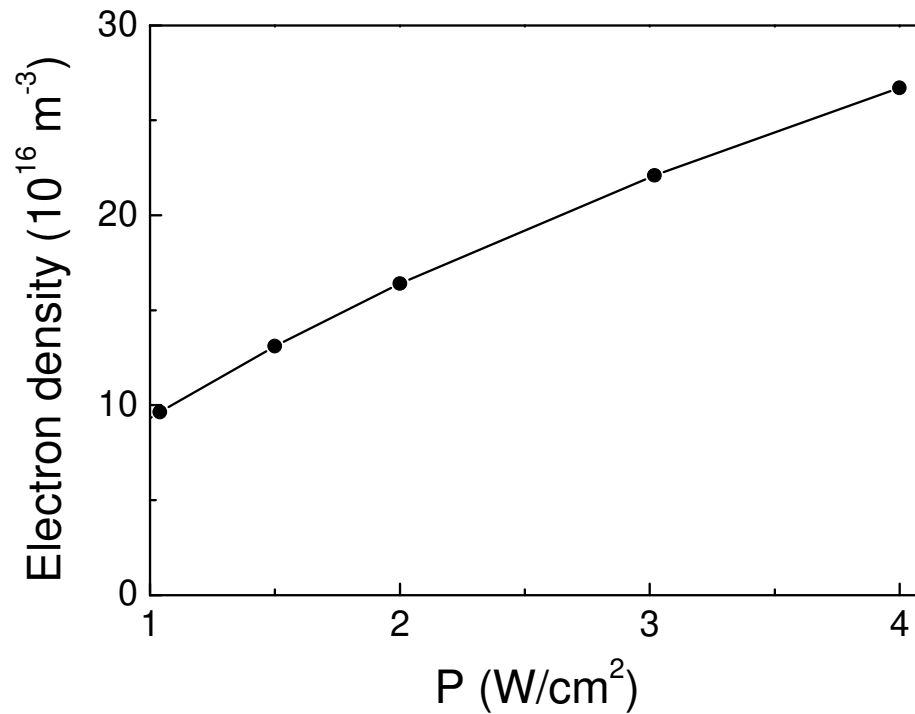


- Good qualitative agreement
- Global model predictions lower than that of fluid model (factor of 2 for O and O₃ and factor of 3 for O* and O₂*)

Global model vs fluid model: charged species



He/O₂

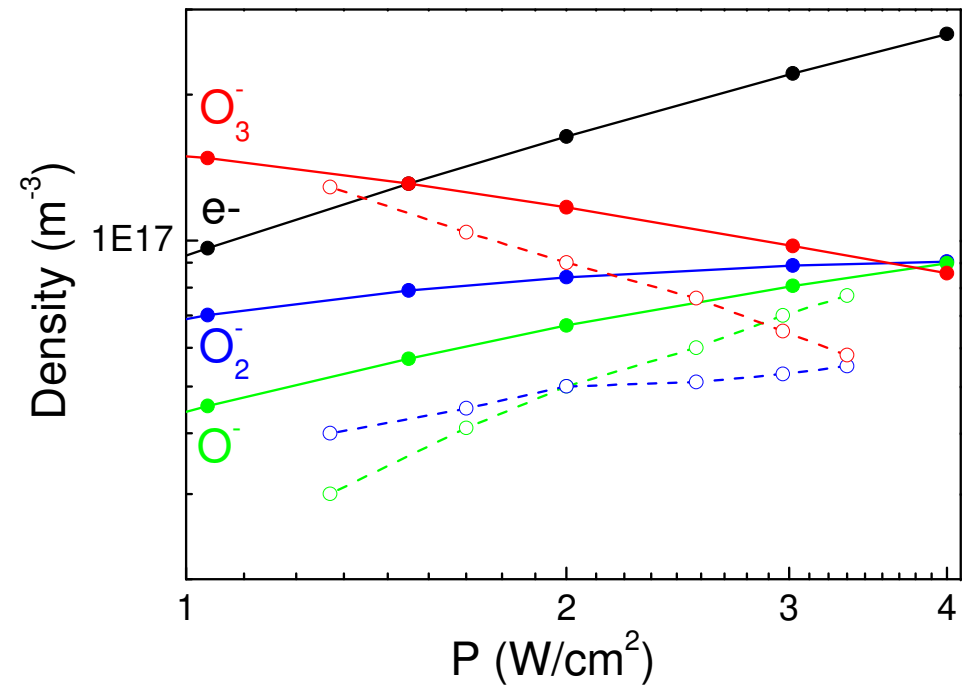
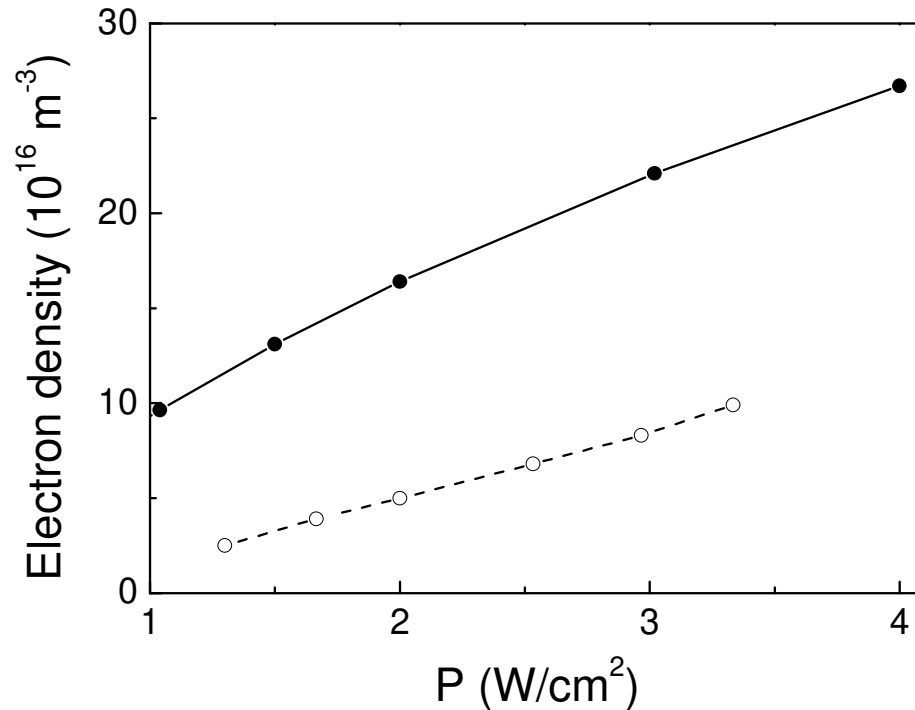


Global model vs fluid model: charged species



Global model → symbols + lines
Fluid model → open symbols + dash lines

He/O₂



- Good qualitative agreement
- n_e (global model) $\approx 3 \cdot n_e$ (fluid model)
- Negative ion densities (global model) $\approx 3/2 \cdot$ Negative ion densities (fluid model)

Hybrid model

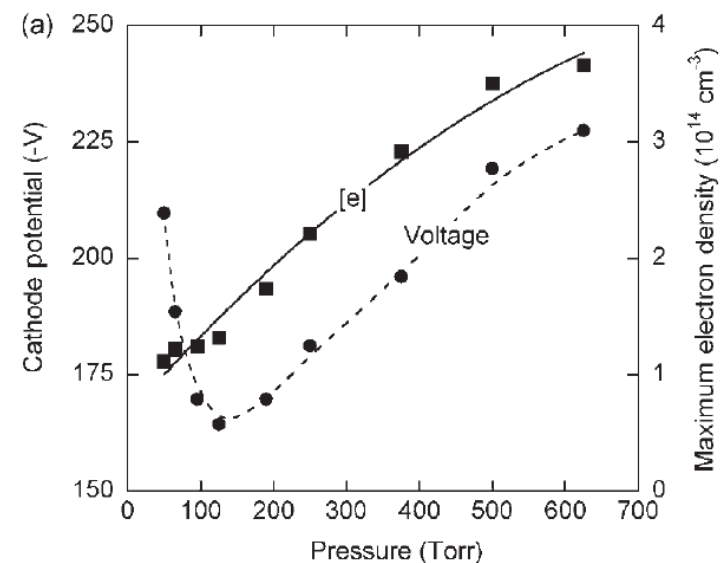
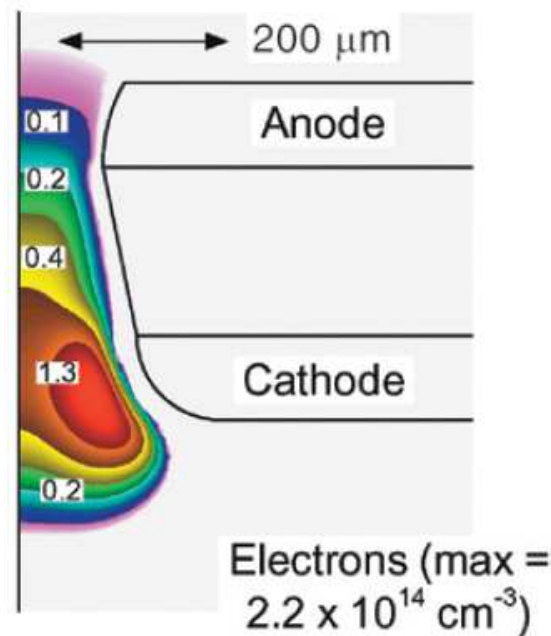
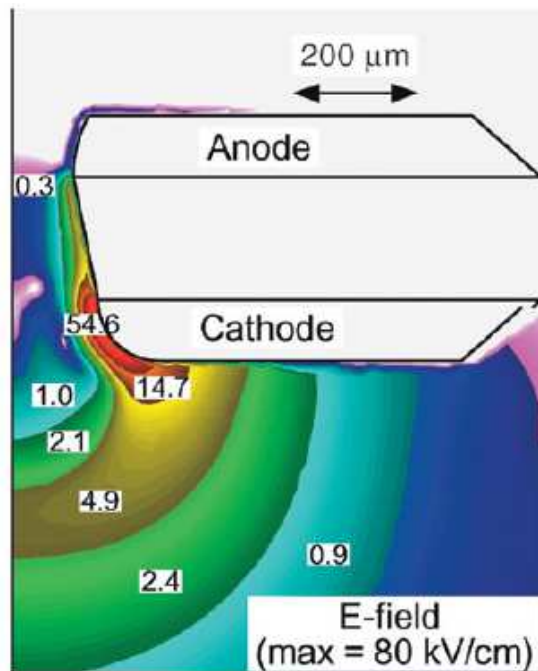


- Some species are fluid-like and others are particle-like → hybrid model
- PIC-MC for particles which need detailed distribution function (often electrons in μ plasmas)
- Fluid model for others species (often ions and neutrals)
- Advantages: accurate + self-consistent + reduced calculation time

Hybrid model: example



- Hybrid model of a MHCD in Ar (hole diameter = 200-300 μm)
- 2D fluid model + MC simulation for beam electrons emitted at the cathode



250 Torr in Ar with $I = 2 \text{ mA}$

Microplasmas for material synthesis

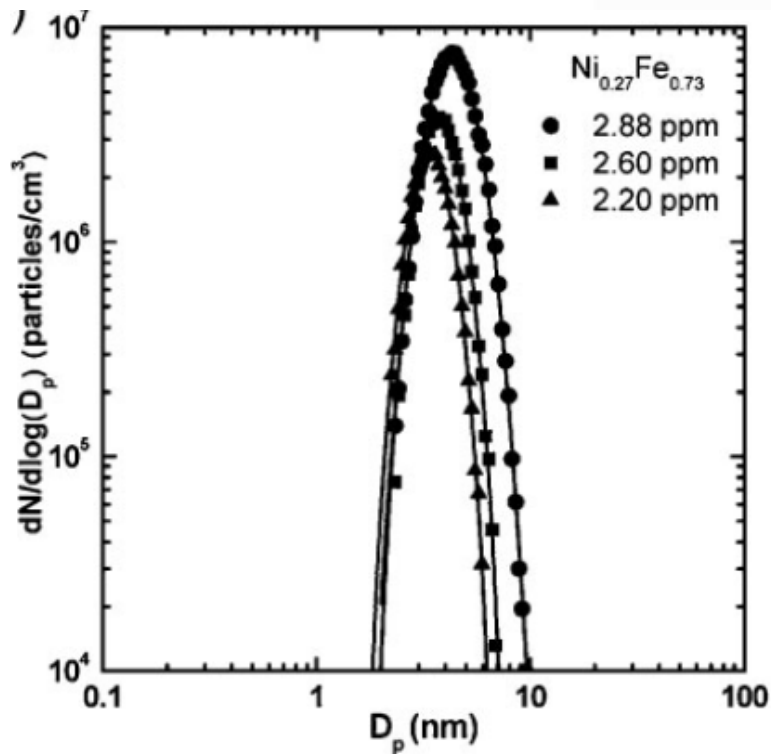
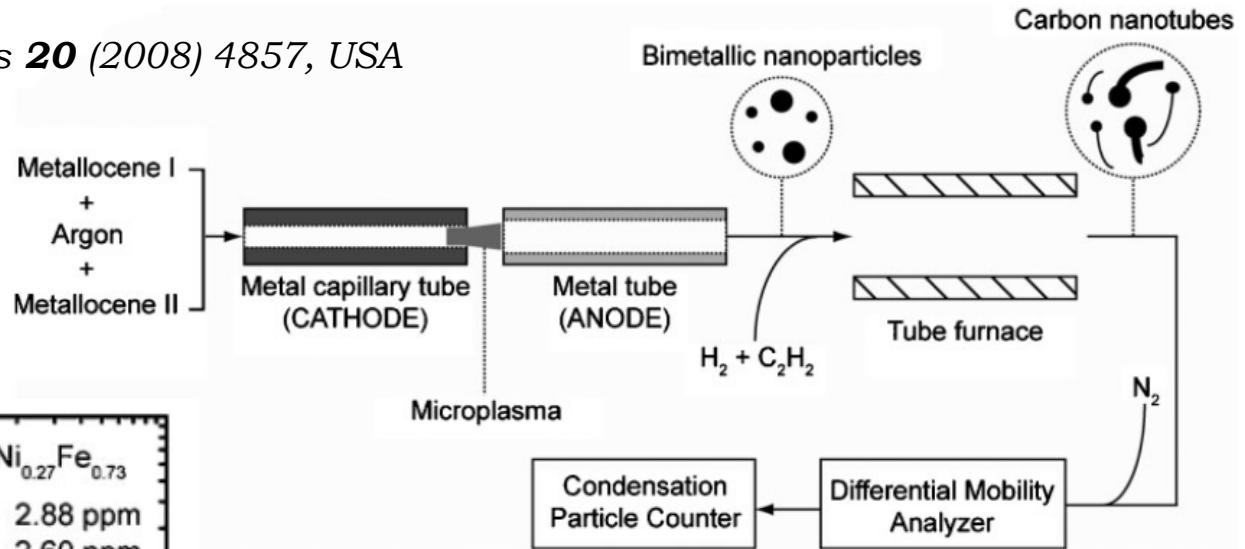


- Synthesis from vapour precursors
- Synthesis from evaporation or sputtering of a sacrificial electrode
- Synthesis from plasma-liquid interactions

Synthesis from vapour precursors: exemple of bimetallic nanoparticles Ni_xFe_{1-x}



Chiang et al., *Advanced Materials* **20** (2008) 4857, USA



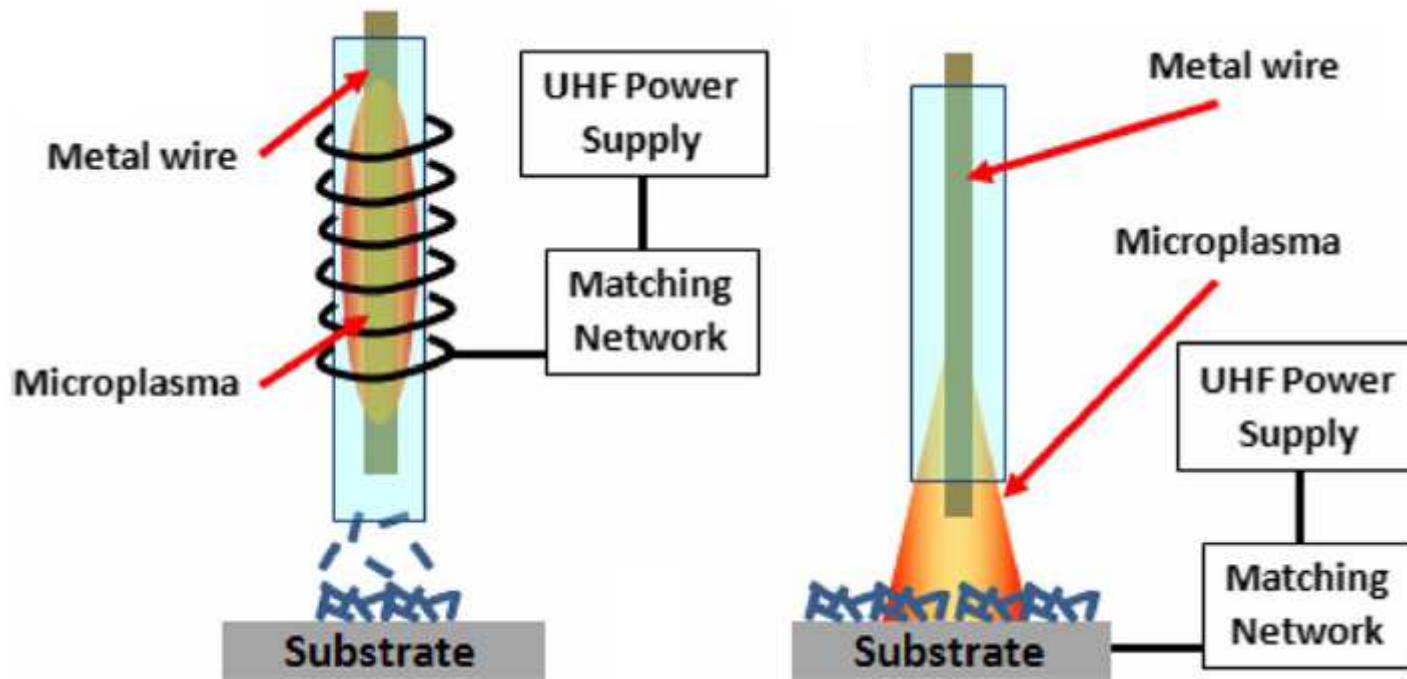
Growth of carbon nanotubes at Patm
 Cathode diameter = 180 μm
 DC power
 Electrode gap = 1 mm
 $I = 7 \text{ mA}$; $V = 400 \text{ V}$

Size distributions
 (differential mobility analyser)

Synthesis from evaporation/sputtering of solid metal electrodes



Mariotti et al., *IEEE Trans. Plasma Sci.* **37** (2009) 1027, UK

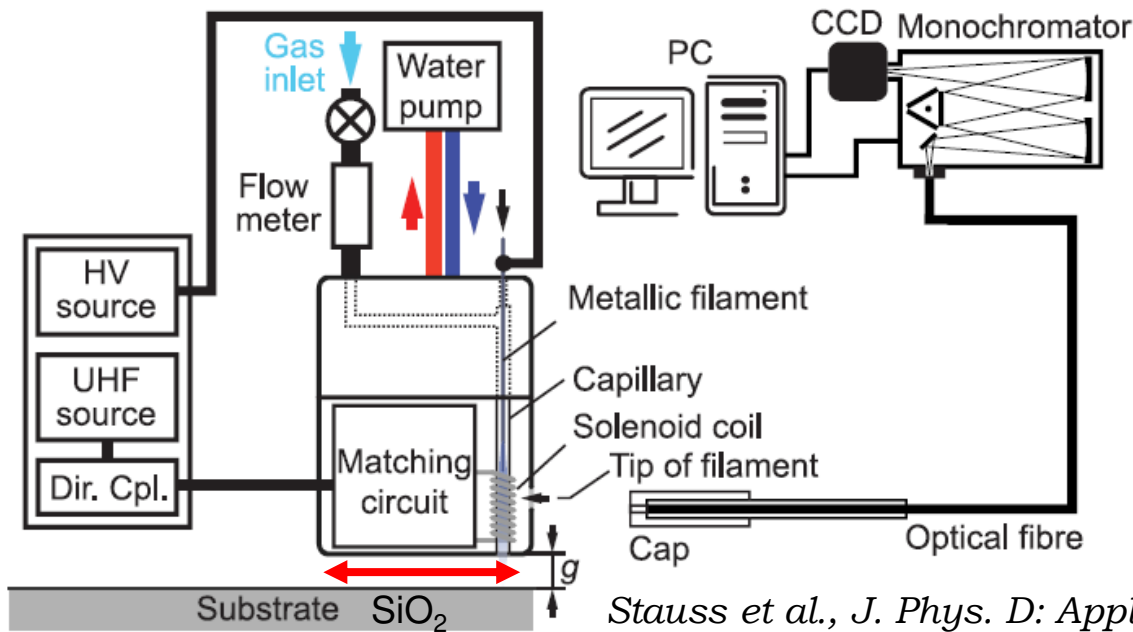


Remote fabrication

Microplasma in direct contact with the processing surface

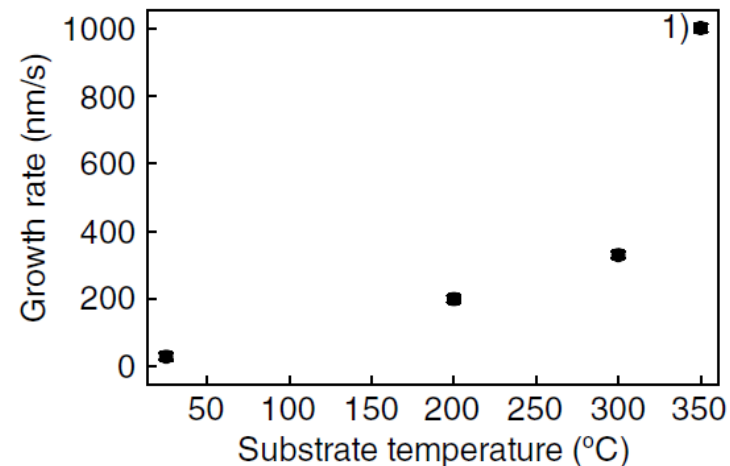
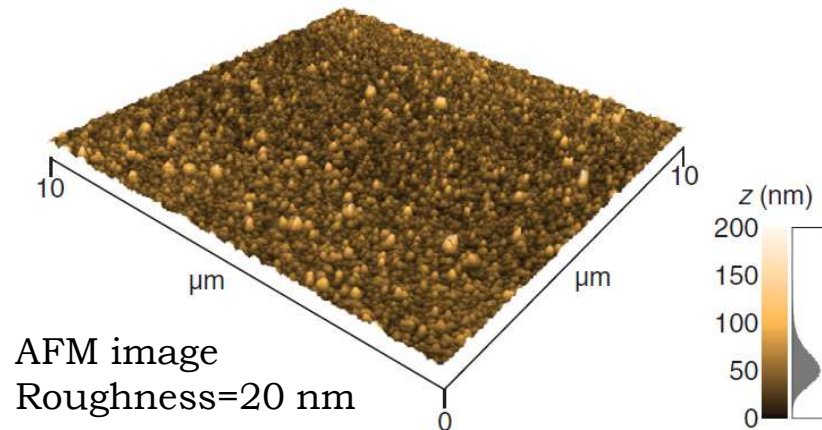
Synthesis of metal and metal-oxide nanostructures

Deposition of ZnO thin films by inductively coupled microplasmas



- Thin filament in quartz capillary
- Zn filament diameter: 0,25 mm
- Ignition: short HV dc (0,5s; 15 kV)
- Plasma sustained by a solenoidal coil (UHF generator, $f=450$ MHz)
- Patm, Ar gas (50-200 sccm), 25W

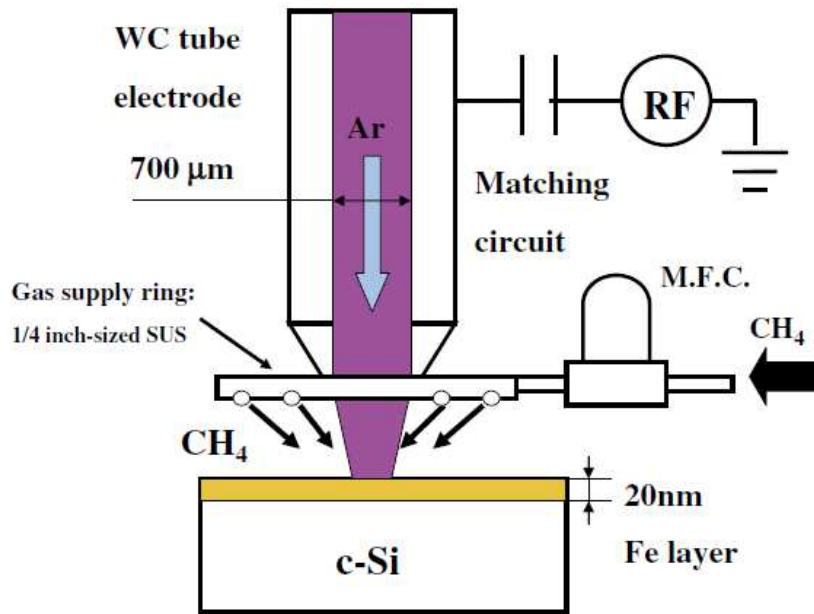
Stauss et al., J. Phys. D: Appl. Phys. 43 (2010) 155203, Japan



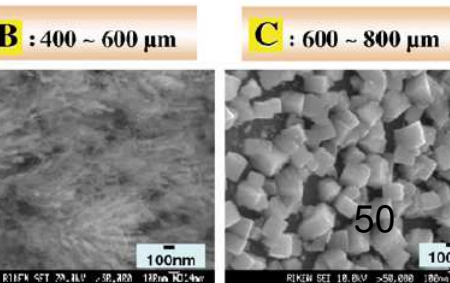
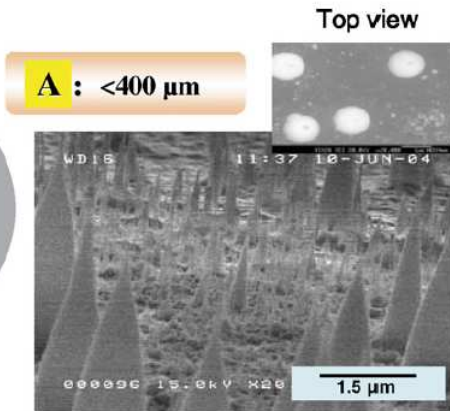
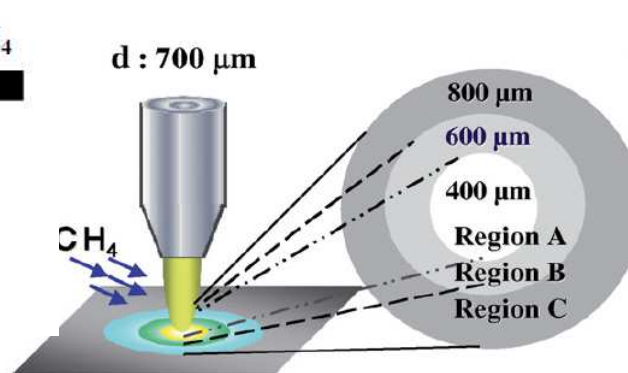
Synthesis of Si nanocones by RF microplasma



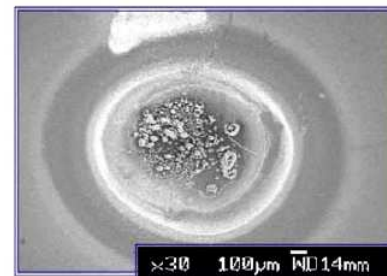
Yang et al., *Thin solid films* **515** (2007) 4158, Japan



- Atmospheric pressure
- Ar/CH₄ mixture
- flow rate: 10-50 sccm
- distance electrode-substrate: 3 mm
- RF power = 35 W



Region A: cone-shaped products (Si)
 Region B: nanowire structures (C)
 Region C: cubic-shape products (tungsten oxide WO_x and diamond particles)

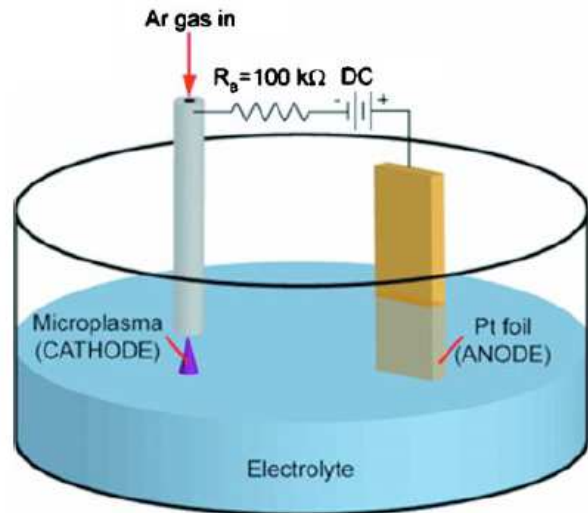


SEM image of products formed underneath the tube electrode

Synthesis by plasma-liquid interactions: example of Ag nanoparticles (SERS)

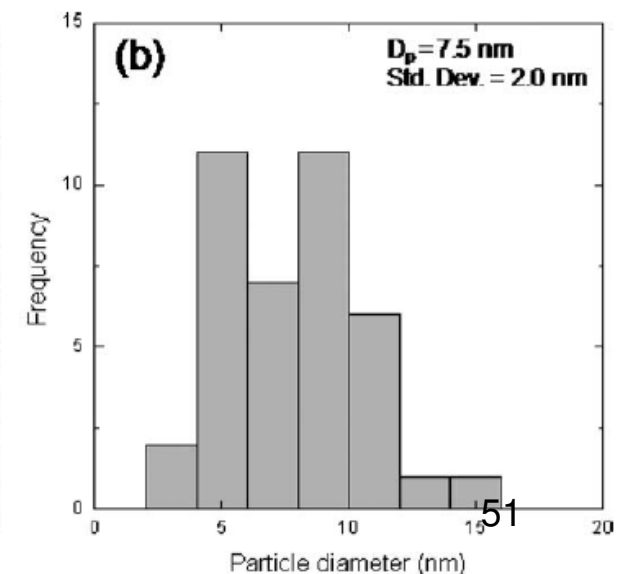
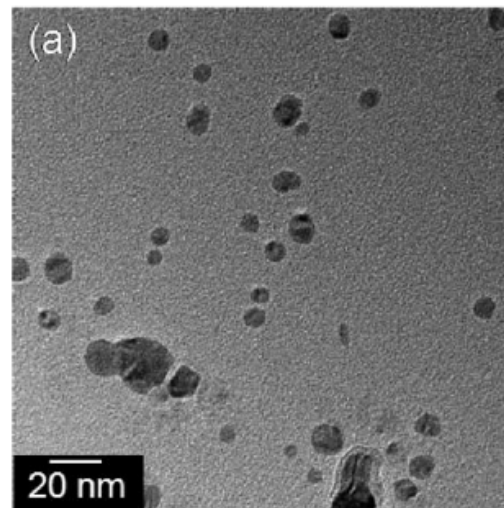


Chang et al., *J. Vac. Sci. Technol.* **28** (2010) L5, USA



- Anode: Pt, 1 cm x 1 cm
- Cathode: stainless-steel capillary, 5 cm x 180 μ m hole diameter, 1 mm from the surface
- Ar flow rate: 25 sccm
- Electrolyte: 1 mmol/L AgNO_3 , target molecule=crystal violet (CV), de-ionized water
- V=290V and I=2mA
- time ranging=1-30 min

TEM images
(process time=10 min)
Mean particle diameter=7,5nm
with standard deviation=2 nm



Microplasmas for environment

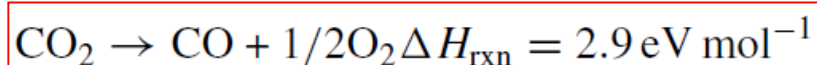
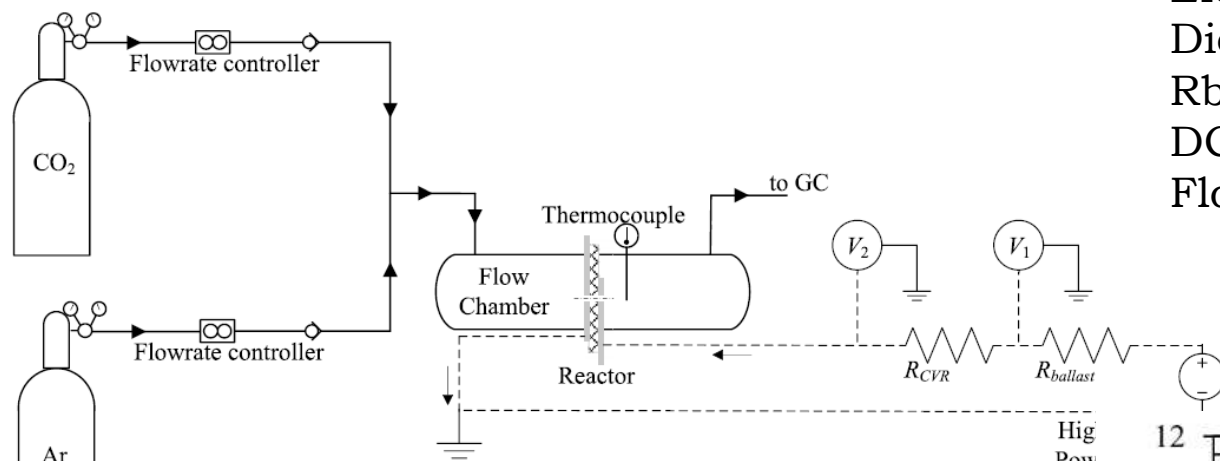


- Dissociation of CO₂
- Decomposition of Toluene

CO₂ dissociation by microplasmas

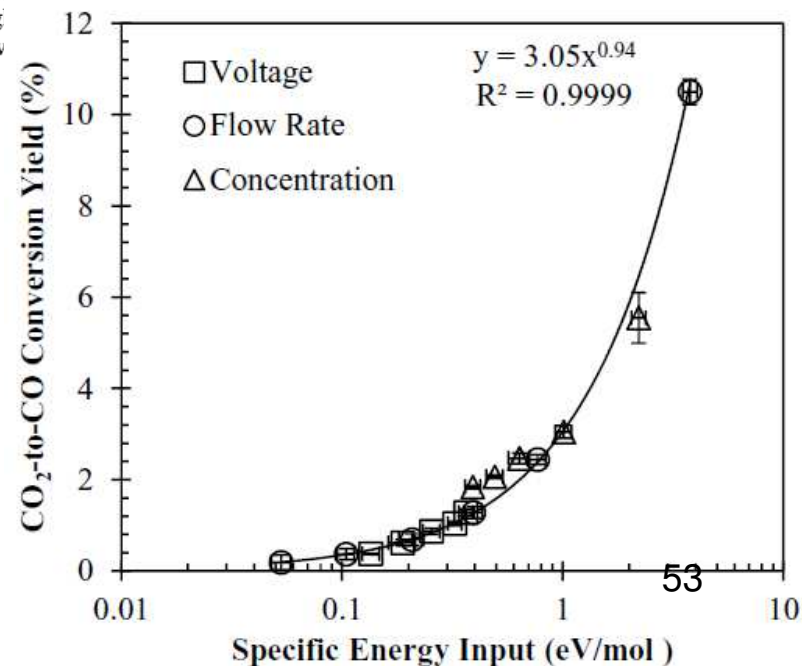
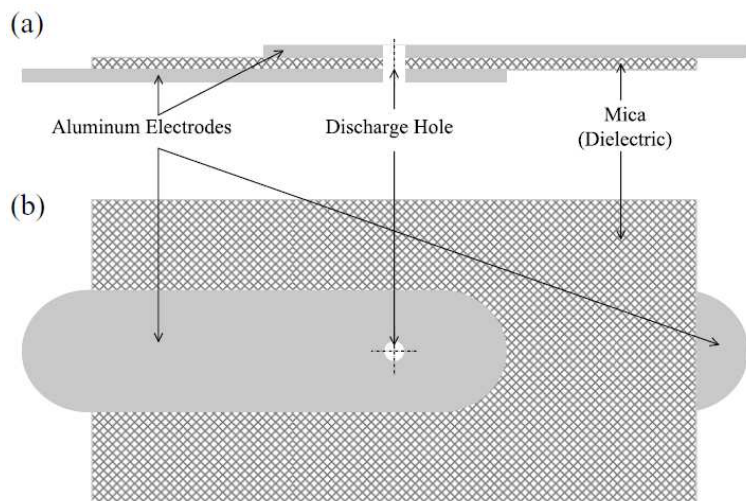


Taylan et al., PSST **24** (2015) 015006, Austin, USA



Patm; hole diameter=400μm
 Electrodes=aluminium 10μm
 Dielectric=mica 150 μm
 Rballast=1 MΩ
 DC voltage=2.5-4.5 kV
 Flow rate=10-800 ml/min

Max yield=10.5%

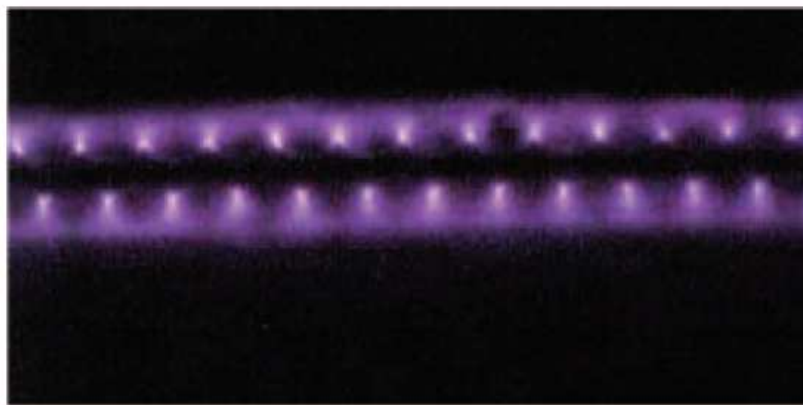
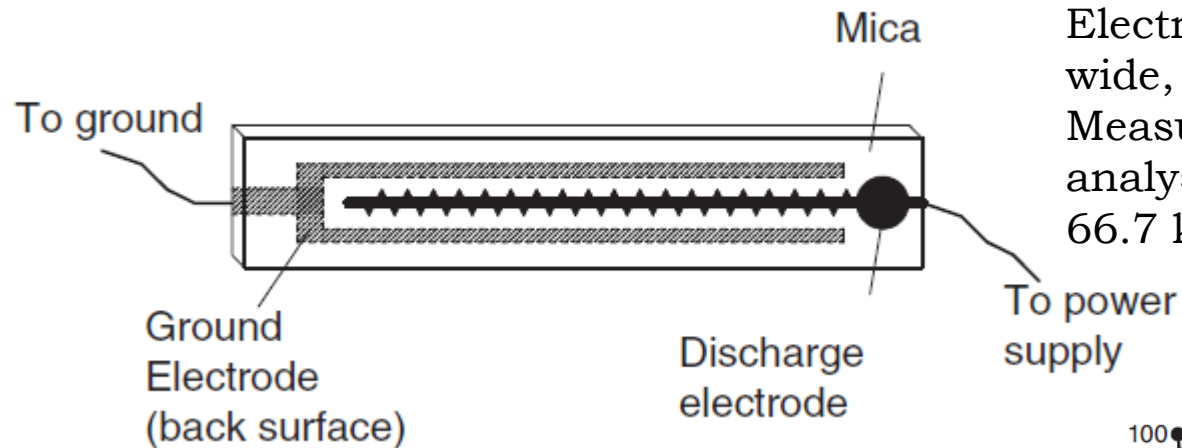


Decomposition of Toluene by microplasmas

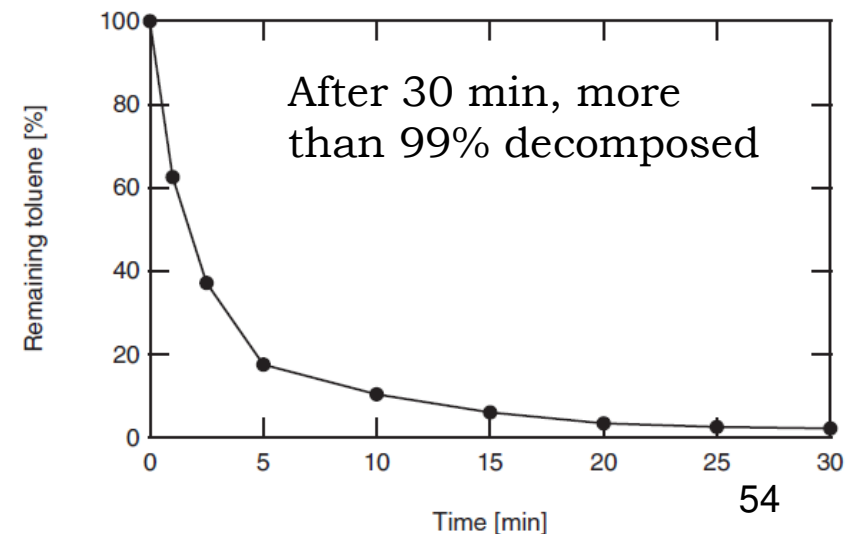


Seto et al., *J. J. Appl. Phys.* **44** (2005) 5206, Japan

Surface discharge microplasma
Rectangular mica sheet: 30 mm long,
20 mm wide, 80 μm thick
Electrodes: stainless steel, 100 μm
wide, 80 μm thick
Measurement with GCMS and CO_2
analysis
66.7 kHz, 3.5 kV



500 μm



Microplasmas for medical applications



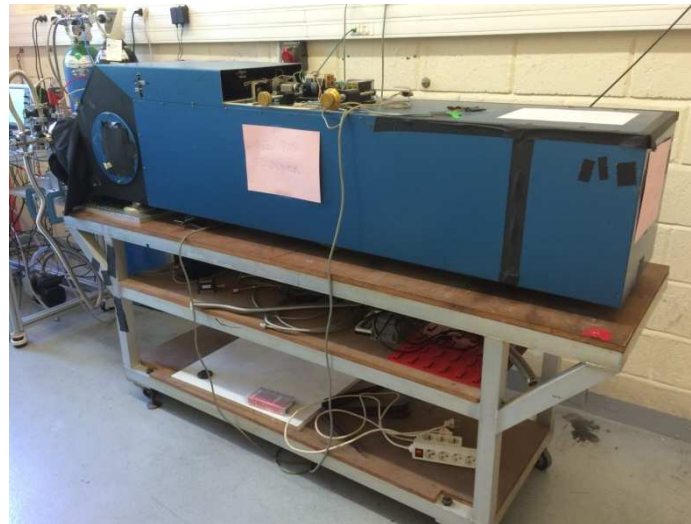
- Joao Santos Sousa's presentation

Merci pour votre attention

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-f = 2 m

-1200 g/mm

-résolution spatiale = 2 μ m

-dispersion spectrale: 0.788
pm/pixel

Spectromètre SOPRA haute-résolution (débitmètre H₂)
(outils réseau plasmas froids)