

# Synthèse de nanoparticules par plasma (essentiellement dans les liquides)

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16 octobre 2018

# 0 - Introduction

## Nanoparticles for what?

Do what chemists cannot do !

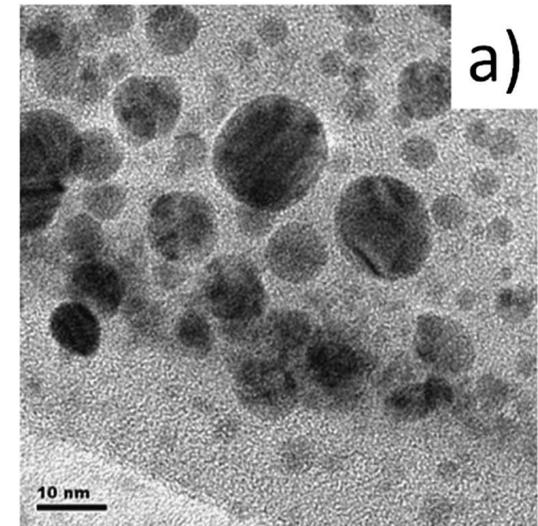
Rules to follow to find the holy grail !

- Rule n°1: benefit from non-equilibrium conditions !
- Rule n°2: benefit from nano-effects, *i.e.* quantum confinement !
- Rule n°3: combine high production yields and narrow size distributions
- Rule n°4: avoid clustering to get isolated nanoparticles.
- Rule n°5: understand to optimize

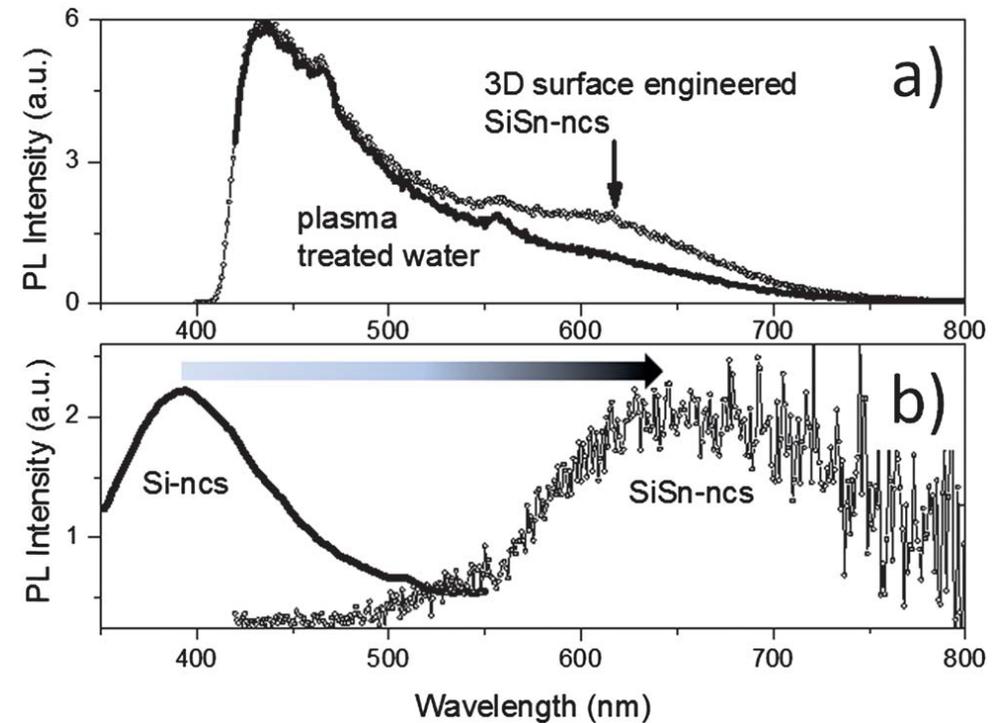
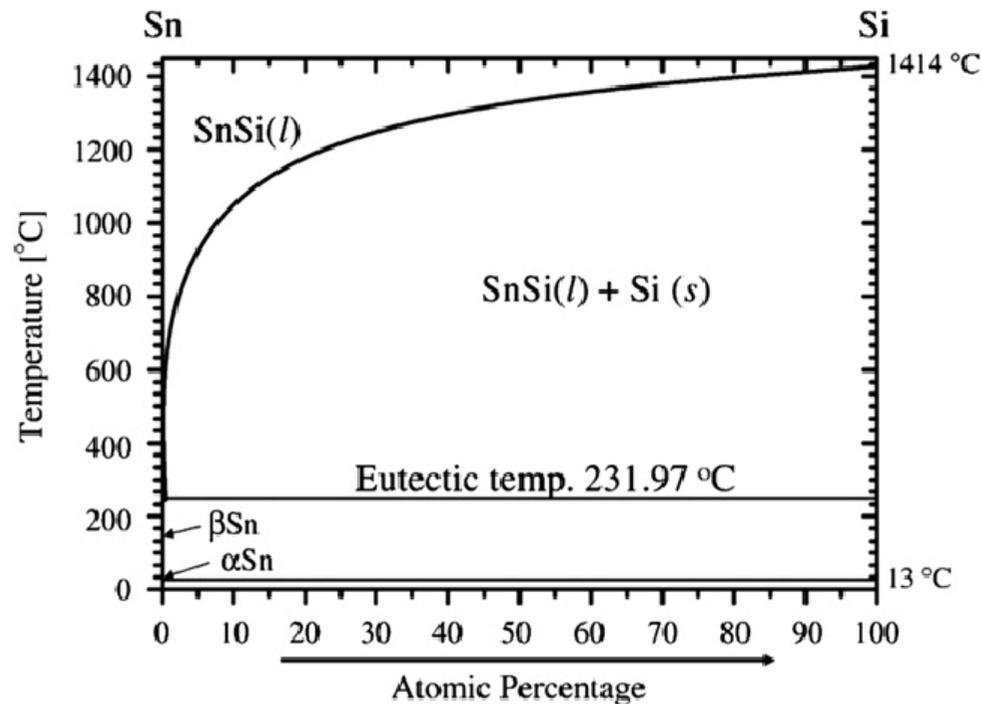
# 0 - Introduction

## Rule n°1: benefit from non-equilibrium conditions !

- Example: mix silicon and tin or silicon and germanium (ns-laser ablation in water)



Švrček et al. Optics express, 17 (2009) 520

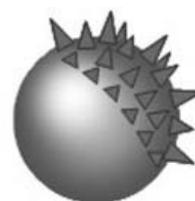


# 0 - Introduction

## Rule n°1: benefit from non-equilibrium conditions !

- Example: exotic nanoparticles

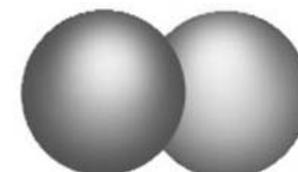
- Alloy NPs
- Bi-materials
- Core-shell structures



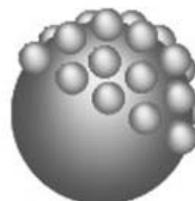
Janus particle



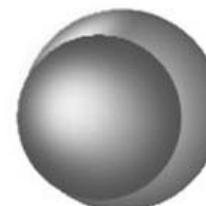
Bicompartmental particle



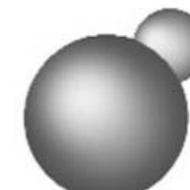
Dumbbell-like particle



Half raspberry-like particle



Acorn-like particle



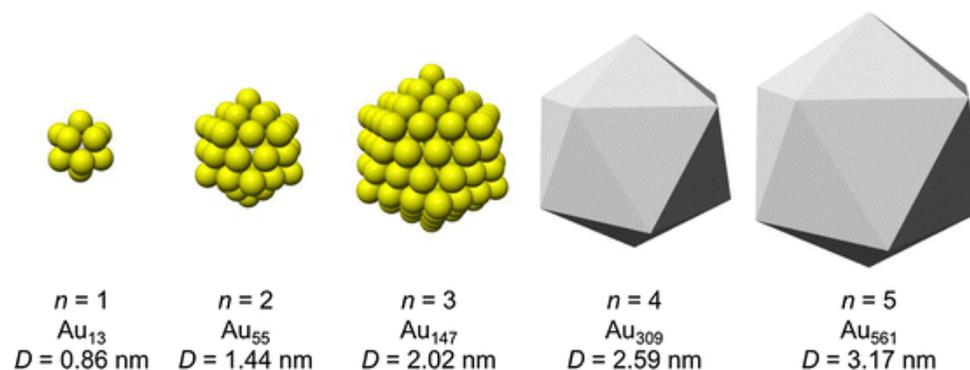
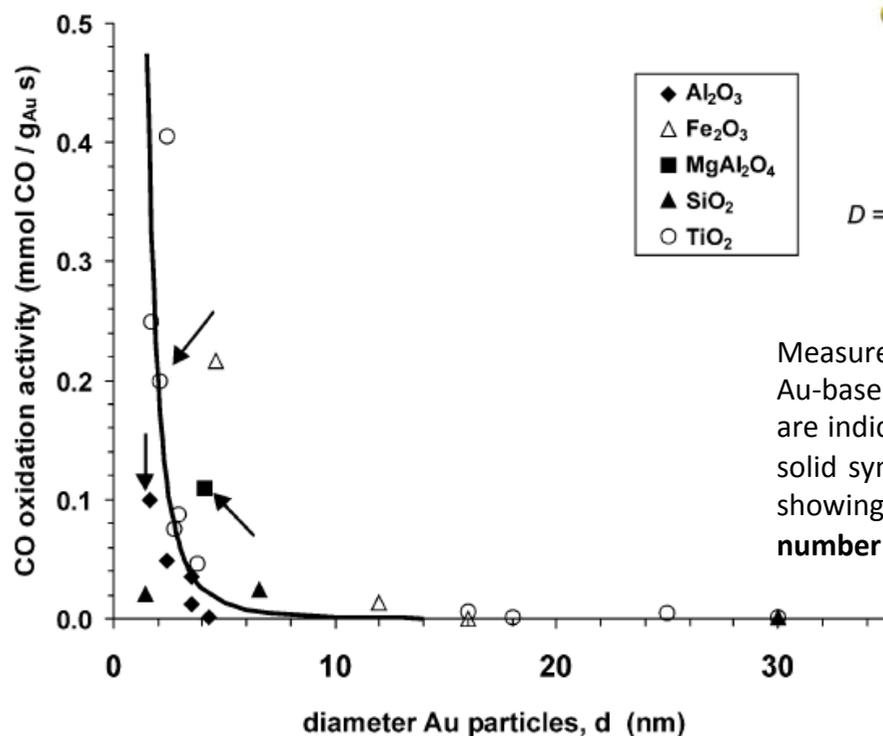
Snowman-like particle

*Perro et al. J. Mater. Chem. 15 3745 (2005)*

# 0 - Introduction

## Rule n°2: benefit from nano-effects, *i.e.* quantum confinement

- Example: Gold catalytic activity

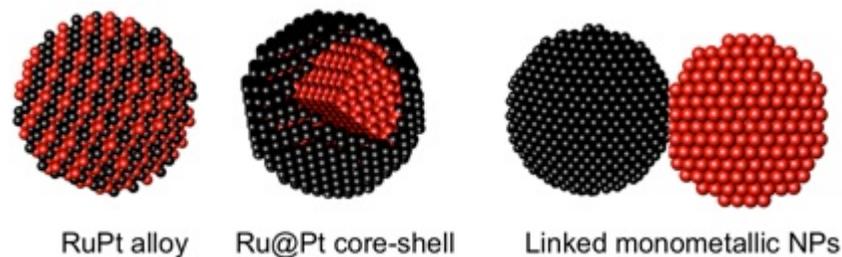
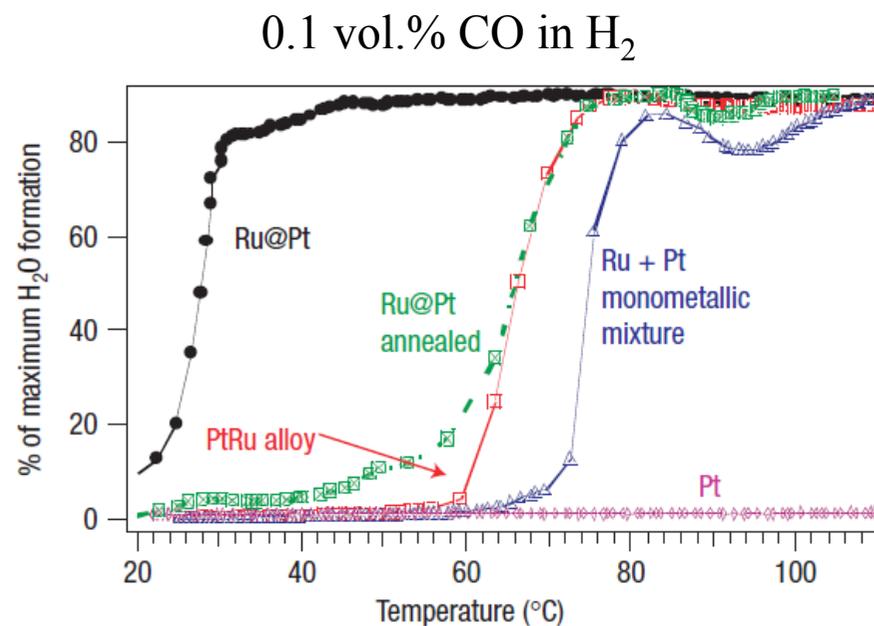


Measured activities (in mmol CO/(gAu s)) for CO oxidation at 273 K over different Au-based catalysts as a function of the average particle size ( $d$ , in nm). Supports are indicated by the symbol shape. Open symbols are used for reducible supports; solid symbols for irreducible supports. The curve shows a  $1/d^3$  guide to the eye, showing that the activity of gold catalysts is approximately **proportional with the number of low-coordinated atoms at the corners of the gold particles..**

*N. Lopez et al., J. Catalysis 223 (2004) 232*

# 0 - Introduction

Rule n°2: benefit from nano-effects, *i.e.* quantum confinement



*Alayoglu et al. Nature Materials 7 333 (2008)*

# 0 - Introduction

## Rule n°3: combine high production yields and narrow size distributions

Monodisperse nanoparticles are arbitrarily defined as a predominately homogeneous population with >90% uniformity in size distribution.

It is not only a matter of size, but also of shape !

Common production yields:

- Laser ablation: 100 mg / h
- Plasma in liquids (liquid conversion): 10 g / h
- Plasma in liquids (electrode erosion): 100 mg / h
- Low pressure CVD : 100 mg / h

# 0 - Introduction

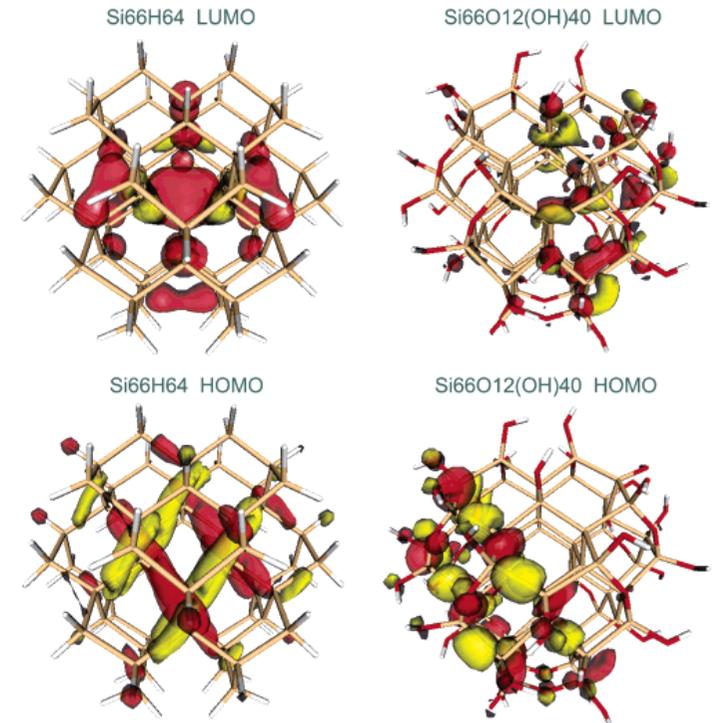
## Rule n°4: avoid clustering to get isolated nanoparticles

- The oxide-passivated nanocrystals have optically forbidden, indirect-gap-type transitions whereas the hydrogen-passivated nanocrystals have optically allowed, direct-gap-type transitions.
- H-passivated Si-nc luminesce in the blue whereas O-passivated Si nanocrystals luminesce in the yellow-red.

DFT calculations for 1.1 to 1.4 nm Si-nc

*Zhou et al. 2003 Nano Lett. 3 163*

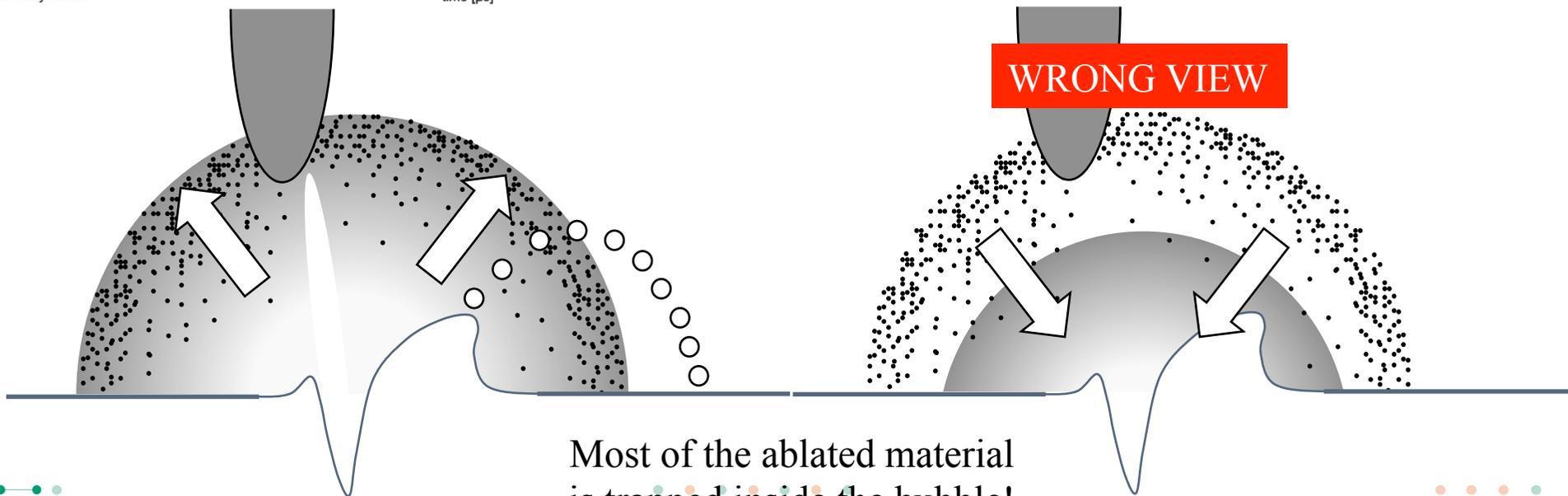
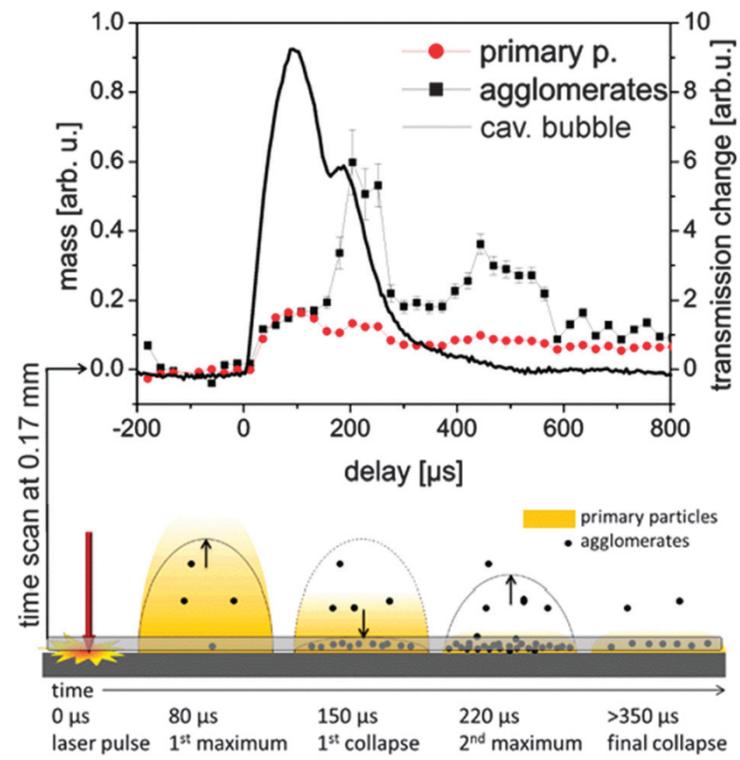
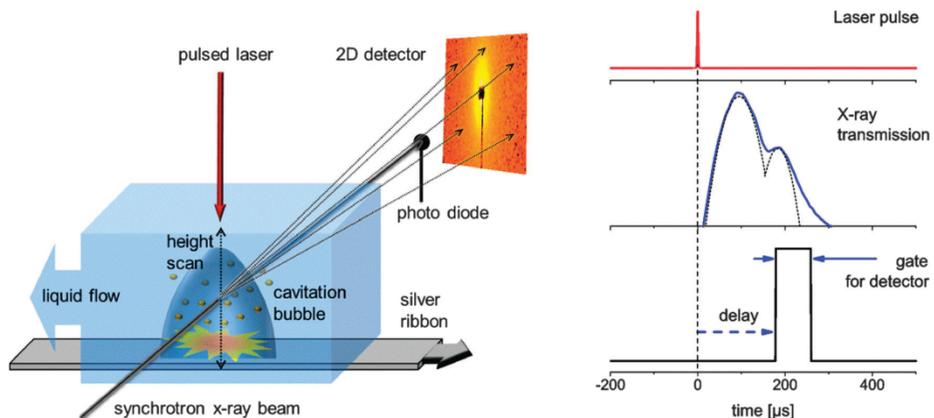
Surface groups strongly affect the final properties of small NPs.



# 0 - Introduction

## Size distributions

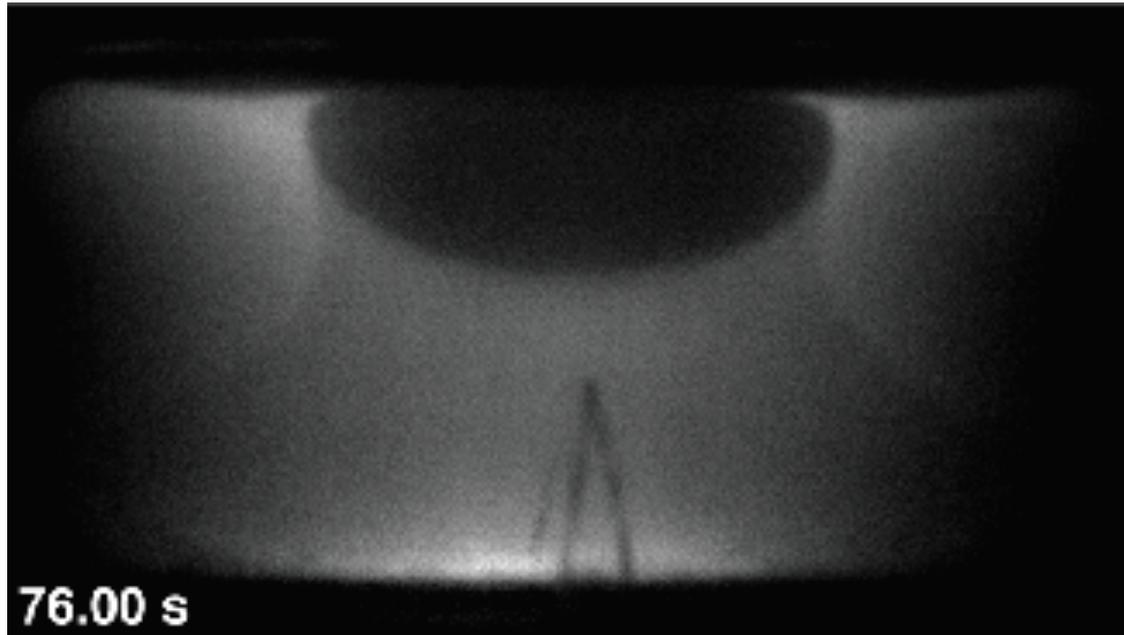
### Rule n°5: understand to optimize



Most of the ablated material is trapped inside the bubble!

# Chap. I

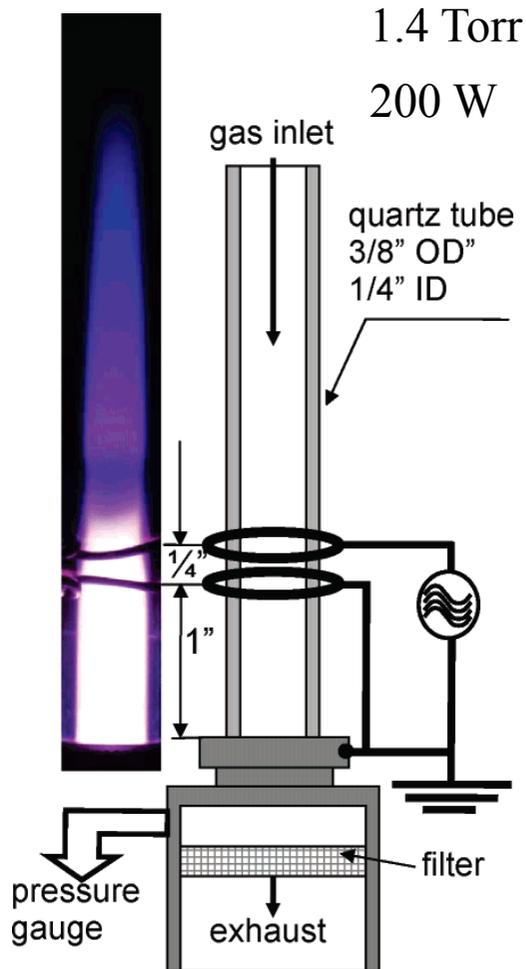
## Nanoparticles and discharges in gases



van de Wetering *et al.* 2015 *J. Phys. D: Appl. Phys.* **48** 035204

# 1 - Nanoparticles and low pressure discharges

High-rate synthesis of Si-nc for PV applications

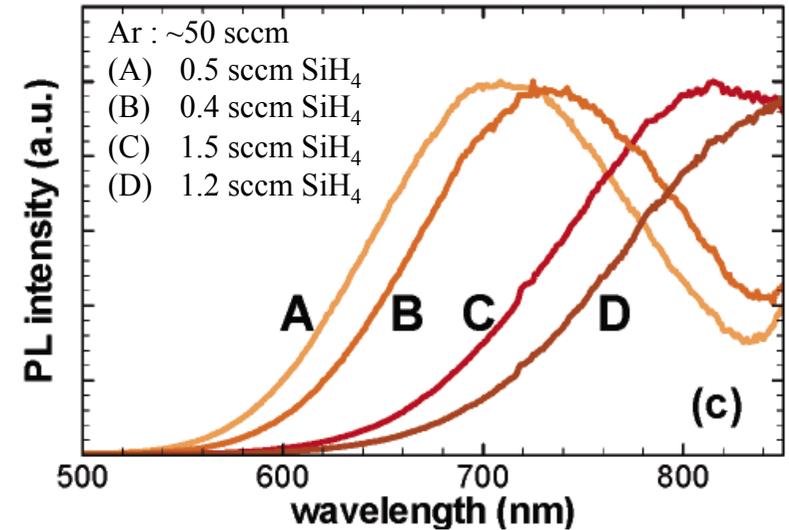


**Photoluminescence quantum yields usually < 30% without post-treatment**

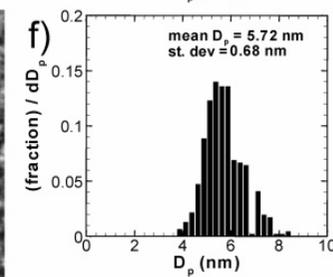
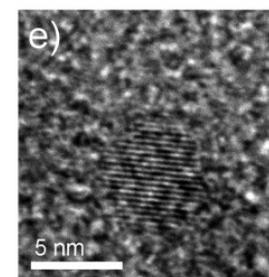
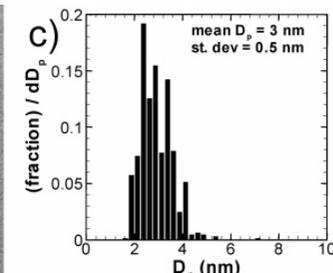
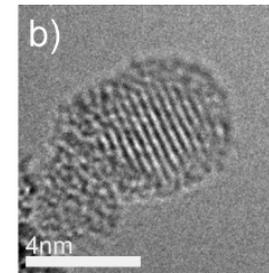


Fast quenching of the particle temperature  
+  
High flux of atomic hydrogen to the nanocrystal

Anthony *et al.* 2011 *Adv. Funct. Mater.* **21** 4042



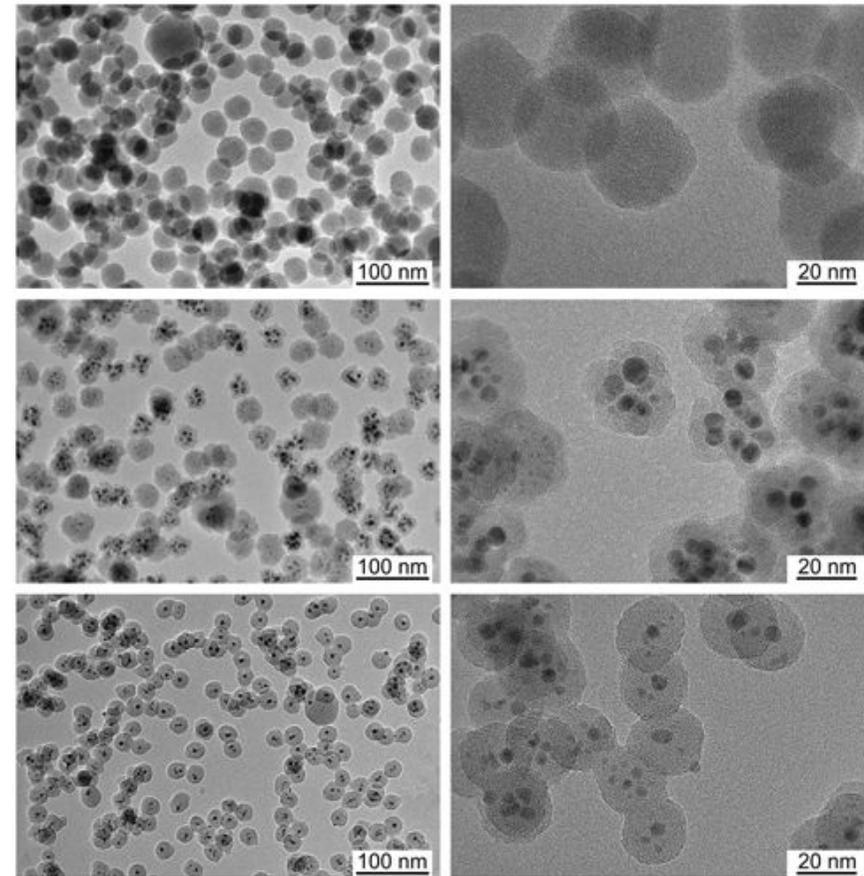
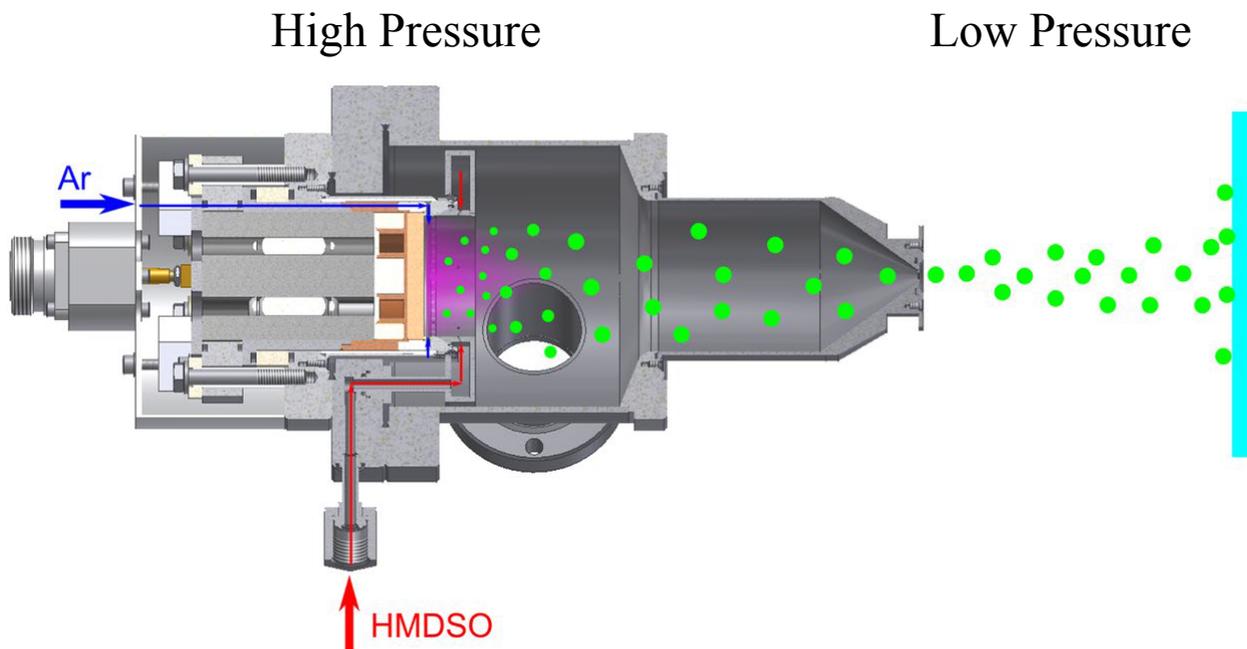
Excellent size distribution



Mangolini *et al.* 2005 *Nano Lett.* **5** 655

# 1 - Nanoparticles and low pressure discharges

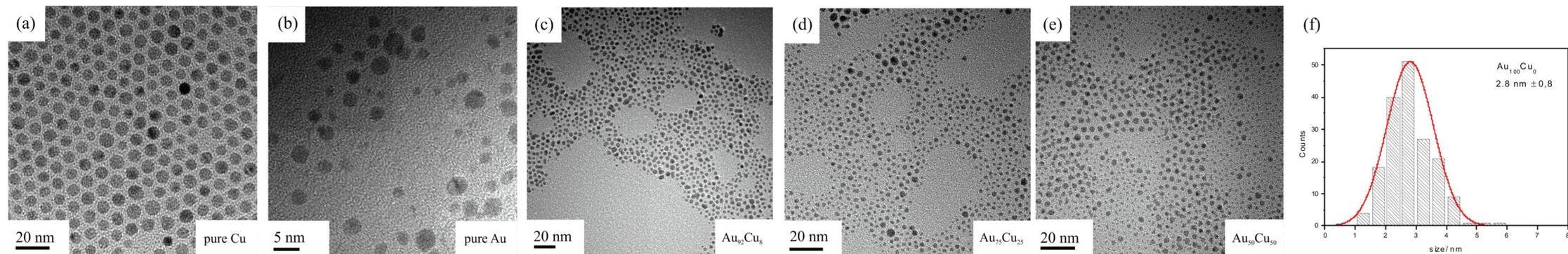
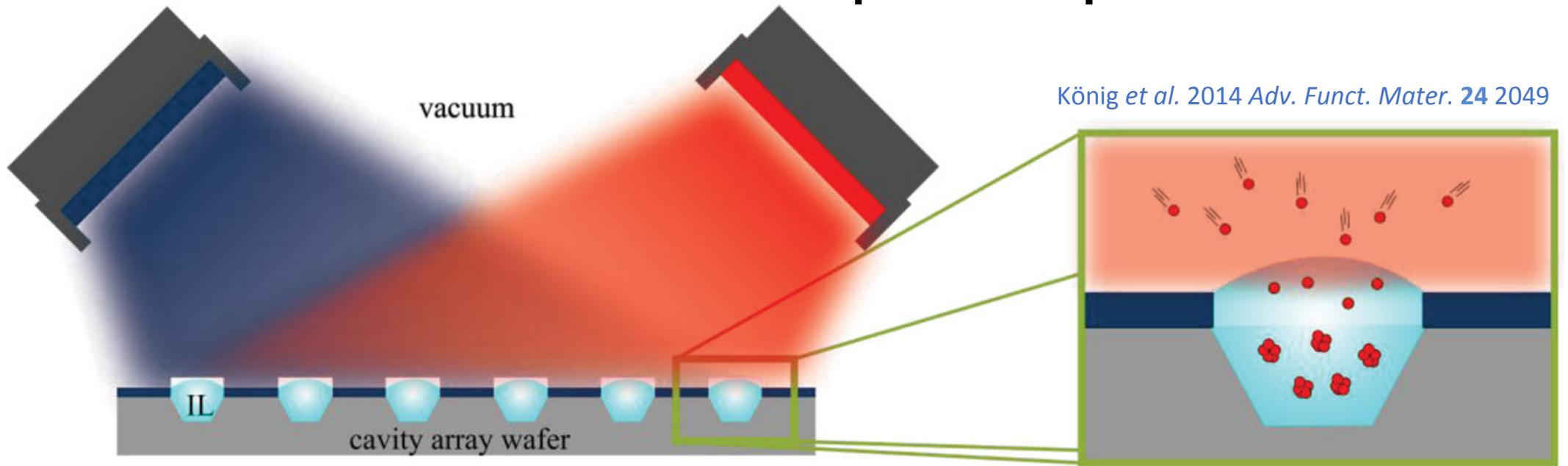
## Dual chamber Physical Vapor Deposition



Solař et al. 2017 *Scientific Reports* 7 8514

# 1 - Nanoparticles and low pressure discharges

Pulvérisation combinatoire sur liquide ionique [C<sub>1</sub>C<sub>4</sub>im][Tf<sub>2</sub>N]



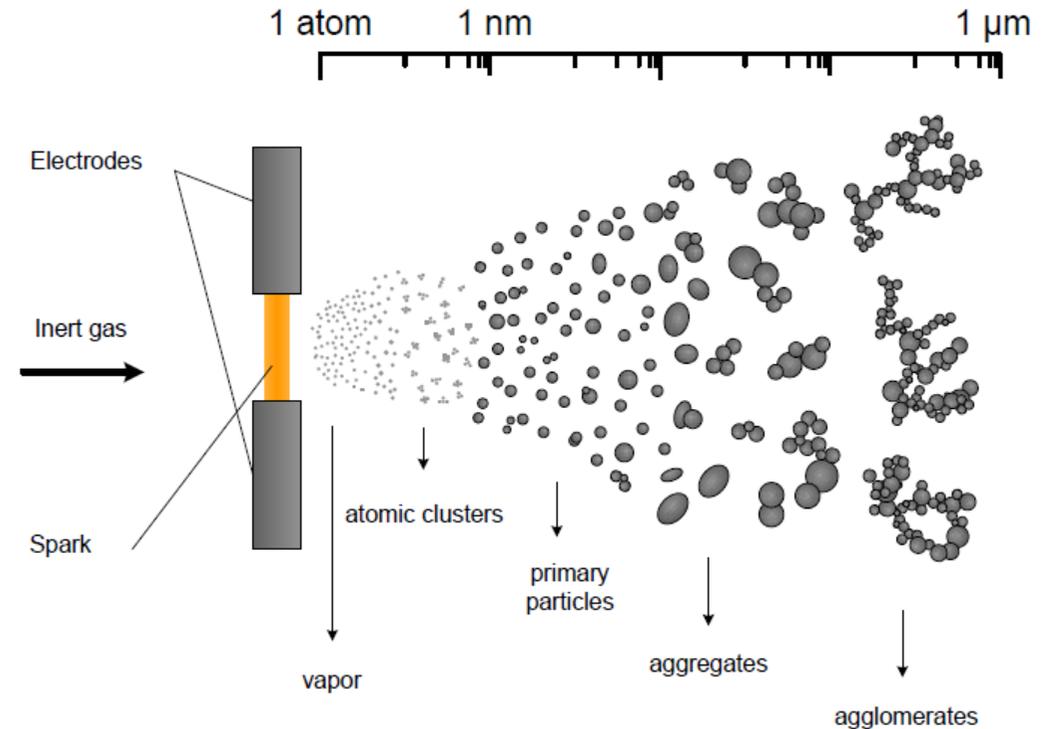
# 1 - Nanoparticles and high pressure discharges

Spark generator  
(thermal plasma - low-powered arc)

Low frequency: yield  $\sim 1$  mg/h

Local melting of electrodes

Why not simply melting electrodes?



T: 015-2786751 [www.patent.tudelft.nl](http://www.patent.tudelft.nl)

evaporation/condensation method:

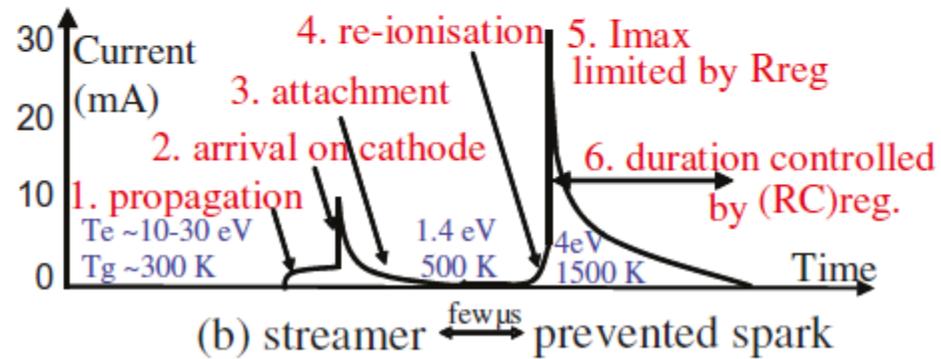
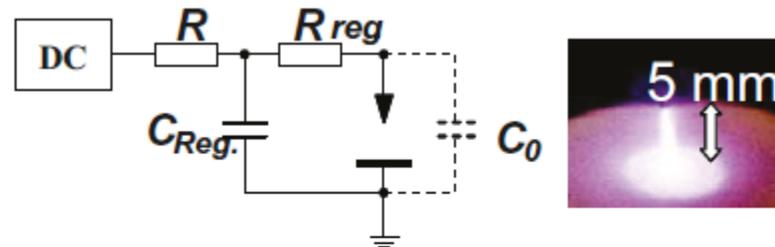
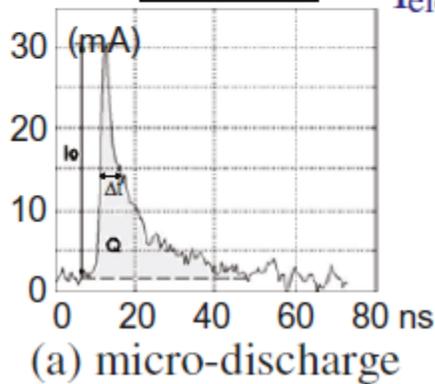
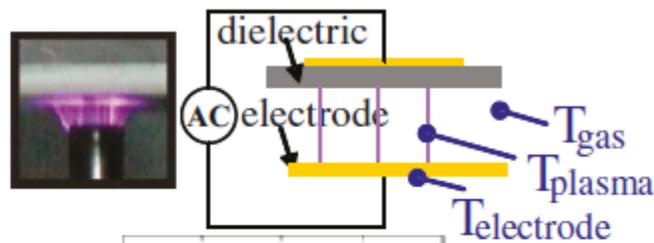
PROS - good yield and control over deposition parameters

CONS compared with spark discharge - furnace energy consumption and heating up and cooling down times, clogging of the furnace due to particle deposition on the walls, etc.

# 1 - Nanoparticles and high pressure discharges

Non-thermal generator  
Stop before or prevent arcing

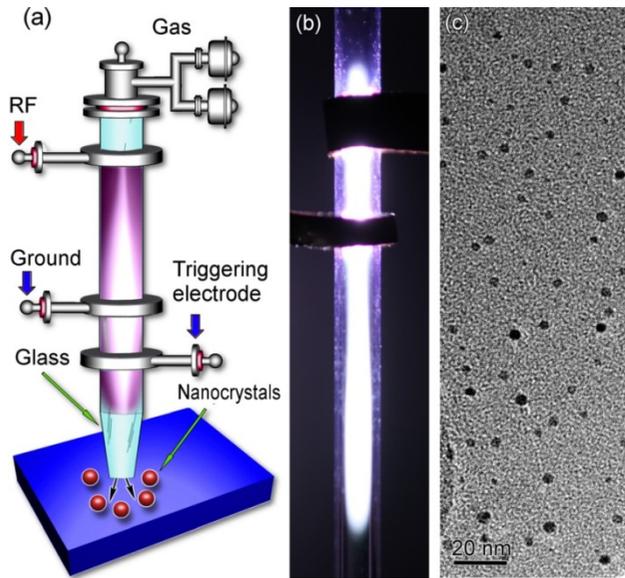
J.P. Borra *et al.* 2006 *J. Phys. D: Appl. Phys.* 39 R19



**LOW THROUGHPUT**

# 1 - Nanoparticles and high pressure discharges

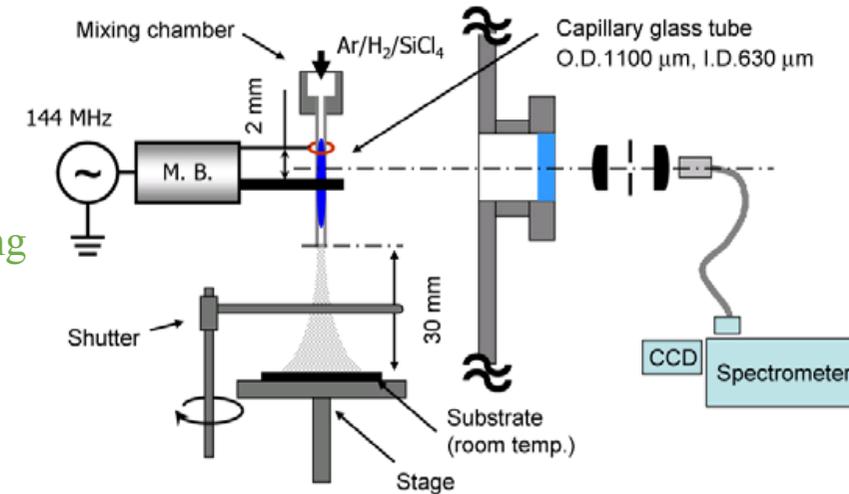
Askari *et al.* 2014 *Appl. Phys. Lett.* **104** 163103



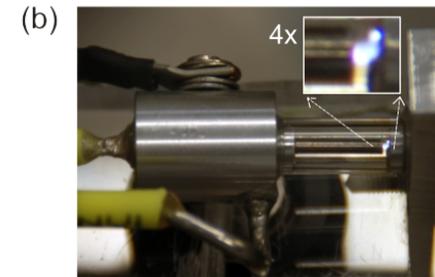
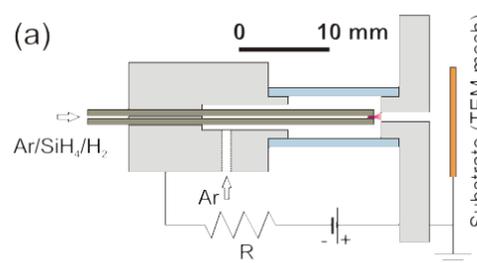
RF power: 100W  
250 sccm Ar  
SiH<sub>4</sub>: 10 ppm  
1mm electrode spacing

$\mu$ wave power: 35W  
200 sccm Ar  
H<sub>2</sub> < 5%  
SiCl<sub>4</sub>: 100 ppm  
2mm electrode spacing

DC power: 2W  
100 sccm Ar  
H<sub>2</sub> < 5000 ppm  
SiCl<sub>4</sub>: 1.5-2.5 ppm  
180  $\mu$ m electrode spacing



Nozaki *et al.* 2007 *Nanotechnol.* **18** 235603

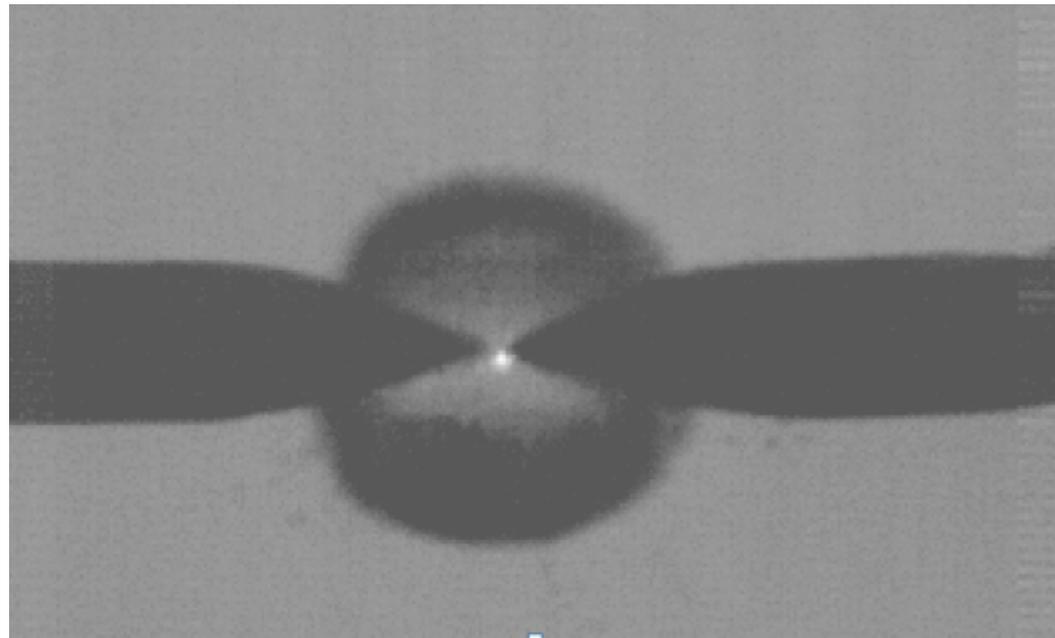


Barwe *et al.* 2015 *Plasma Process. Polym.*, **12** 132

**HIGH THROUGHPUT**

# Chap. II

## Discharges in liquids



# 2.1 – Spark discharge in dielectric liquids

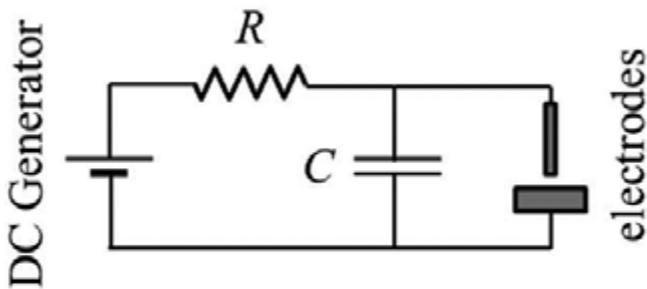
Several possibilities:

*by electrical characteristics*

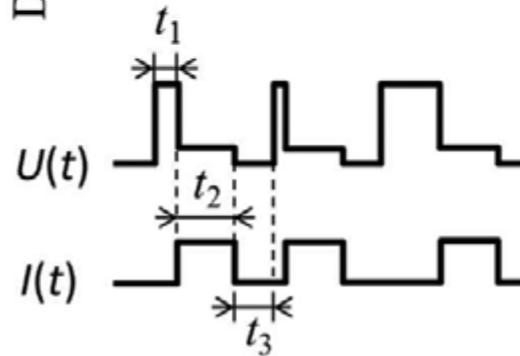
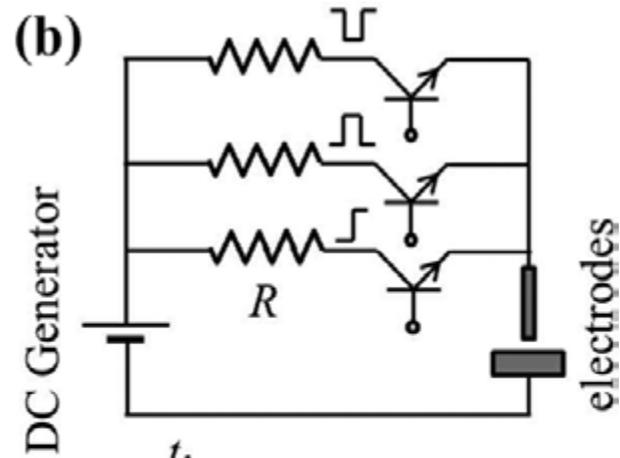
*by liquids*

*by inter-electrode gap distances*

(a)



(b)



## 2.1 – Spark discharge in dielectric liquids

- igniting a discharge,
- maintaining the discharge for a given time and, eventually, control the voltage or the current during this step,
- setting an off-time to recharge capacitors and flush the liquid,
- repeating the process as long as necessary

### Dielectric or conductive liquid?

The conductance of the liquid is NOT an intrinsic characteristic if the liquid contains non negligible impurity levels.

Pre-treatments of both the liquid and its container to remove dust, particles, dissolved gases, electrolytes, etc.

If the rate constant for the electron–impurity interaction is, say,  $10^{-8} \text{ cm}^3 \text{ s}^{-1}$ , an electron lifetime of typically  $10^{-4} \text{ s}$ , the impurity concentration should not exceed  $10^{-9} \text{ mol l}^{-1}$ , which is particularly weak ( $6 \times 10^{11} \text{ cm}^{-3}$  vs  $3.3 \times 10^{22} \text{ cm}^{-3}$  for water)

# 2.1 – Spark discharge in dielectric liquids

Several possibilities:

*by electrical characteristics*

*by liquids*

*by inter-electrode gap distances*

**Table 1.** Examples of electrical characteristics found for discharges in liquids.

Shape	Liquid	Applied voltage (kV)	Current (A) and pulse width	Gap distance (mm)	Reference
Relaxation	Distilled water	+112	Unknown (~500 ps)	4	Starikovskiy <i>et al</i> (2011)
Relaxation	Distilled water	-35	250 (~200 ns)	0.4	Schoenbach <i>et al</i> (2008)
Relaxation	Oil ITO 100	+4-6	100 (~1.2 $\mu$ s)	0.3	Kudelcik <i>et al</i> (2010)
Relaxation	Liquid Helium	23	280 (2.5 ms)	5	Hayakawa <i>et al</i> (1995)
Relaxation	n-Hexane	15	140 (10 ns)	1	Fuhr <i>et al</i> (1986)
Relaxation	n-Decane	0.022	$8 \times 10^{-8}$ (0.6 s)	$2 \times 10^{-5}$	Virwani <i>et al</i> (2007)
Transistor	Oil Daphne Cut HL-25	+0.110	4.2 (~250 ns)	0.25-1	Kurnia <i>et al</i> (2008)
Transistor	Liquid nitrogen	$\pm 0.220$	10.5 (32 $\mu$ s)	0.3	Muttamara and Fukuzawa (2012)

# 2.1 – Spark discharge in dielectric liquids

Several possibilities:

*by electrical characteristics*

*by liquids*

*by inter-electrode gap distances*

Dividing dielectric liquids into three groups:

- Non-polar liquids such as liquefied gases, hexane, benzene and mineral oils.
- Polar liquids with very high permittivity, which have to be distinguished depending on whether they can be self-dissociated (self-ionized) or not.
- Others
  - The former ones, such as water or ethanol, have a high permittivity. A part of their molecules spontaneously dissociates and recombines partially, giving a certain resistivity.
  - The latter ones have either a high permittivity, such as nitrobenzene or propylene, or a low one, such as chlorobenzene.

# 2.1 – Spark discharge in dielectric liquids

Several possibilities:

*by electrical characteristics*

*by liquids*

*by inter-electrode gap distances*

**Table 2.** Some physical data of dielectric liquids at 1 atm and 293 K (except for liq.N<sub>2</sub> (75 K), liq.He (5 K) and liq.Ar (85 K)); m.p.: melting point; b.p.: boiling point; B.S.: breakdown strength.

Liquid	$\epsilon_r$	Ion mobility (m <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )	$\lambda$ (W m <sup>-1</sup> K <sup>-1</sup> )	m.p.–b.p. (K)	$\mu$ (kg m <sup>-1</sup> s <sup>-1</sup> )	B.S. (MV cm <sup>-1</sup> )	Reference
<i>Non-polar liquids</i>							
Liquid N <sub>2</sub>	1.44	$2.50 \times 10^{-7}$	0.137	63.1–77.4	$1.66 \times 10^{-4}$	1.6–1.9	Henson (1964)
Liquid He	1.05	$4.62 \times 10^{-6}$	0.0198	b.p.: 4.2	$3.32 \times 10^{-6}$	0.7	Donnelly and Barengi (1998)
Liquid Ar	1.6	$6.00 \times 10^{-8}$	0.132	83.8–87.2	$2.78 \times 10^{-4}$	1.1–1.42	Henson (1964)
Hexane	2.0	$1.90 \times 10^{-8}$	0.124	177.9–341.8	$3.09 \times 10^{-4}$	1.1–1.3	Gray and Lewis (1969)
Benzene	2.3	$2.76 \times 10^{-7}$	0.167	279.2–353.2	$6.50 \times 10^{-4}$	1.1	Huang and Freeman (1980)
Vegetal (castor) oil	3.2	$1.60 \times 10^{-10}$	0.17	260 – >660	0.985	1.0	Yang <i>et al</i> (2012)
Mineral (white) oil	2.2	$2.60 \times 10^{-9}$	0.11	264–648	0.020	1.0	Yang <i>et al</i> (2012)
<i>Polar liquids</i>							
Pure water	80	$2.00 \times 10^{-7}$	0.609	273.2–373.2	$1.00 \times 10^{-3}$	0.65	Light and Licht (1987)
Ethanol	23	$2.10 \times 10^{-7}$	0.171	159.2–351.5	$1.20 \times 10^{-3}$	1.65	Atten and Gosse (1969)
Chlorobenzene	5.69	$2.54 \times 10^{-8}$	0.132	228.2–404.6	$1.06 \times 10^{-3}$	0.95	Barret <i>et al</i> (1975)
Nitrobenzene	35	$2.00 \times 10^{-7}$	0.149	278.8–484.1	$1.96 \times 10^{-3}$	0.5–1	Atten and Gosse (1969)
Propylene carbonate	64.4		0.165	224–515	$2.48 \times 10^{-3}$	2.2	Atten and Gosse (1969)

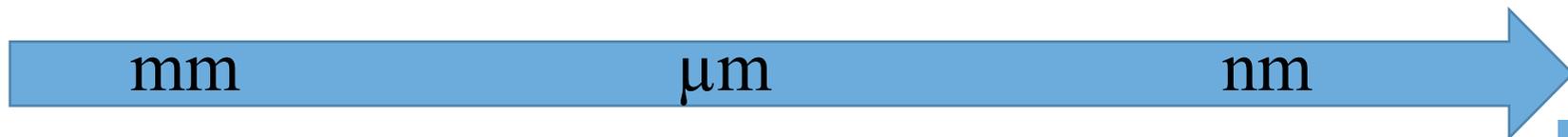
# 2.1 – Spark discharge in dielectric liquids

Several possibilities:

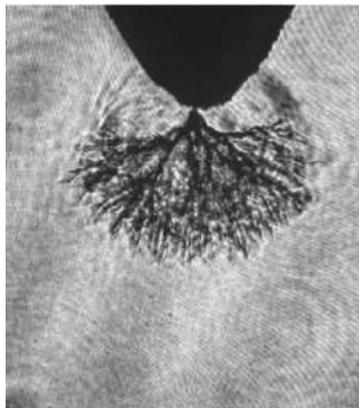
*by electrical characteristics*

*by liquids*

*by inter-electrode gap distances*



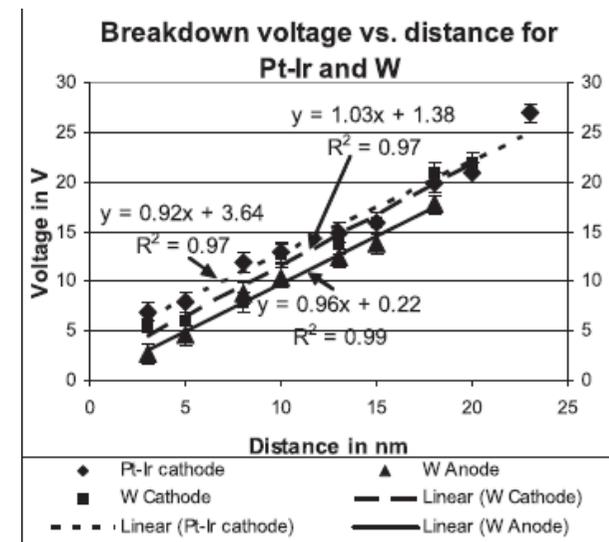
Bush-like structures



filamentary structures



Unknown structures



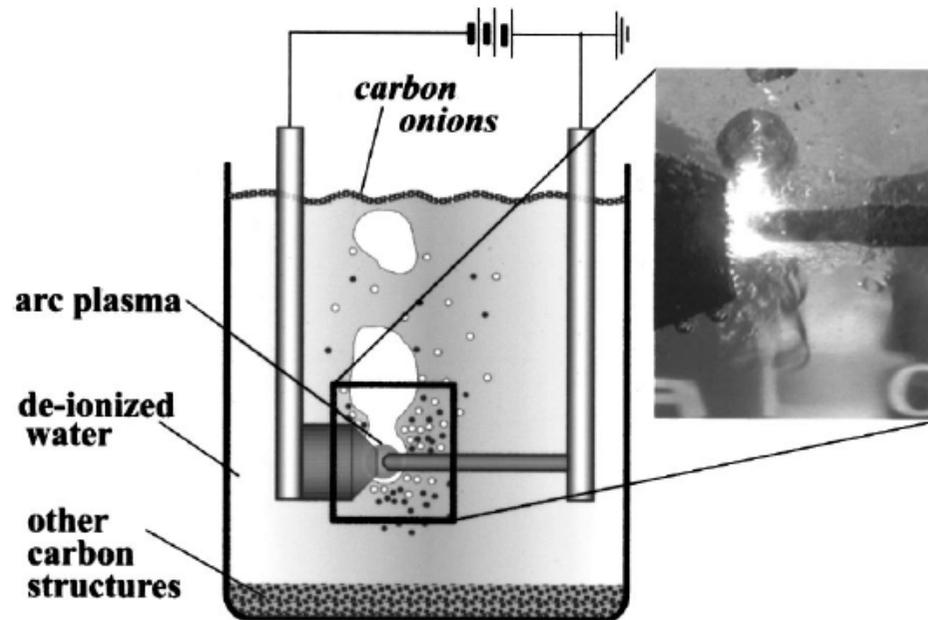
field-emission-assisted avalanche

Virwani et al. PRL 99, 017601 (2007)

# 2.1 – Spark discharge in dielectric liquids

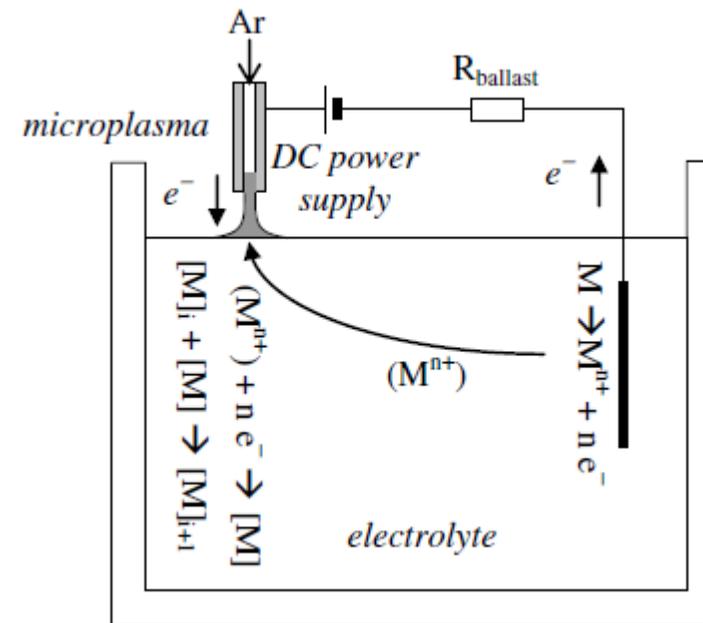
## Submerged discharges

Low voltage: breakdown by contact



Sano et al. 2002 *J. Appl. Phys.* 92 2784

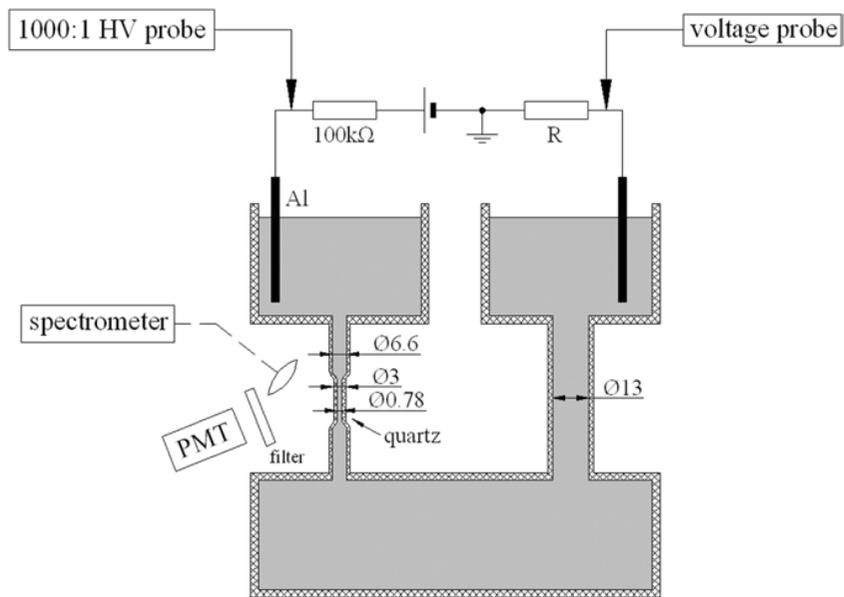
## Discharges in contact with liquids



Chiang et al. 2010 *Plasma Sources Sci. Technol.* 19 034011

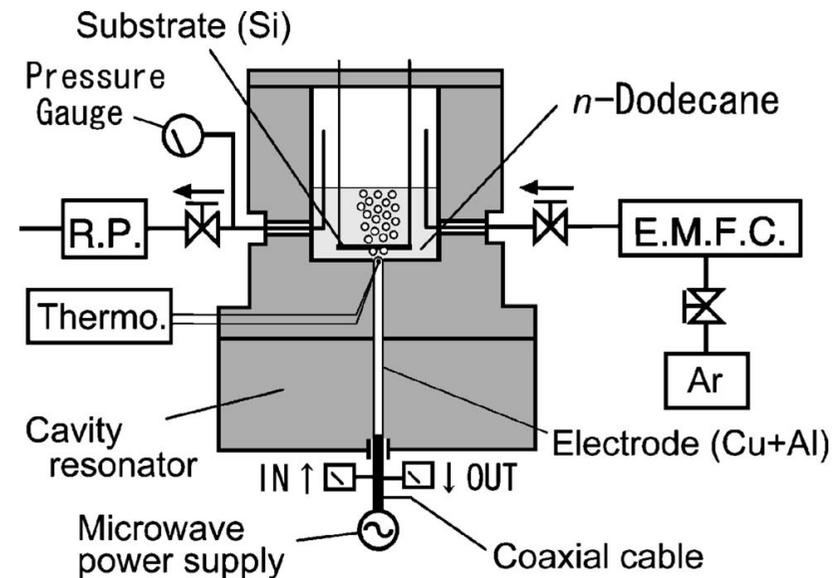
# 2.1 – Spark discharge in dielectric liquids

## Submerged discharges



P. Bruggeman *et al.* 2008 *J. Phys. D: Appl. Phys.* **41** 194007  
See also F. Krčma's works

## Discharges in bubbles in liquids



Nomura *et al.* 2006 *Appl. Phys. Lett.* **88** 211503  
See also A. Hamdan's works

# 2.1 – Spark discharge in dielectric liquids

Model : SR15-R-1200

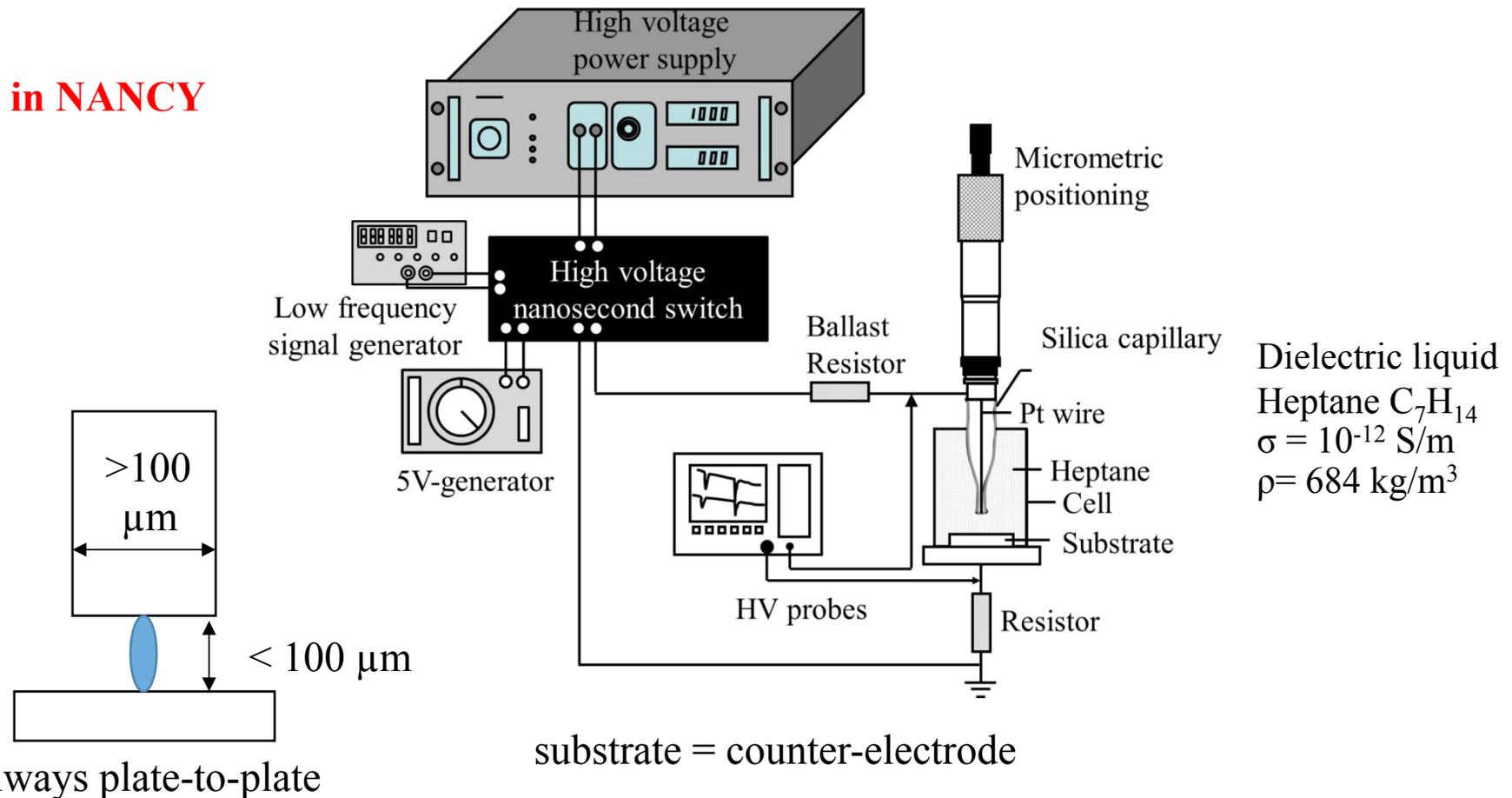
Continuous mode : 0 - 15 kV ; 0 – 80 mA

Pulsed by a Low Frequency Generator

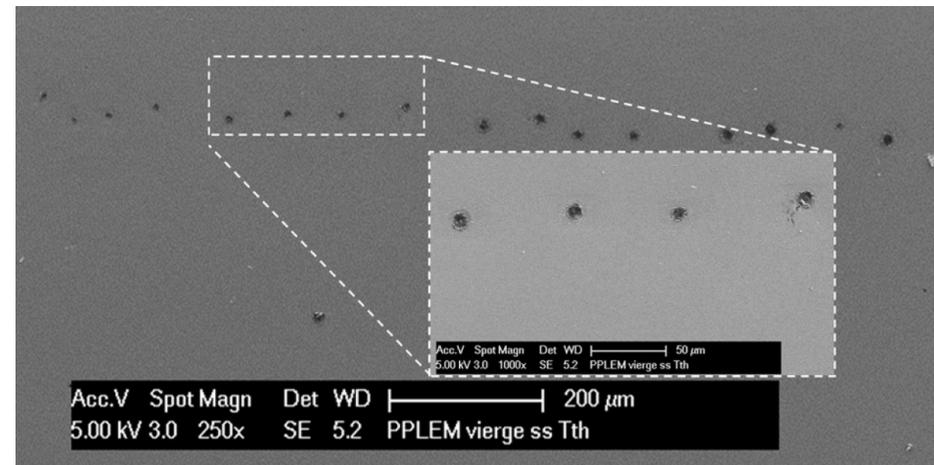
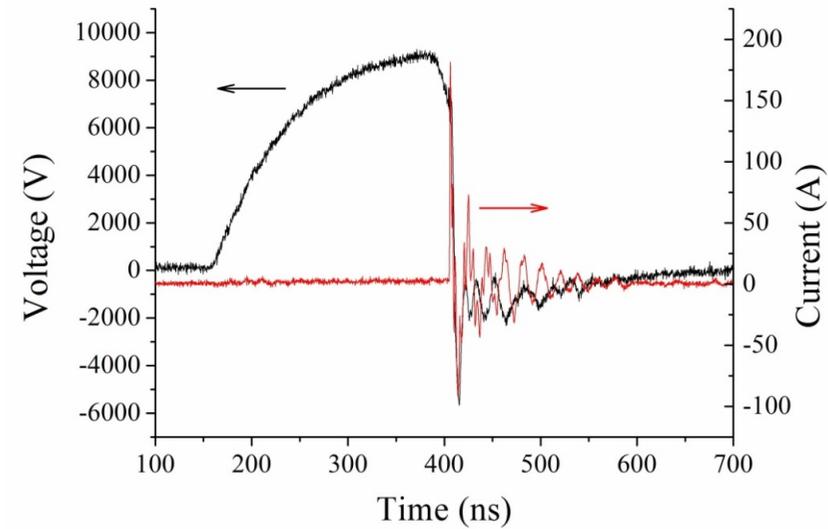
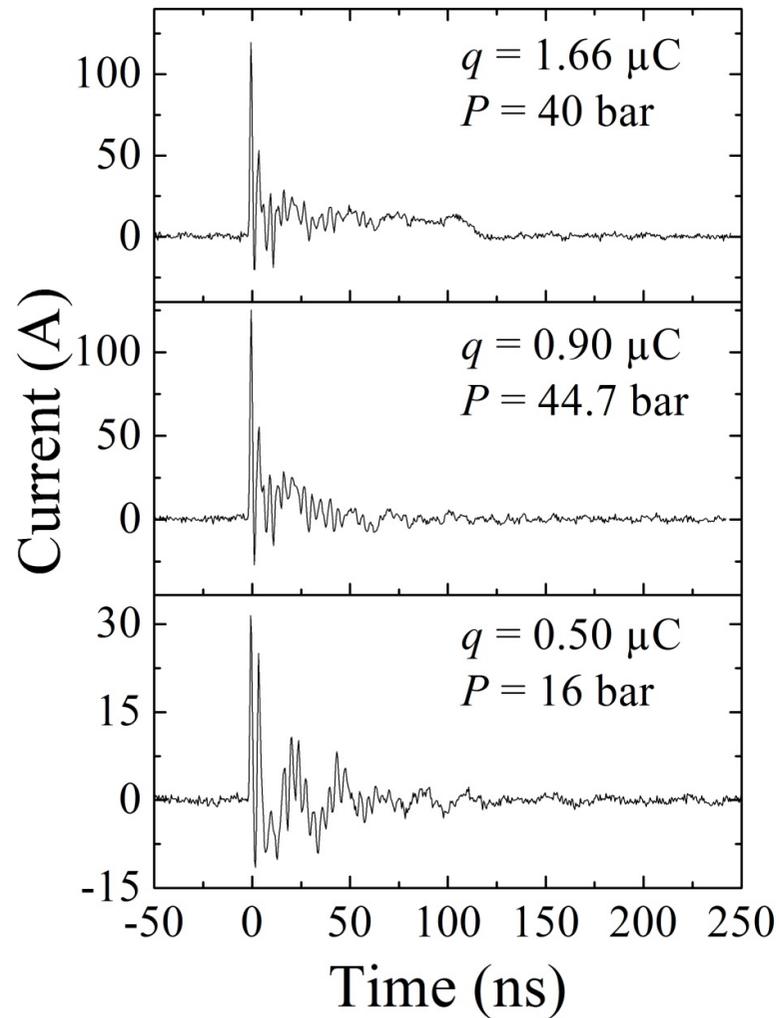
Micrometric positioning

screw ( $\Delta x = 10\mu\text{m}$ )

**Our system in NANCY**

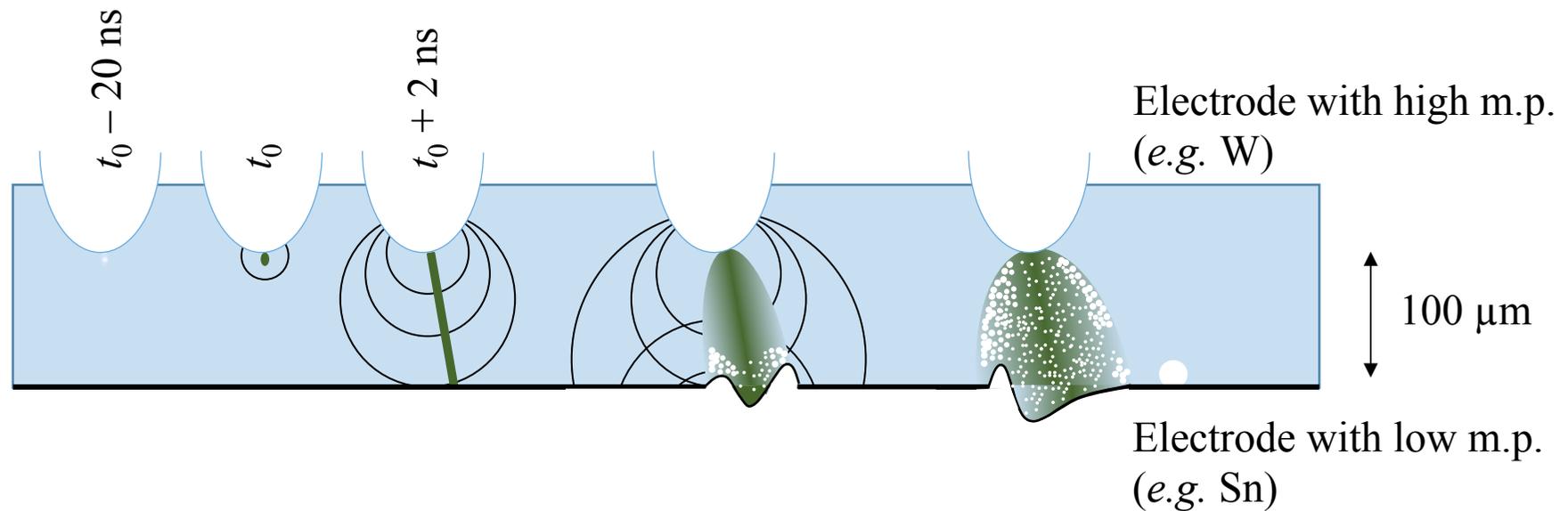


# 2.1 – Spark discharge in dielectric liquids



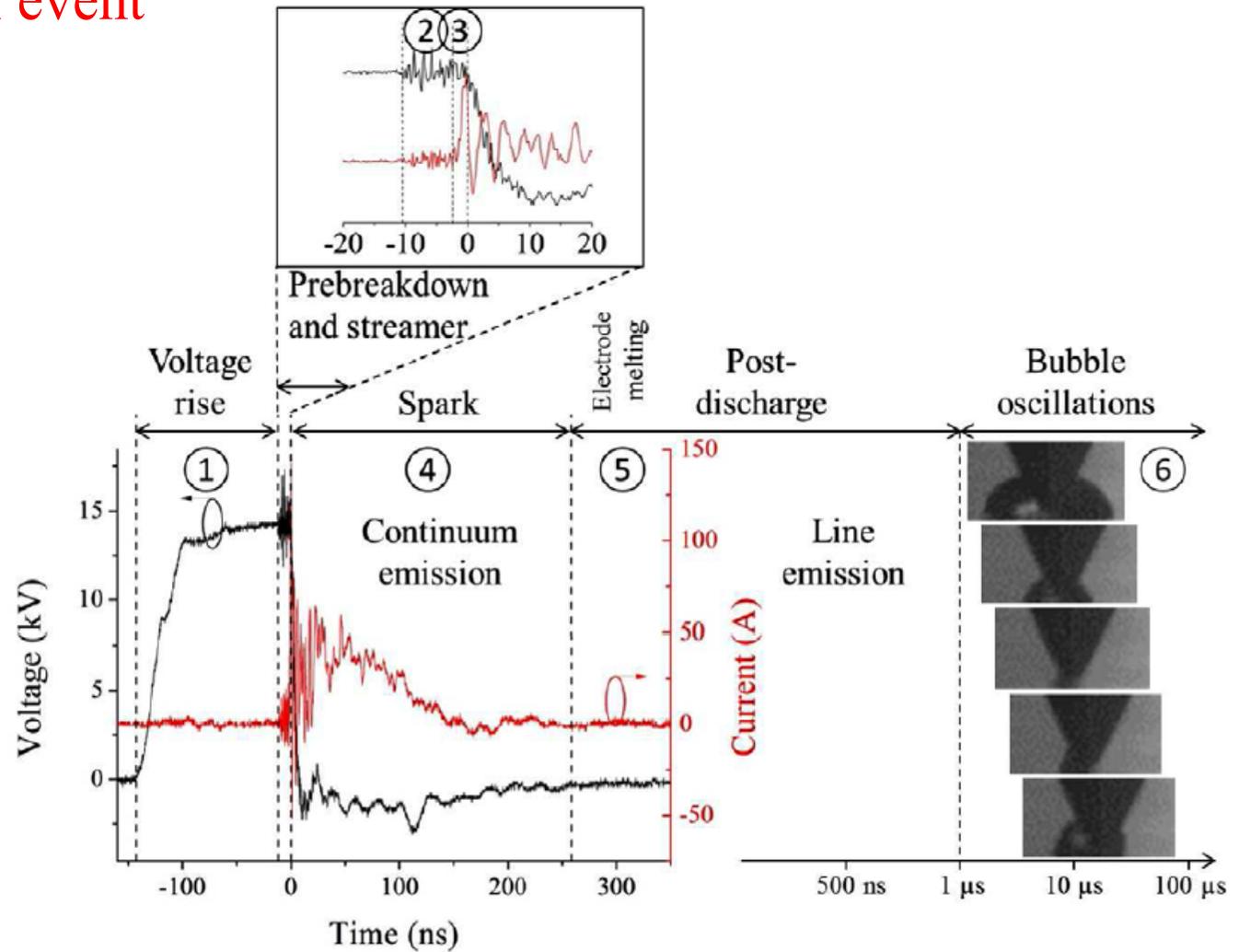
## 2.2 – Streamers

A simplified sketch of event



# 2.2 – Streamers

A simplified sketch of event



# 2.2 – Streamers

## Charge injection

- For very weak fields, the regime is purely ohmic.

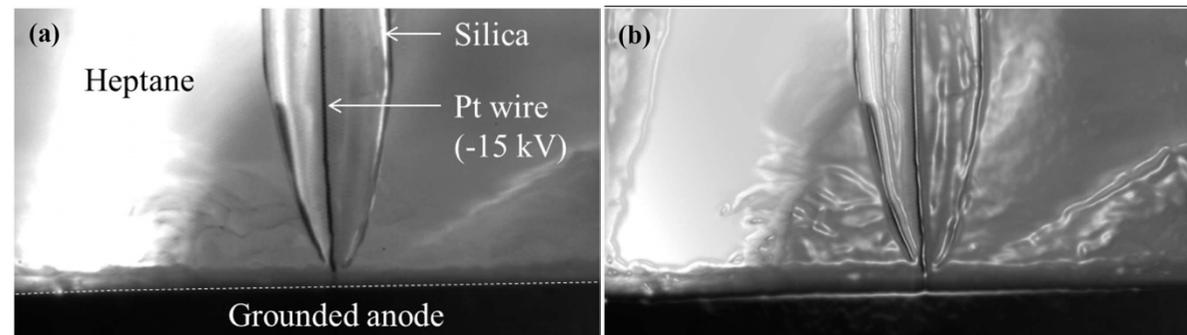
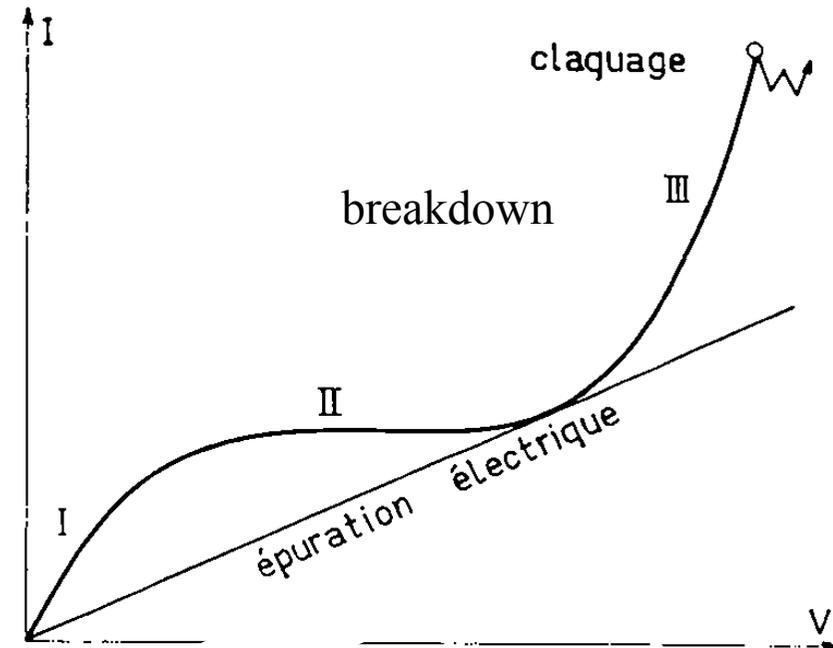
$$I = \sigma \cdot E$$

The electric field extracts impurity ions (limited in quantity) from the liquid.

- For intermediate fields, the current saturates. Sometimes, in the  $I-V$  curve, this saturation regime is reduced to an inflection point or can even be missing.

- When the electric field exceeds a certain threshold, the current rises until breakdown. In this regime, injection of charge carriers occurs via different mechanisms.

Massive injection of charge carriers in the liquid from one electrode induces a space charge gradient which causes turbulence.



## 2.2 – Streamers

### Charge injection

Injection\* mechanisms\*\* are many and as much dependent on the electric field (which is typically of the order of several tens to several hundreds of  $\text{MV}\cdot\text{m}^{-1}$ ) as they are of the liquid/electrode interfaces. One finds the following:

- Fowler–Nordheim (field emission) tunnelling from bulk metallic electrodes (such as platinum or gold),
- tunnelling through electrodes with semiconductive surfaces (a Schottky barrier such as iron or copper oxides),
- Poole–Frenkel emission by electrodes with insulating surfaces (such as alumina or silica),
- Auger electron emission (due to the relaxation of a hole created by the emission of tunnelling electrons), etc.

--

\* injection of electrons, *i.e.* for negative HV at the anode and grounded cathode. For positive HV at the cathode and grounded anode, the story is different (not told here).

\*\* depend on ion distribution in the Helmholtz double layer in electrochemistry. This double layer lowers the liquid–electrode potential barrier, making electronic exchanges easier.

## 2.2 – Streamers

### Breakdown

A matter of time scale!

1° "long" (microsecond) electric pulse:

As a result of electrostatic repulsion, the formation of low density channels occurs. Consequently, the discharge propagates through the low-density regions.

2° "Intermediate" (nanosecond) electric pulse:

The electrostatic forces support the expansion of nanoscale voids behind the front of the ionization wave. In the wave front, the extreme electric field provides a strong negative pressure in the dielectric fluid due to the presence of electrostriction forces, forming the initial micro-voids in the continuous medium.

3° "short" (picosecond) electric pulse:

The regions of reduced density cannot form because of the extremely short duration of the applied electric pulse. Ionization in the liquid phase occurs as a result of direct electron impact without undergoing a phase transition.

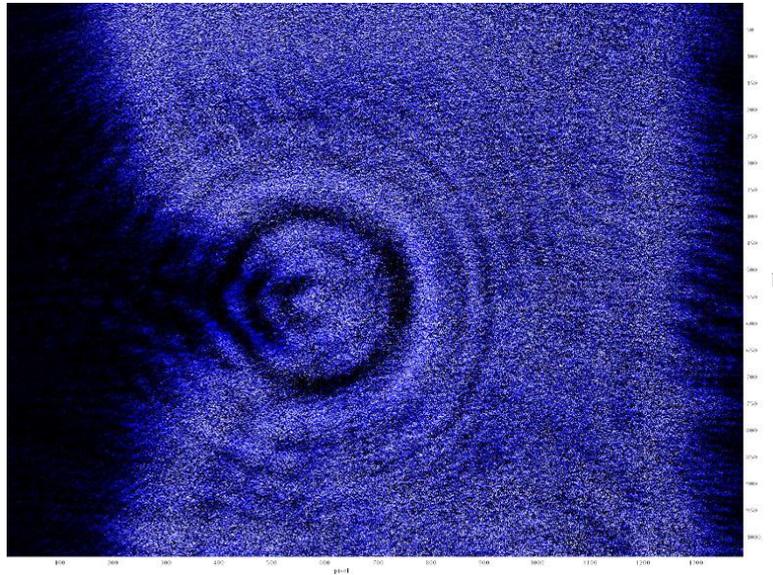
## 2.2 – Streamers

Breakdown

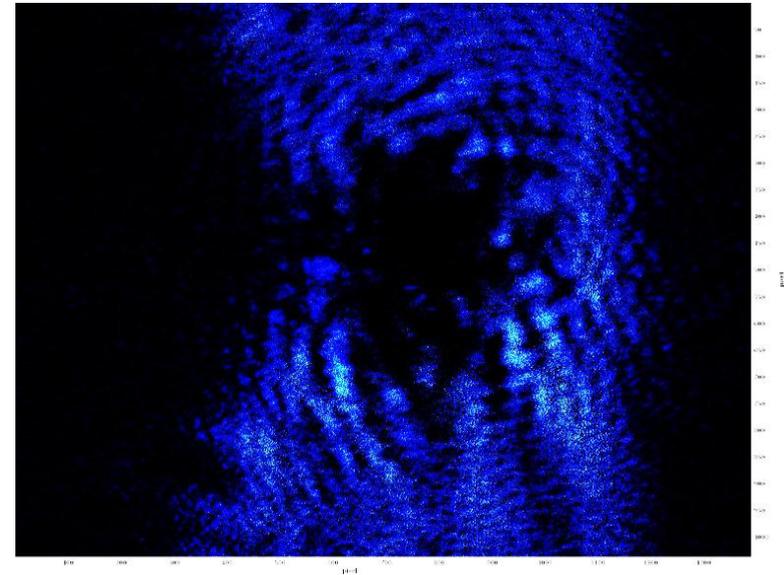
$$\vec{F} = e \delta n \vec{E} - \frac{\epsilon_0}{2} E^2 \nabla \epsilon + \frac{\epsilon_0}{2} \nabla \left( E^2 \frac{\partial \epsilon}{\partial \rho} \rho \right)$$

*Free charges*

*Electrostriction*



The electrostriction rarefaction wave without breakdown. Water, U = 9 kV, 20 ns. The image was collected 200 ns after the HV pulse.

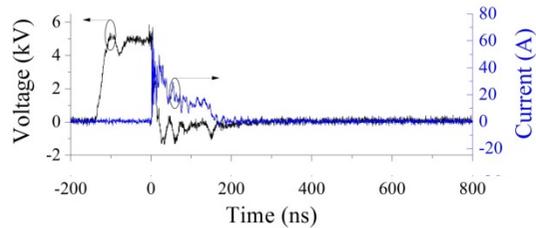
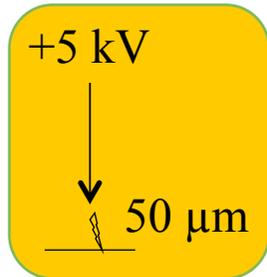


The streamer-leader structure formation. Water, U = 12 kV, 20 ns. The image was collected 200 ns after the HV pulse.

A. Starikovskiy, 32nd ICPIG, July 26-31, 2015, Iași, Romania

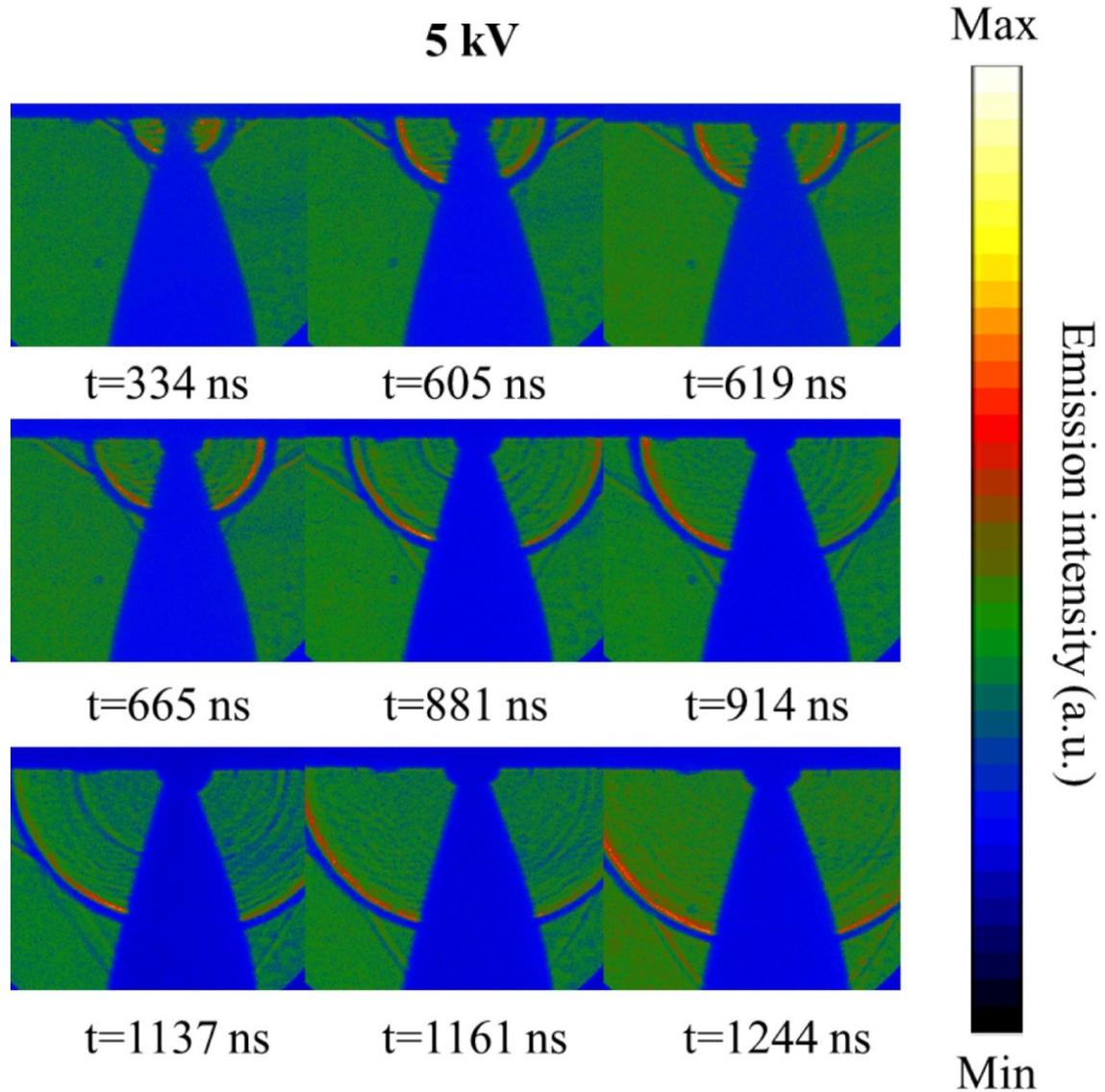
# 2.2 – Streamers

Emission of shock waves

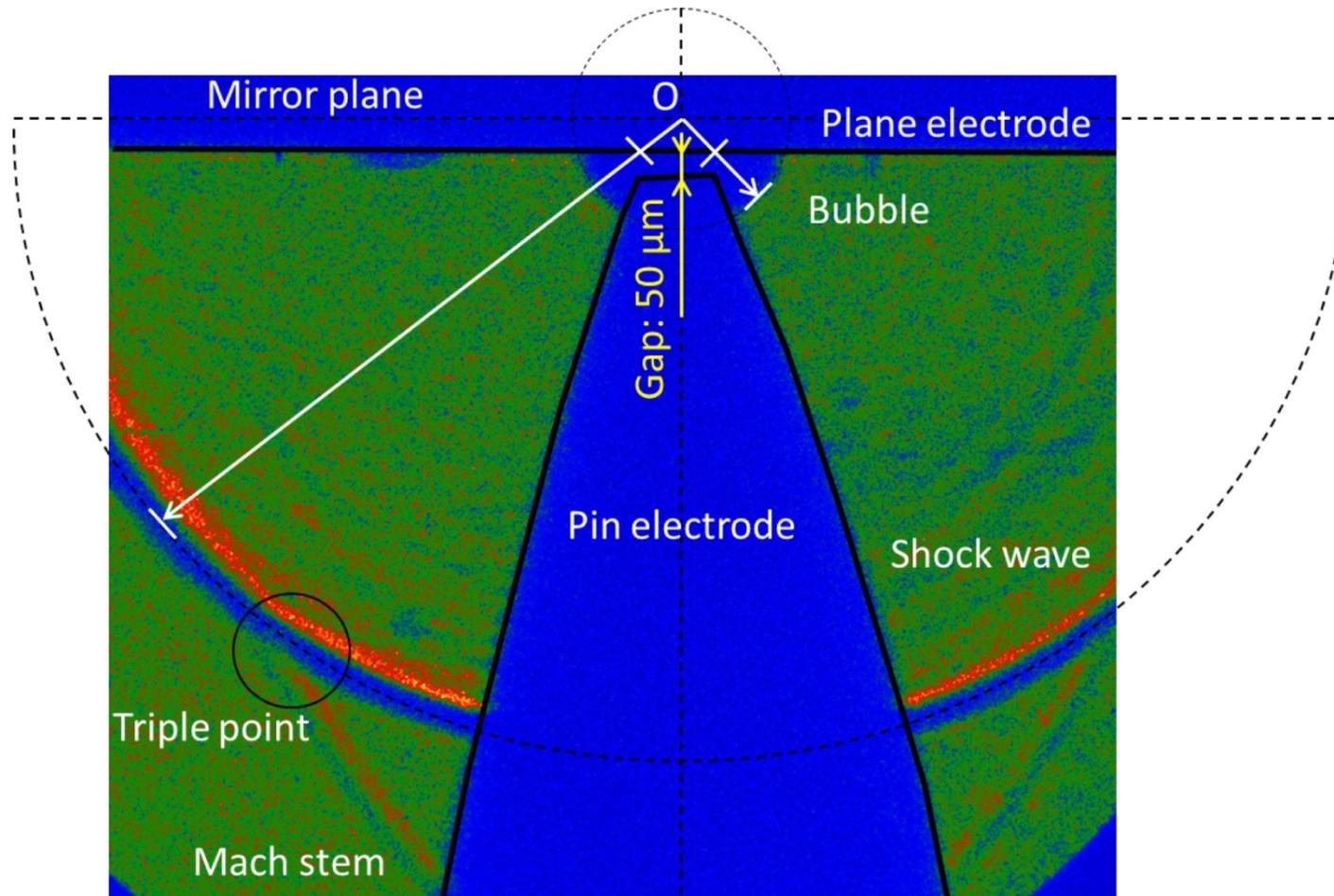


Bubble expansion step

The afterglow lasts about 100 ns



## 2.2 – Streamers



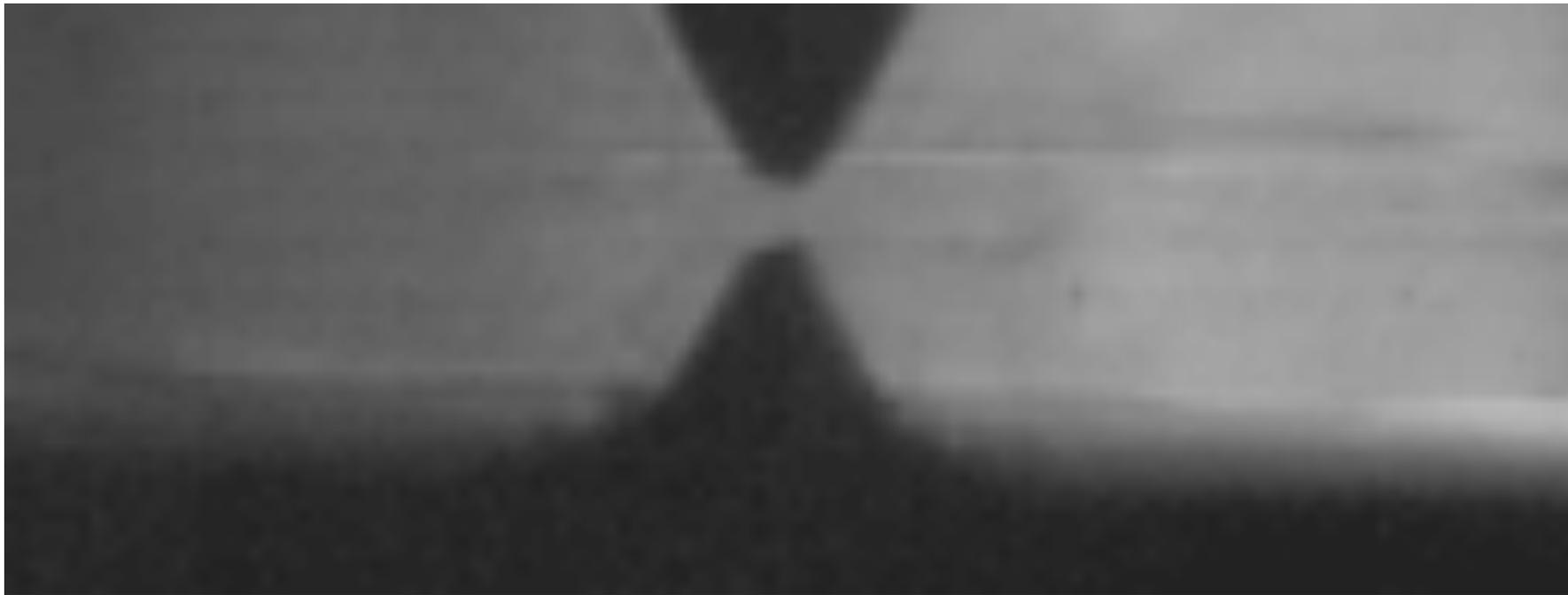
Hugoniot equations  $\rightarrow P = \rho_0 u_s \frac{u_s - A}{B}$

Liquid density
Acoustic sound velocity
Parameter

Usually a fraction of GPa

## 2.2 – Streamers

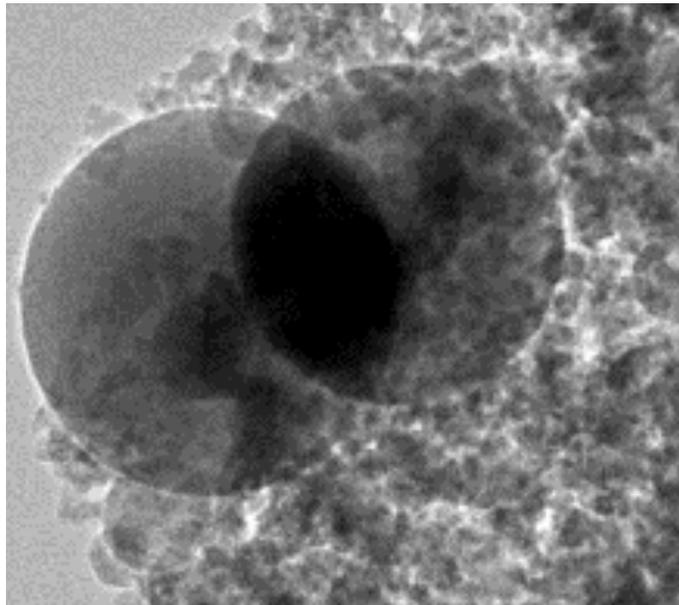
### Dynamics of bubbles



1.92  $\mu\text{s}$  / image (movie duration:  $\sim 500 \mu\text{s}$ )

# Chap. III

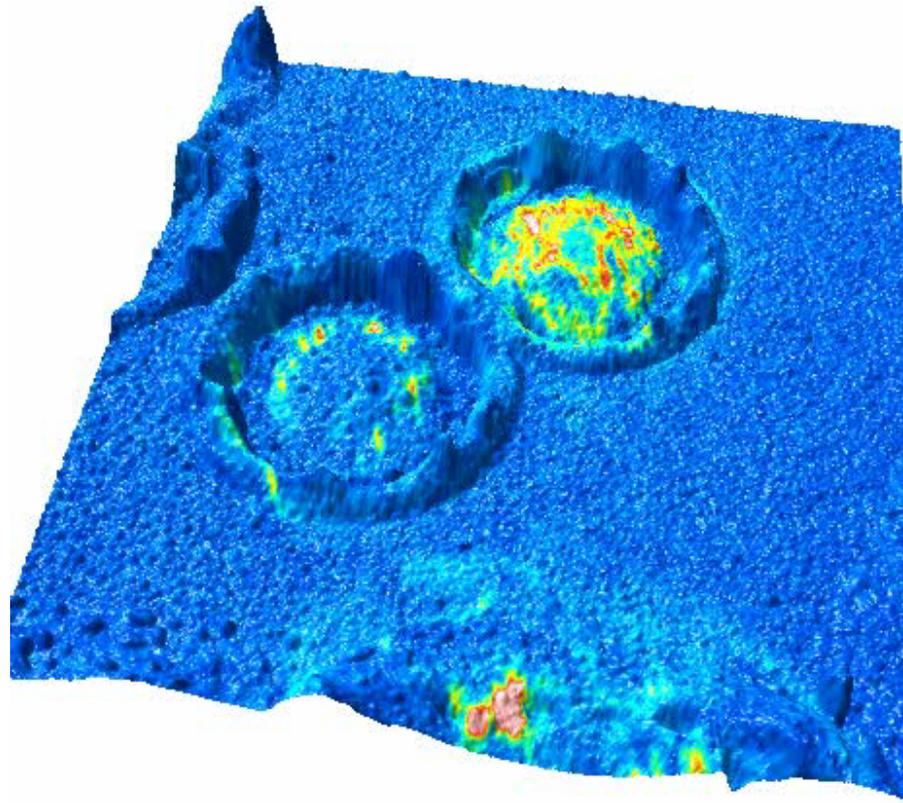
## Nanoparticles



# 3.1 – Electrode erosion

General mechanism: Melting

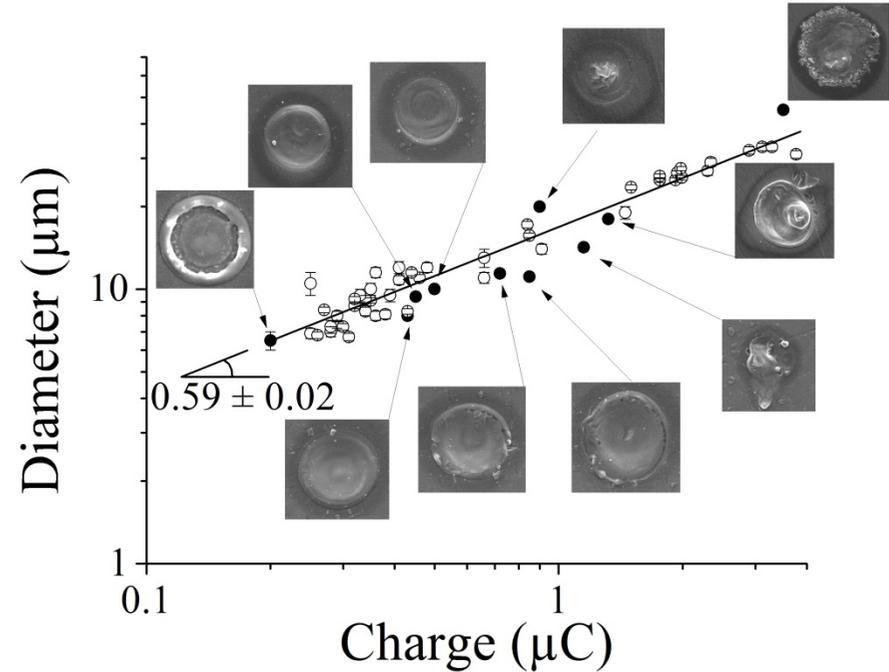
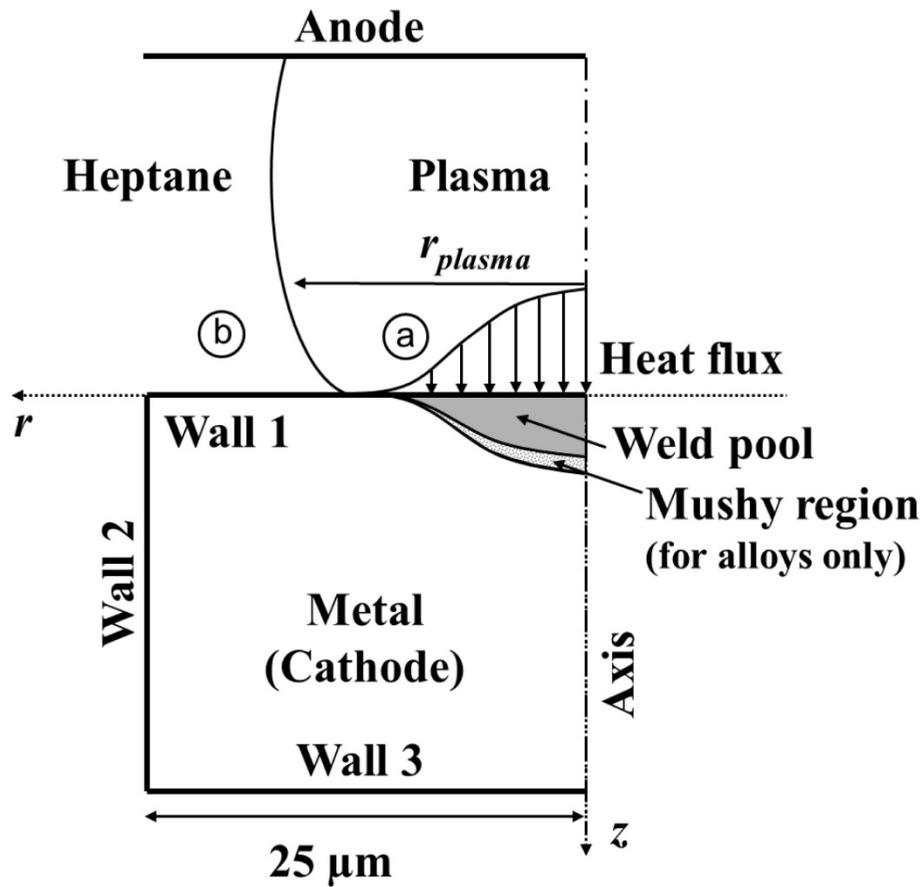
Shape of impacts



Nano-SIMS analyses of AlO mass peak projected onto the corresponding AFM image

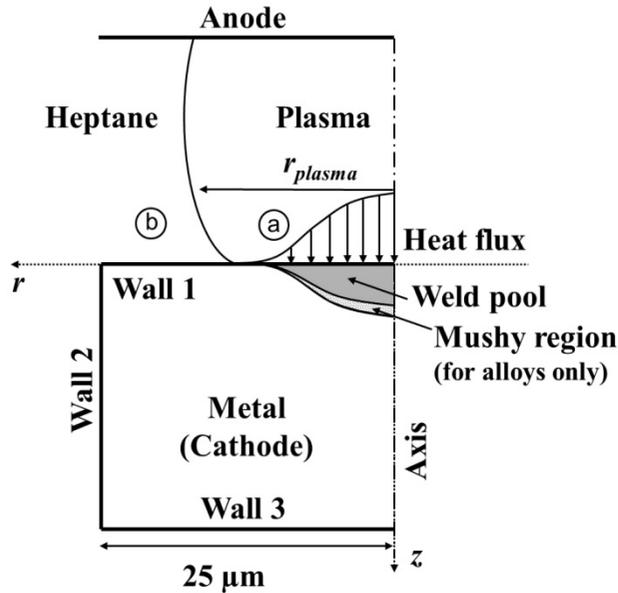
# 3.1 – Electrode erosion

Modelling plasma-surface interaction: example of Al



# 3.1 – Electrode erosion

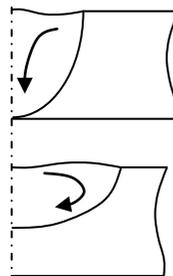
## Boundary conditions



Boundary	$u$	$v$	$T$
Wall 1a [<math>r_{plasma}</math>]	$(\tau_{\perp\perp}; \tau_{\perp\parallel})_{1a}$	0	$\phi_{1a} = -\alpha \frac{Q_{plasma}}{\pi r_{plasma}^2} \exp\left(-\frac{r^2}{r_{plasma}^2}\right) + h_{1a}(T - T_{plasma})$ $+ \sigma_e \varepsilon_{1a}(T^4 - T_{plasma}^4) + J^{(Al)} L_{SL}^{(Al)}$
Wall 1b [>math>r_{plasma}</math>]	$(\tau_{\perp\perp}; \tau_{\perp\parallel})_{1b}$	0	$\phi_{1b} = h_{1b}(T - T_{liquid}) + \sigma_e \varepsilon_{1b}(T^4 - T_{liquid}^4)$
Wall 2	0	0	$\phi_2 = 0$
Wall 3	0	0	$\phi_3 = 0$
Axis	$u=0$	$\frac{\partial v}{\partial r}\Big _{axis} = 0$	$\frac{\partial T}{\partial r}\Big _{axis} = 0$

$$\frac{\partial \sigma}{\partial T} = +10^{-4} N m^{-1} K^{-1} \quad \text{Marangoni positif}$$

$$\frac{\partial \sigma}{\partial T} = -10^{-4} N m^{-1} K^{-1} \quad \text{Marangoni négatif}$$



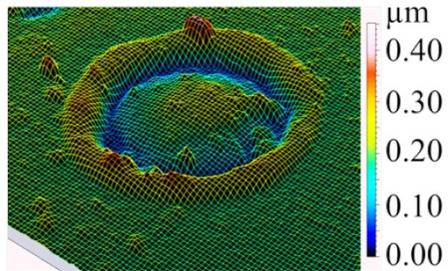
$$P_{surf} = P_{plasma} - \sigma \nabla \cdot \left( \frac{\vec{n}}{\|\vec{n}\|} \right)$$

$$\tau_{\perp\perp} = \mu \frac{\partial u}{\partial e_{\perp}} \approx K \gamma e_{\perp}$$

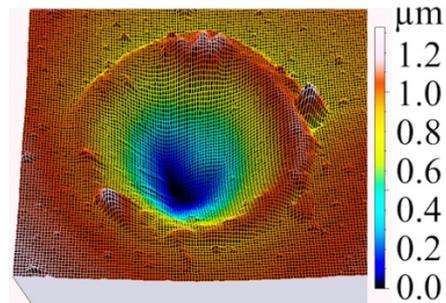
$$\tau_{\perp\parallel} = \mu \frac{\partial u}{\partial e_{\parallel}} = \frac{d\gamma}{dT} \frac{\partial T}{\partial e_{\parallel}}$$

# 3.1 – Electrode erosion

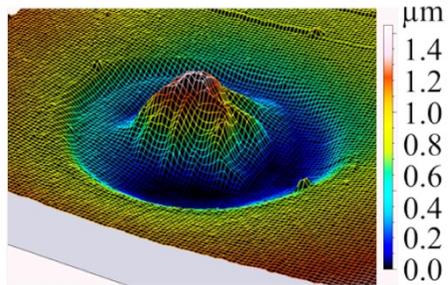
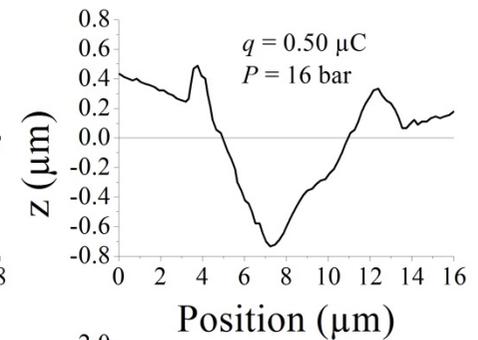
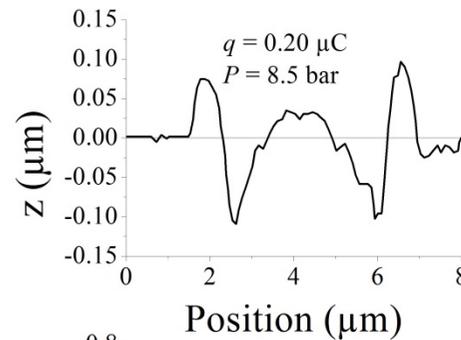
Selection of impacts  
Weakly asymmetric



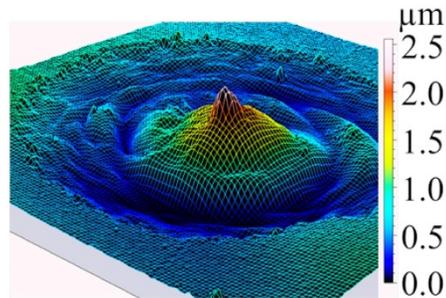
$q = 0.20 \mu\text{C} - P = 8.5 \text{ bar}$



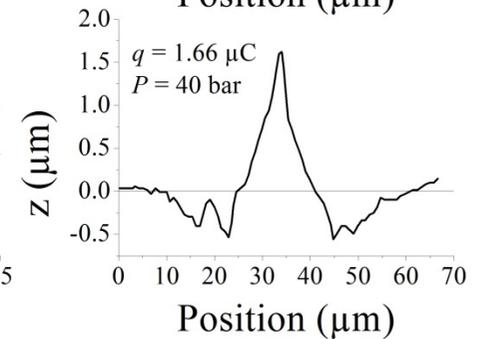
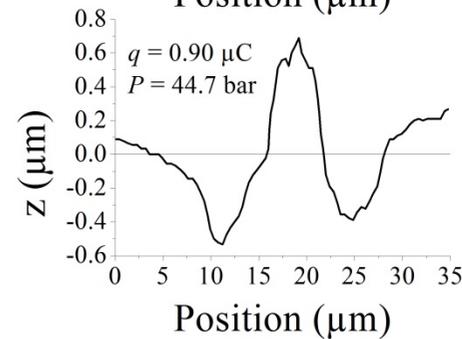
$q = 0.50 \mu\text{C} - P = 16 \text{ bar}$



$q = 0.90 \mu\text{C} - P = 44.7 \text{ bar}$

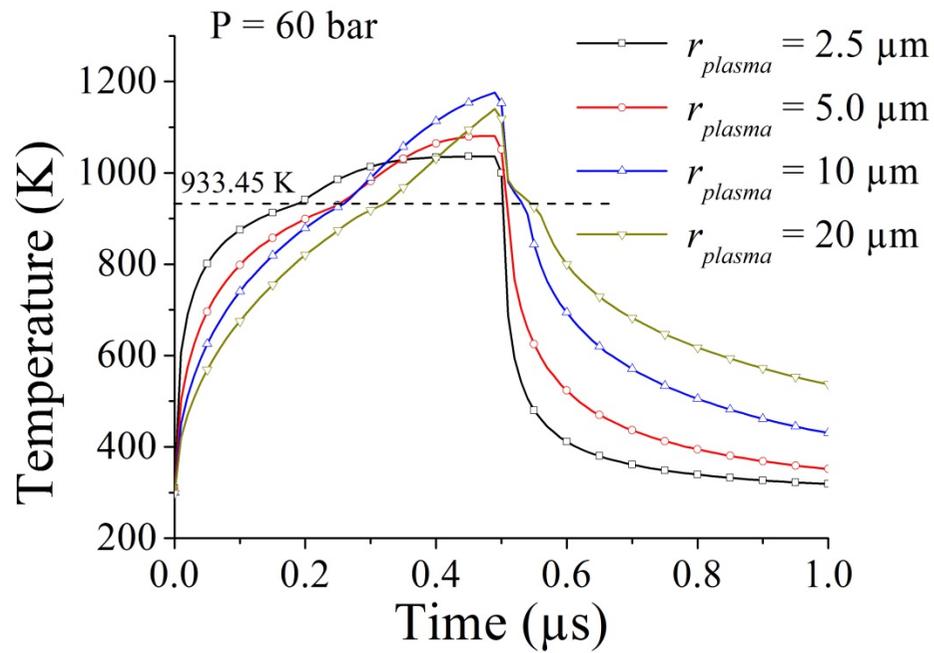


$q = 1.66 \mu\text{C} - P = 40 \text{ bar}$



$V_+/V_- = 1$  but practically  $V_+/V_- > 0.8$

# 3.1 – Electrode erosion



$Q=1.6 \text{ W/m}^2$   
 $R_p=2.5 \mu\text{m}$   
 $P=60 \text{ bar}$

# 3.1 – Electrode erosion

## Thomson-Marangoni's effect

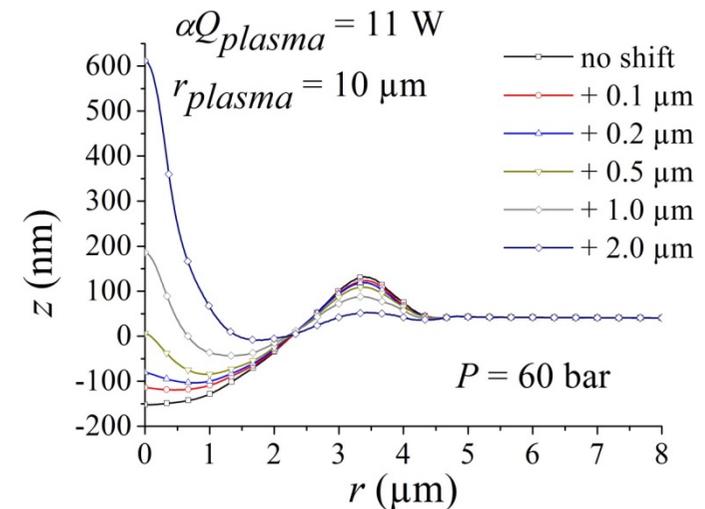
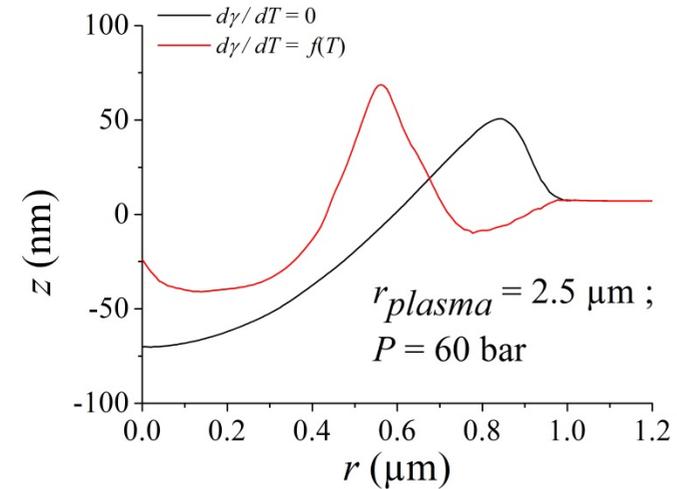
- Only negligible on large impacts ( $> 10 \mu\text{m}$ ) where pressure dominates
- Strong influence on  $d\gamma/dT$

## Pressure effect

- Dominant on large impacts ( $> 10 \mu\text{m}$ )
- Strong effect of the plasma orientation (tilt)

The plasma cannot be hollow ( $\varepsilon = 1.92$ )

*Babaeva et al., J. Phys. D: Appl. Phys. (2009) 42, 132003*



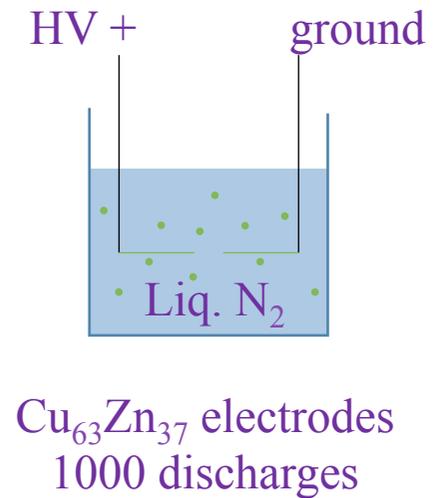
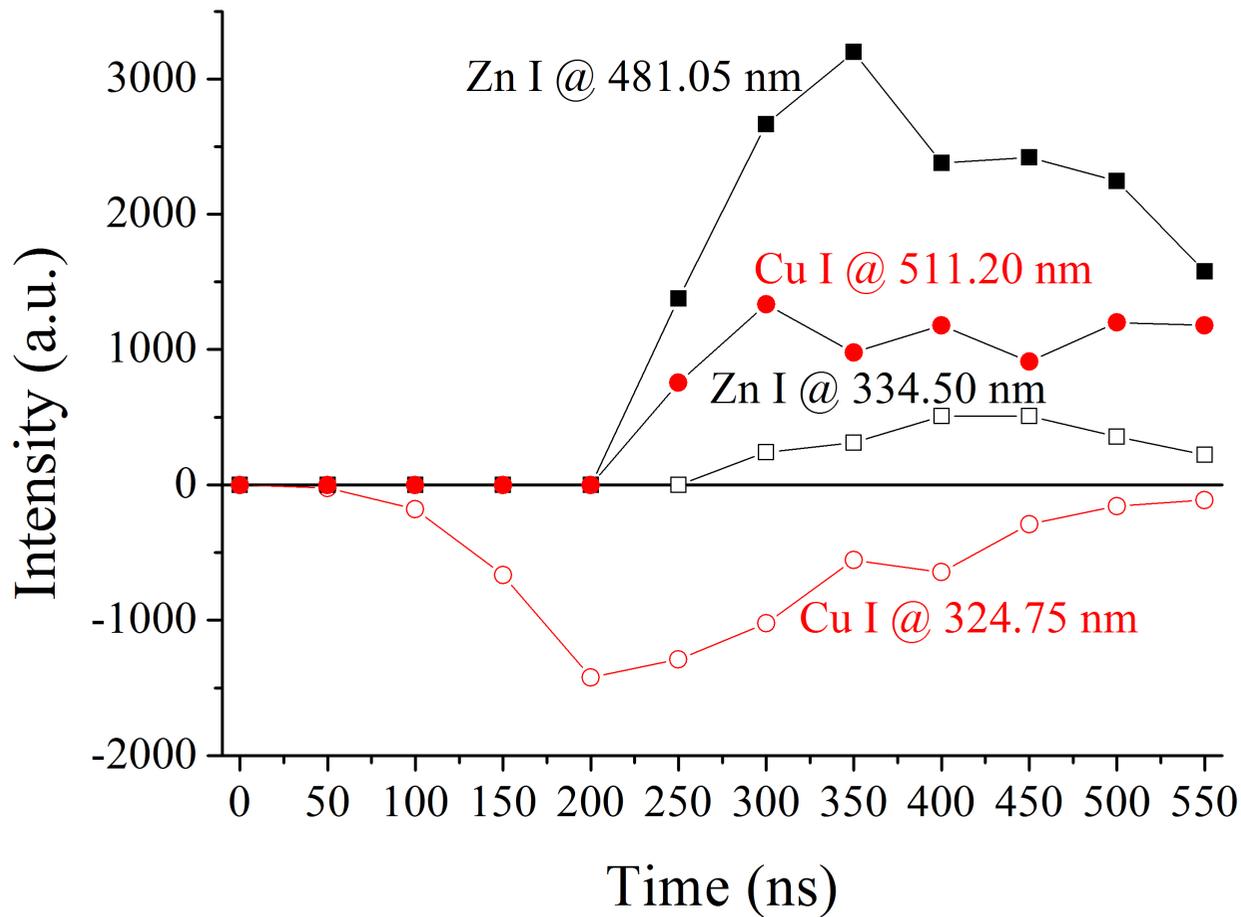
# 3.1 – Electrode erosion

Plasma radius $r_{plasma}$ [ $\mu\text{m}$ ]	Heat flux $\alpha Q_{plasma}$ [W]	Maximal temperature [K]	Impact radius [ $\mu\text{m}$ ]	Energy to create impact [ $\mu\text{J}$ ]	Total energy (exp. data) [mJ]
2.5	1.6	1036	0.96	0.80	<b>0.02</b>
5.0	3.8	1080	2.05	1.90	<b>0.50</b>
10	11	1176	4.67	5.50	<b>0.90</b>
20	35	1140	8.69	17.5	<b>1.66</b>

From 0.4 to 4% of the total energy is used in electrode erosion

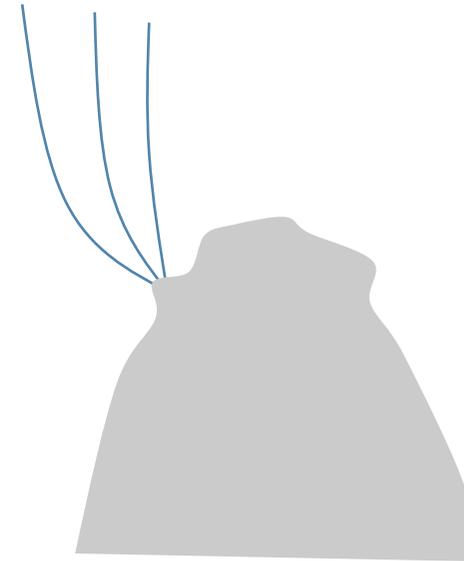
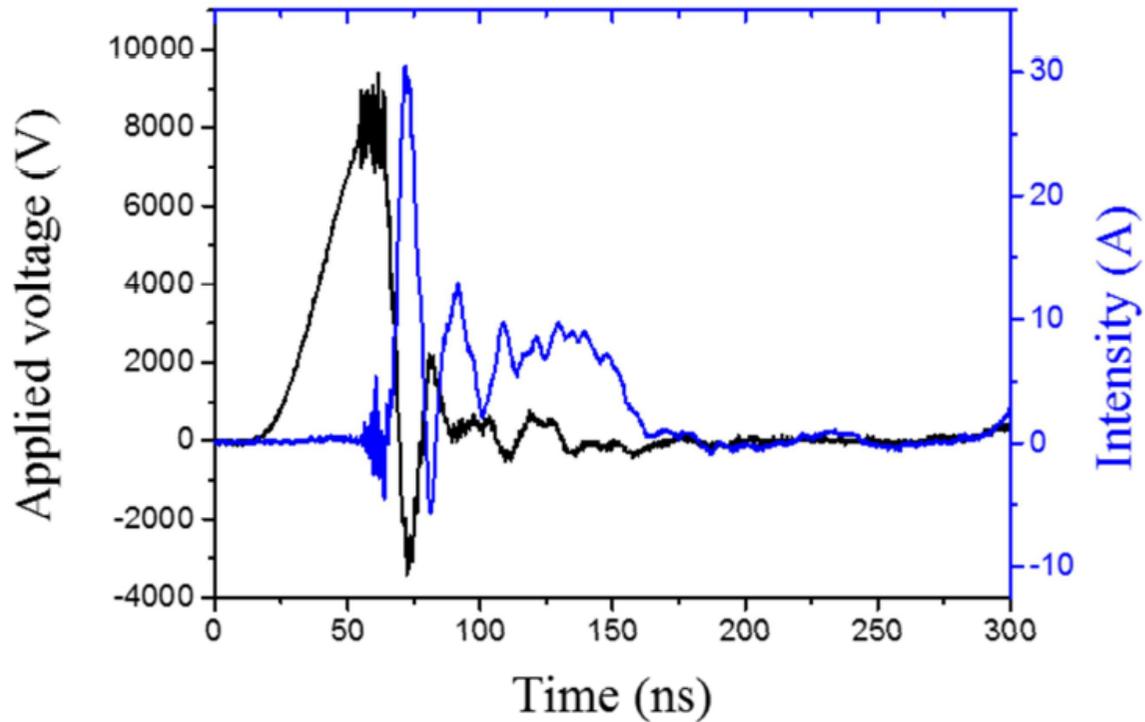
# 3.1 – Electrode erosion

Appearance of Cu I lines much before melting

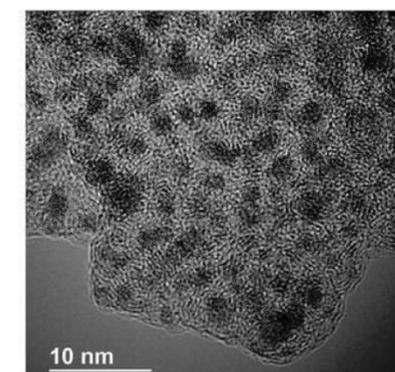
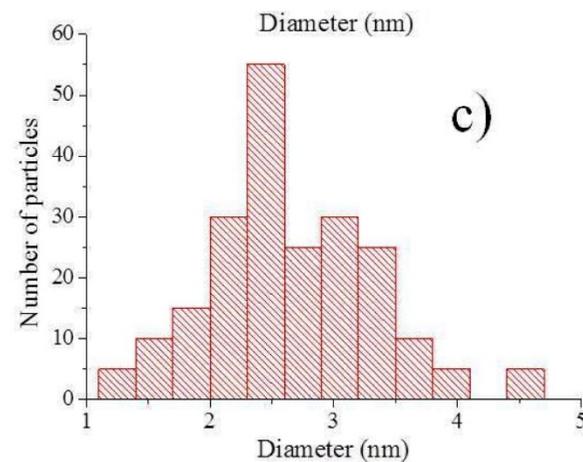
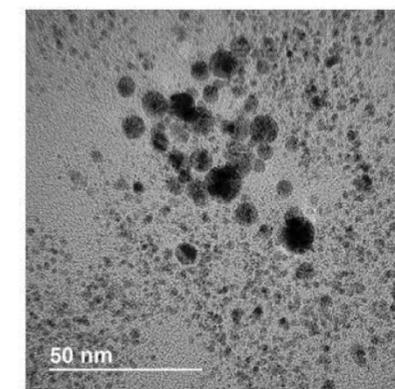
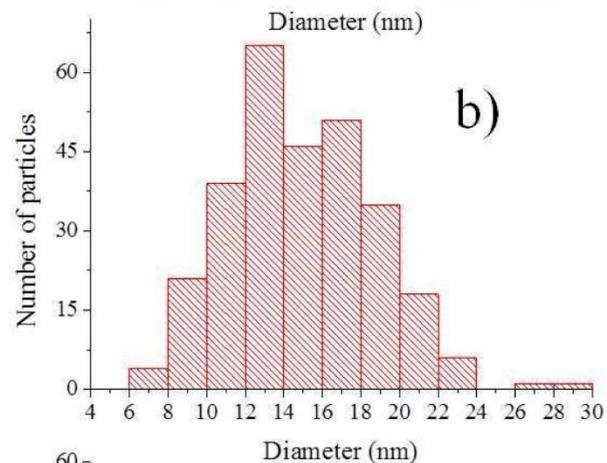
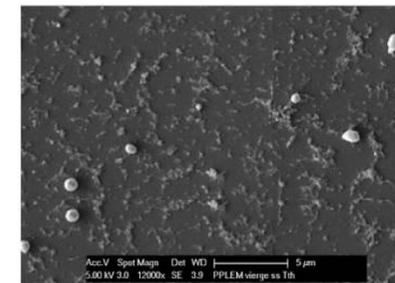
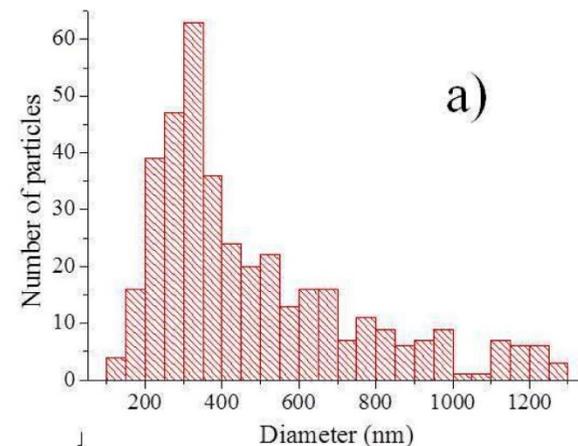


# 3.1 – Electrode erosion

Ohmic overheating when a high-current density is conducting via a protrusion



# 3.2 – Size distribution



# 3.3 – Nanoparticles

## Shape

### Spheres

Al, Cu, Ag, Pt, Au, Fe, W, Sn

### Sheets

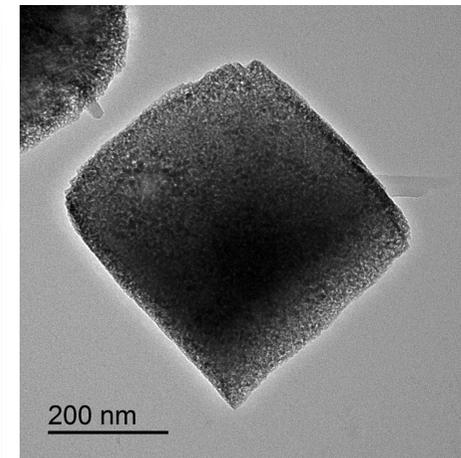
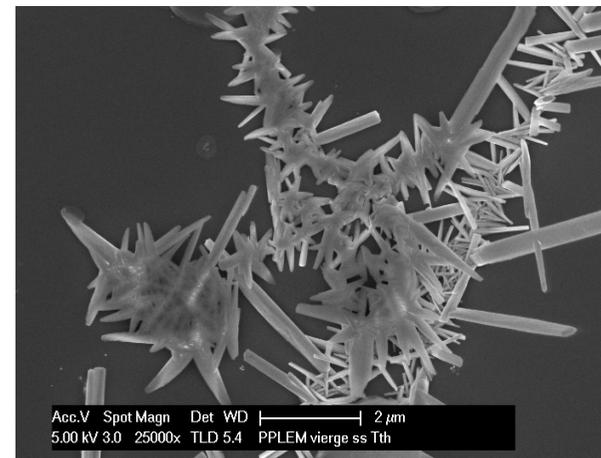
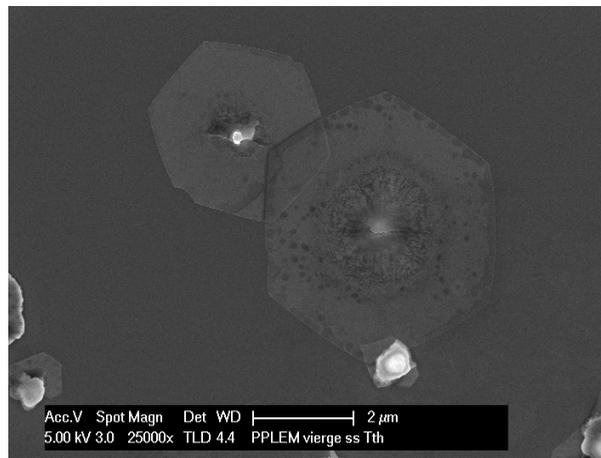
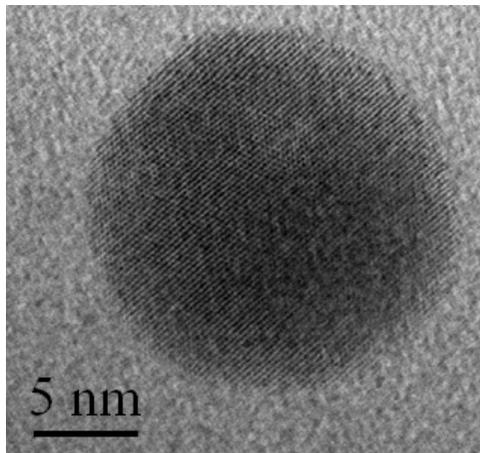
Ag, Bi, **Pb**, Zn

### Wires

Ag, Pb

### Cubes

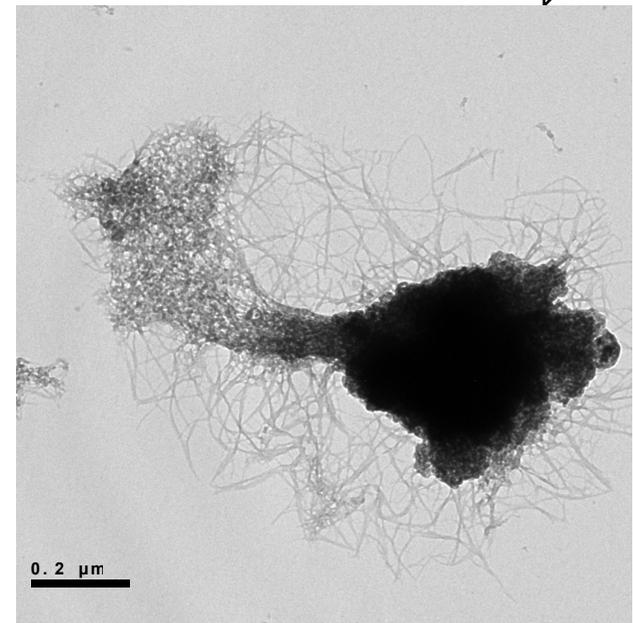
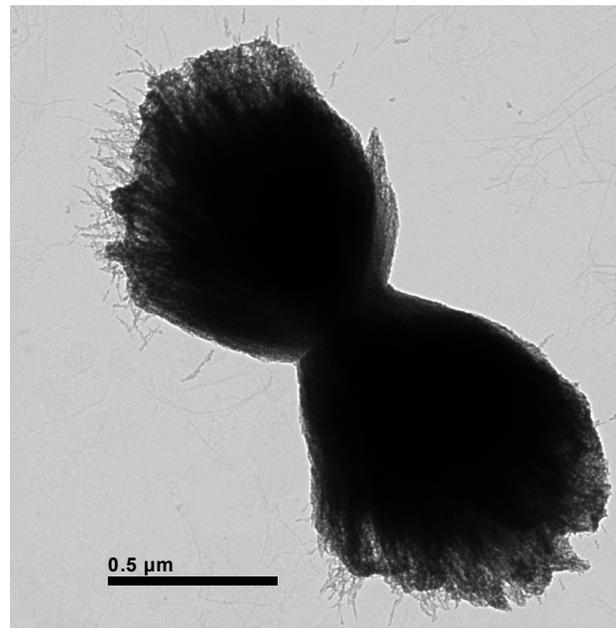
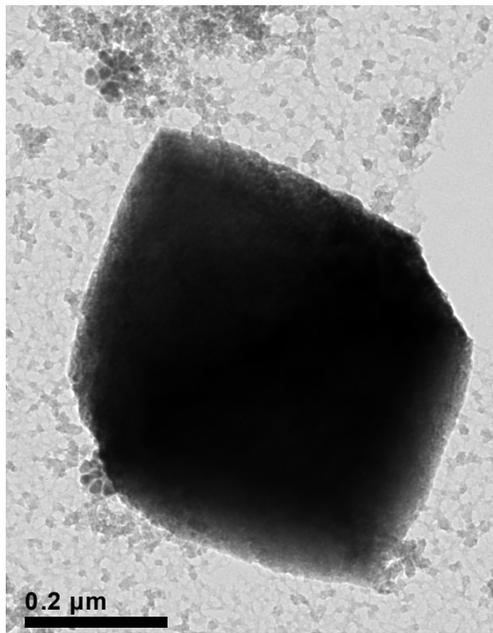
Cd, Ge



# 3.3 – Nanoparticles

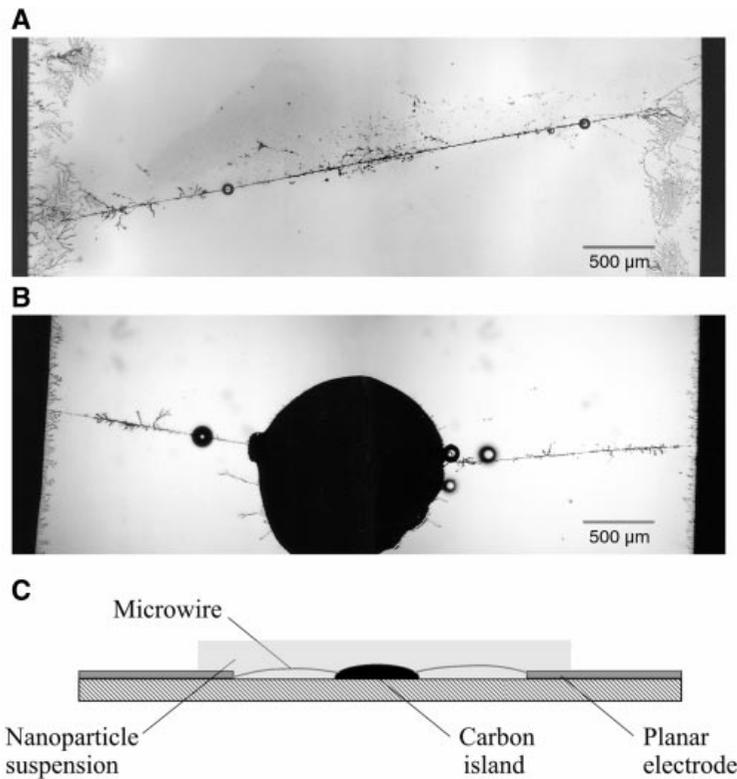
## Role of the electric field

Increasing applied voltage



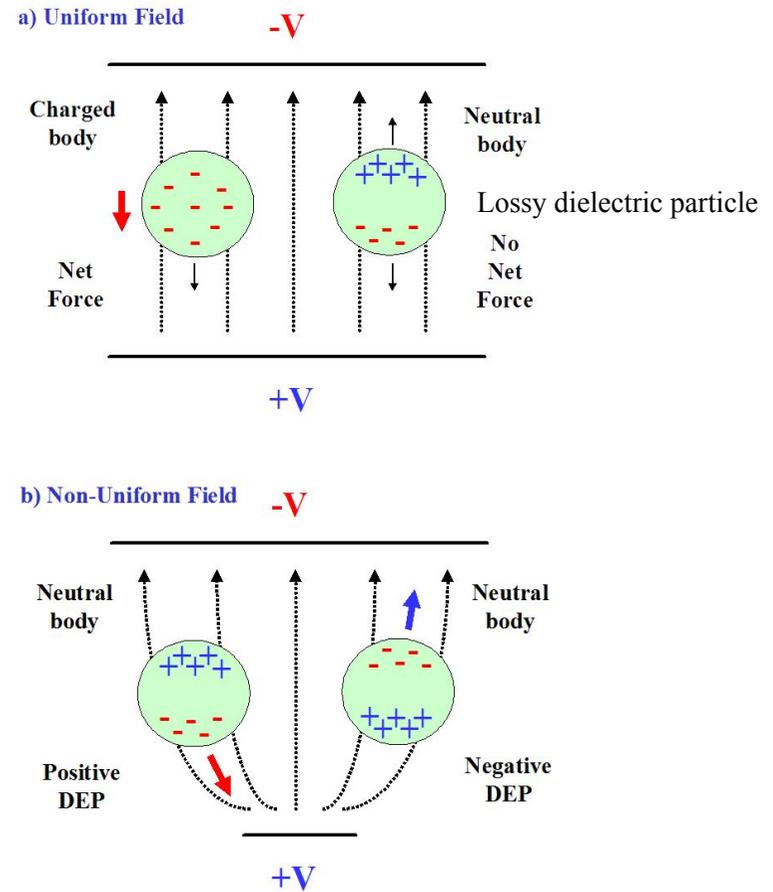
# 3.3 – Nanoparticles

## Role of the electric field



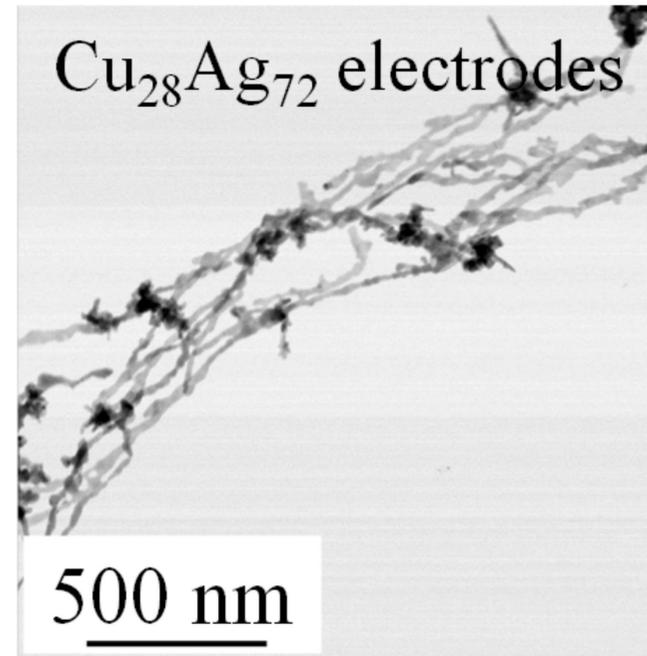
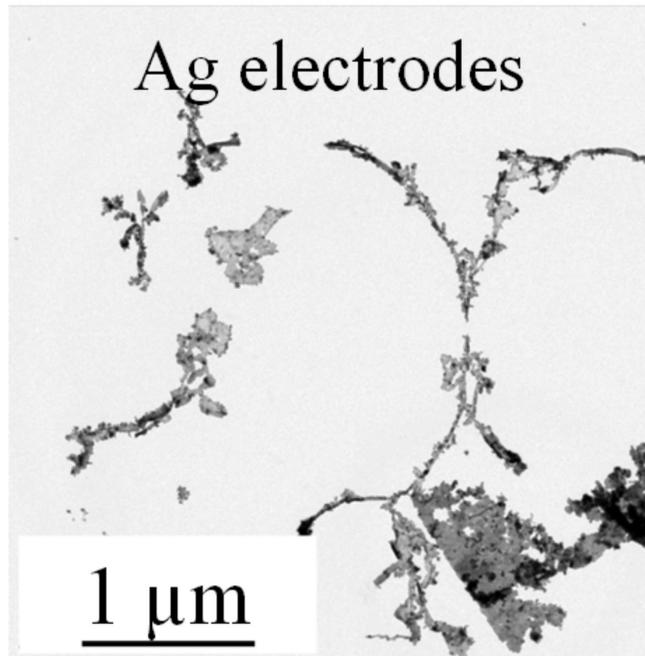
Hermanson *et al.* 2001 *Science* 294 1082

Dielectrophoresis might be used if insulating materials are synthesized



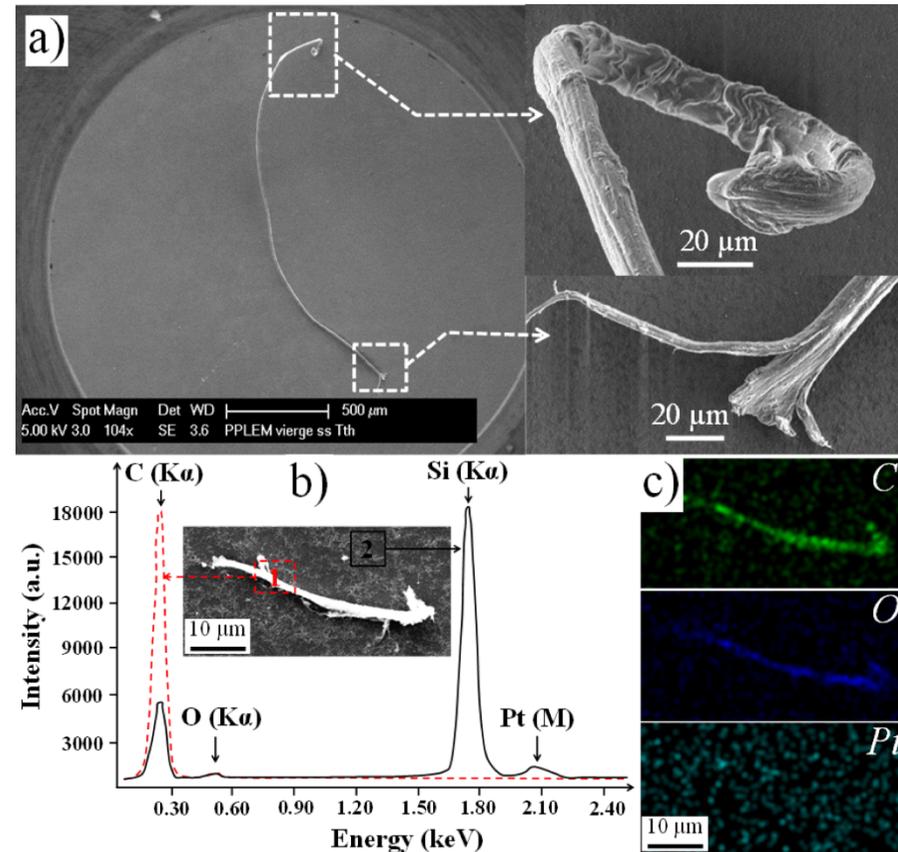
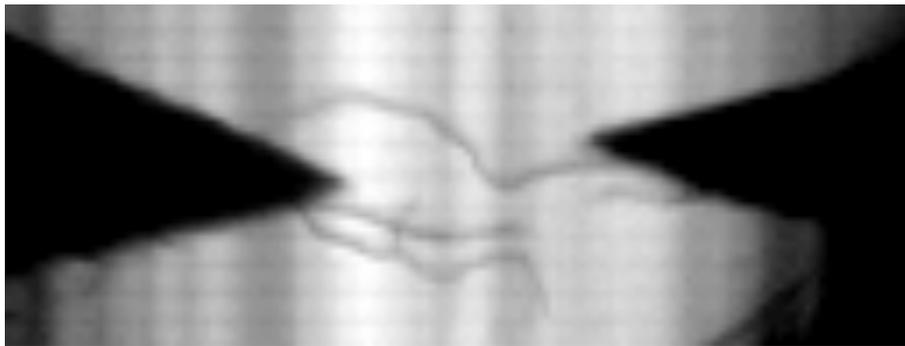
# 3.3 – Nanoparticles

## Role of the electric field



# 3.3 – Nanoparticles

## Role of the electric field

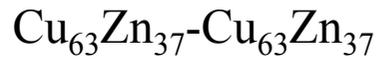


Hamdan et al. 2014 Mater. Lett. 135 115

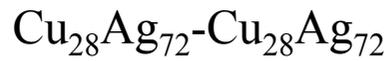
# 3.3 – Nanoparticles

$m.p. (Cu) = 1356 \text{ K}$   
 $m.p. (Ag) = 1234 \text{ K}$

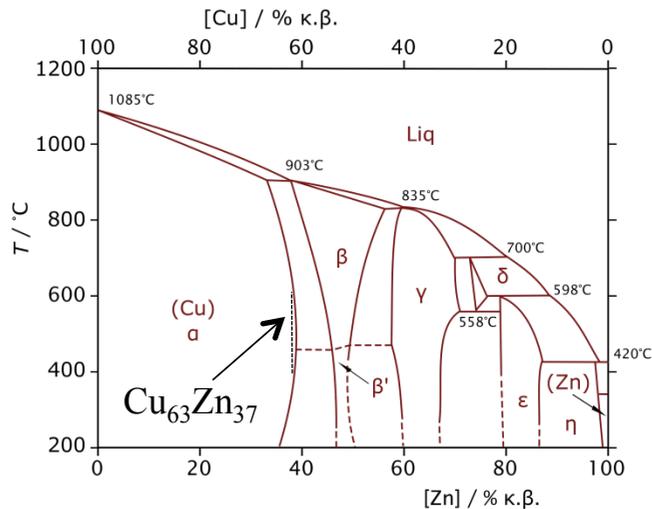
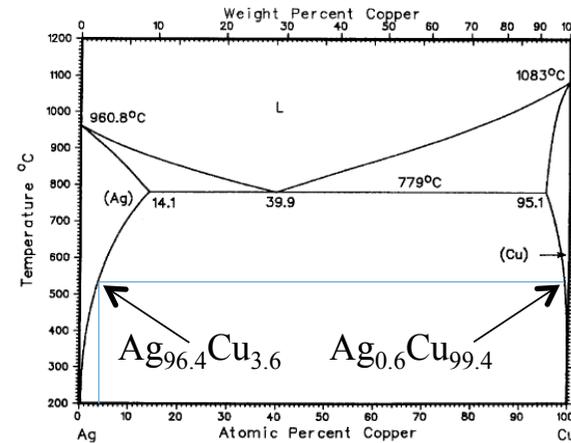
## Alloy nanoparticles



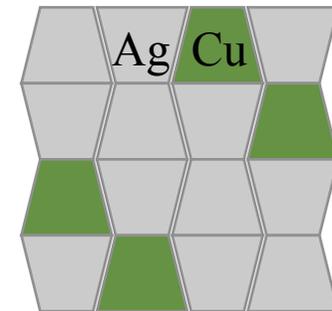
Cu-Zn



Cu-Ag



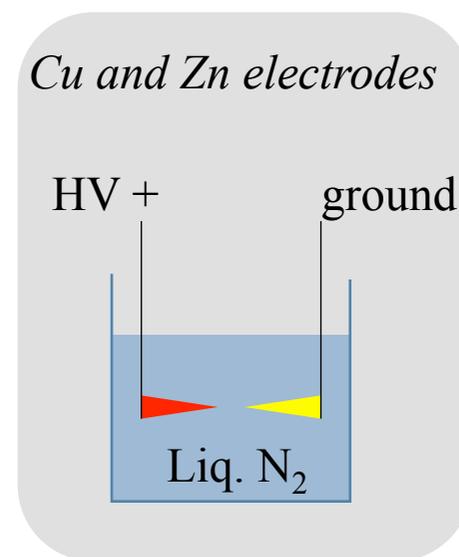
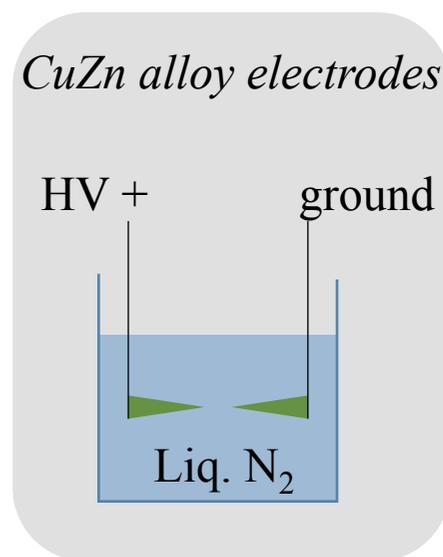
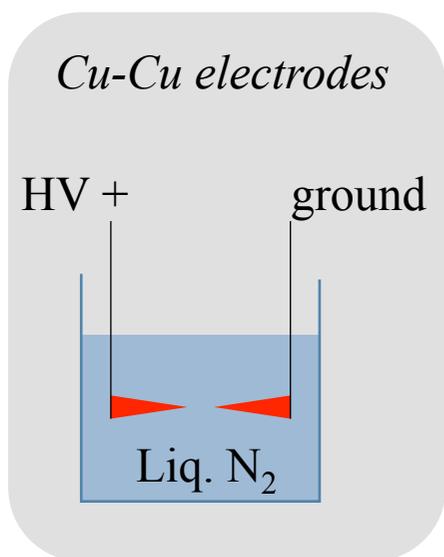
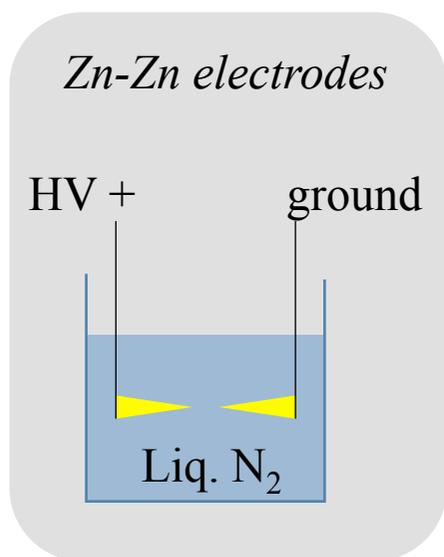
$Cu_{28}Ag_{72} \text{ (at.\%)} =$   
 $73\text{wt.\%} Ag_{96.4}Cu_{3.6} + 27\text{wt.}$   
 $\% Ag_{0.6}Cu_{99.4}$   
 as determined by XRD



$m.p. (Cu) = 1356 \text{ K}$   
 $m.p. (Zn) = 907 \text{ K}$

# 3.3 – Nanoparticles

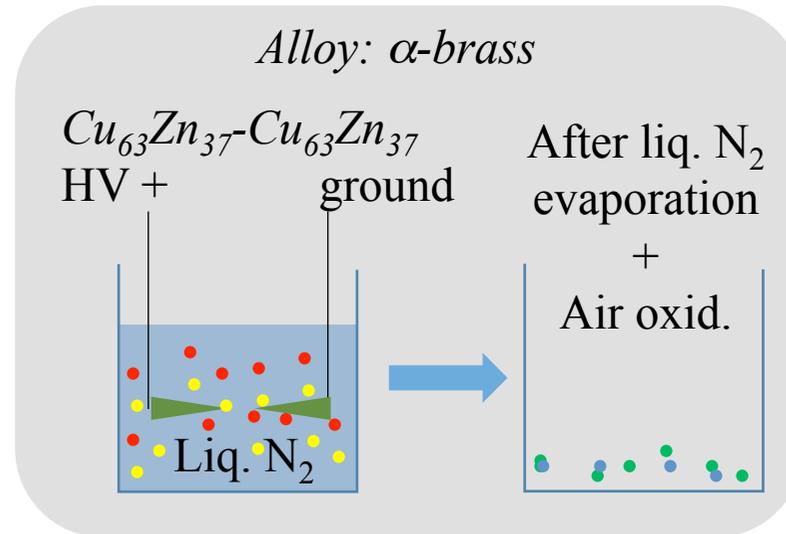
## Alloy nanoparticles



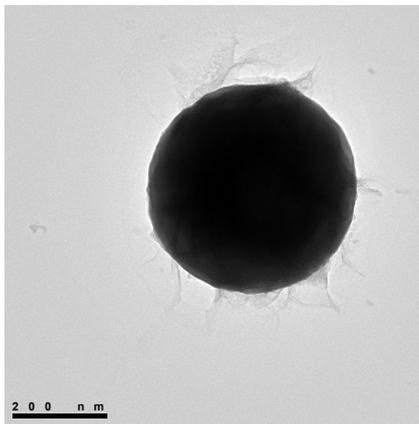
# 3.3 – Nanoparticles

## Alloy nanoparticles

Shift in the phase diagramme



$Cu_{63}Zn_{37}$  wires  
 $\varnothing$  1 mm

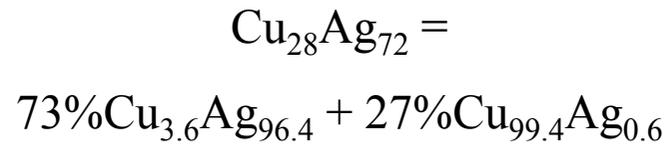


Particles are enriched in copper (dezincification)

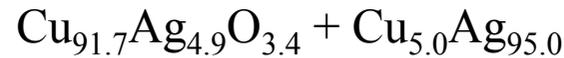
Cu (73 wt.%), Zn(25 wt.%) & O(2 wt.%)

# 3.3 – Nanoparticles

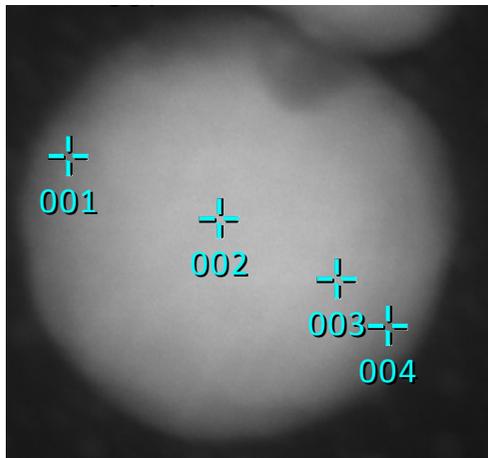
## Alloy nanoparticles



vs.



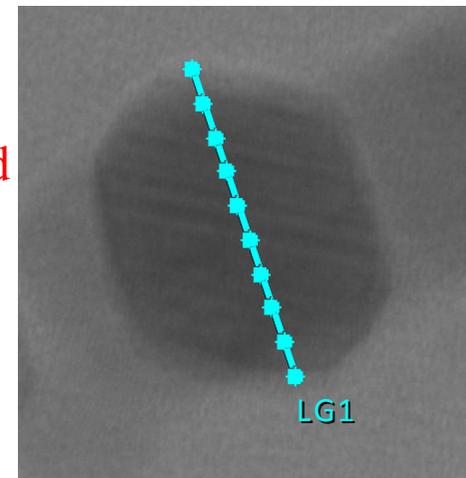
Memo	O	Cu	Ag
1	2.39	19.91	77.7
2	0	8.38	91.62
3	0	6.91	93.09
4	0	5.73	94.27
5	0	4.95	95.05
6	0	5.73	94.27
7	0	4.18	95.82
8	0	5	95
9	0	8.58	91.42
10	0	17.81	82.19



No significant shift in the phase diagramme

Only single-phase nanoparticles are obtained

Memo	O	Cu	Ag	Total(Atom %)
1	11.43	84.27	4.3	100
2	3.38	91.73	4.89	100
3	6.55	90.04	3.41	100
4	14.7	80.65	4.66	100



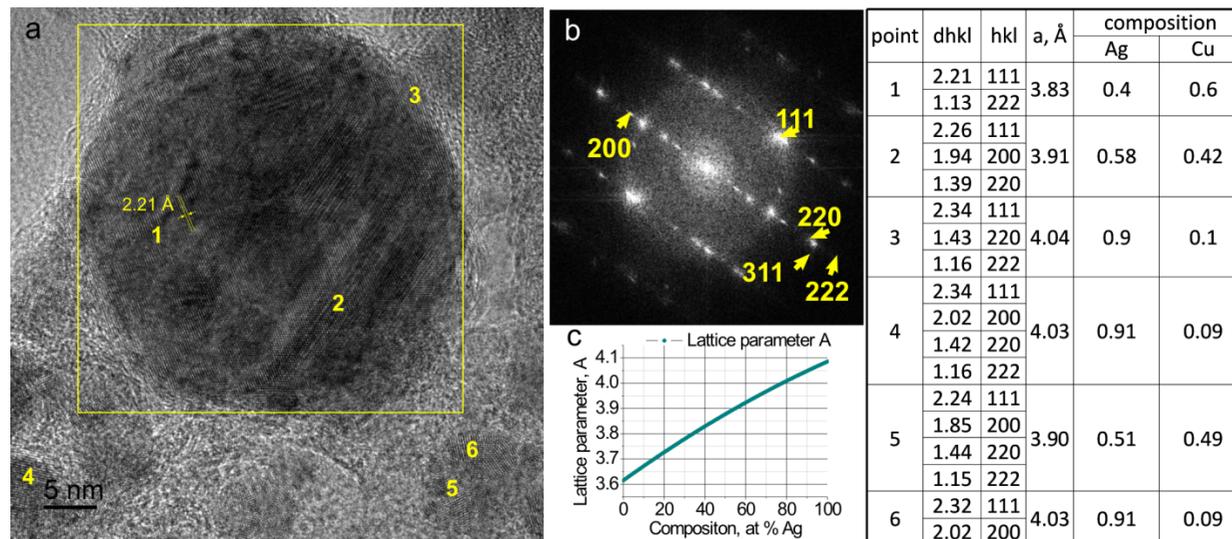
# 3.3 – Nanoparticles

No alloying from metallic vapours (no mixing?)

## Alloy nanoparticles

Shift in the phase diagramme for elements with different melting points

The HRTEM image of as-prepared alloyed nanoparticle (a) and results of FFT processing of the corresponding area (b). The interplanar spacings found correspond to the (111) plane of the cubic AgCu alloy of the composition  $Ag_{0.4}Cu_{0.6}$  as determined from the lattice parameter – composition dependence (c).



A  $Ag_{0.4}Cu_{0.6}$  NP cannot be predicted by the phase diagram

# 4 – Conclusion

Safe process for the user

Specificity due to pressure gradients

Several erosions processes at stake, early processes having to be clarified

Producing alloy nanoparticles is possible directly from alloy electrodes with pre-determined composition to account for element removals at high temperature but the final composition usually depends on the size of the produced NP.

The control of compositions is still to be studied, but likely possible. Anyway, getting very small NPs is a real challenge.

Phase stability seems to be surprisingly high. New tests are in progress for confirmation...

# 5 – Acknowledgements

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